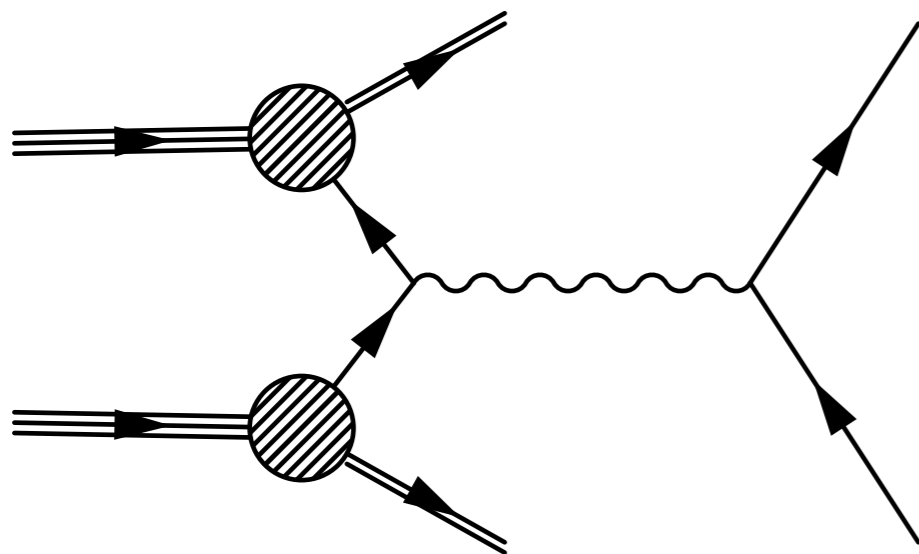


# Precision EW Measurements from ATLAS

## Extracting $\sin^2\theta_{\text{eff}}$



Introduction

Why measure  $\sin^2\theta_{\text{eff}}$  ?

New triple-diff<sup>l</sup> Drell-Yan Cross Sections  $d^3\sigma$

Systematic Uncertainties

Extraction of  $\sin^2\theta_{\text{eff}}$

Measurement of the Drell-Yan triple differential cross section

in pp collisions at  $\sqrt{s} = 8$  TeV

[http://dx.doi.org/10.1007/JHEP12\(2017\)059](http://dx.doi.org/10.1007/JHEP12(2017)059)

[arXiv:1710.05167](https://arxiv.org/abs/1710.05167)

HepData tables:

<https://www.hepdata.net/record/ins1630886>

Eram Rizvi



UCL Seminar  
26<sup>th</sup> October 2018



Queen Mary  
University of London



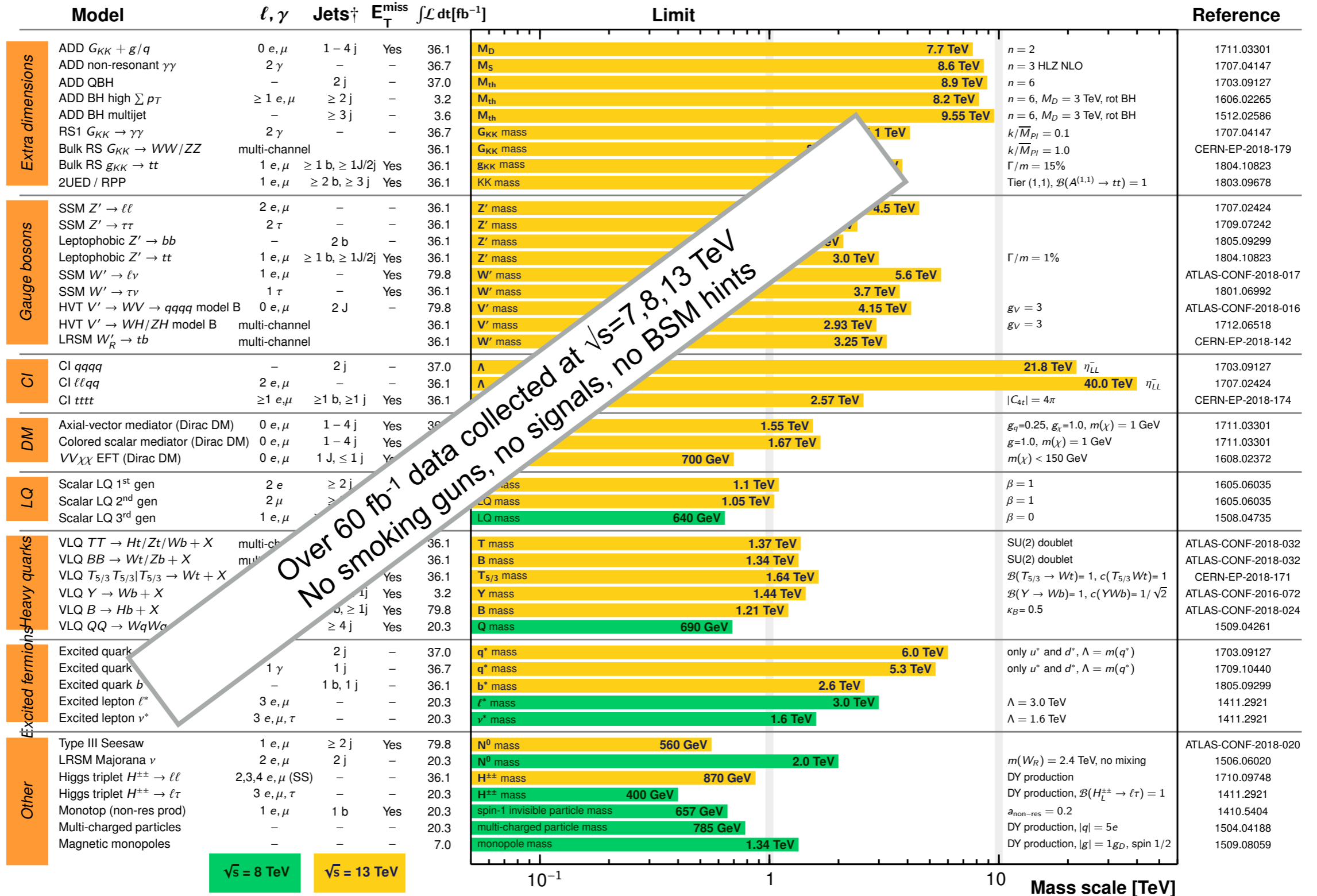
## ATLAS Exotics Searches\* - 95% CL Upper Exclusion Limits

Status: July 2018

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 79.8) \text{ fb}^{-1}$$

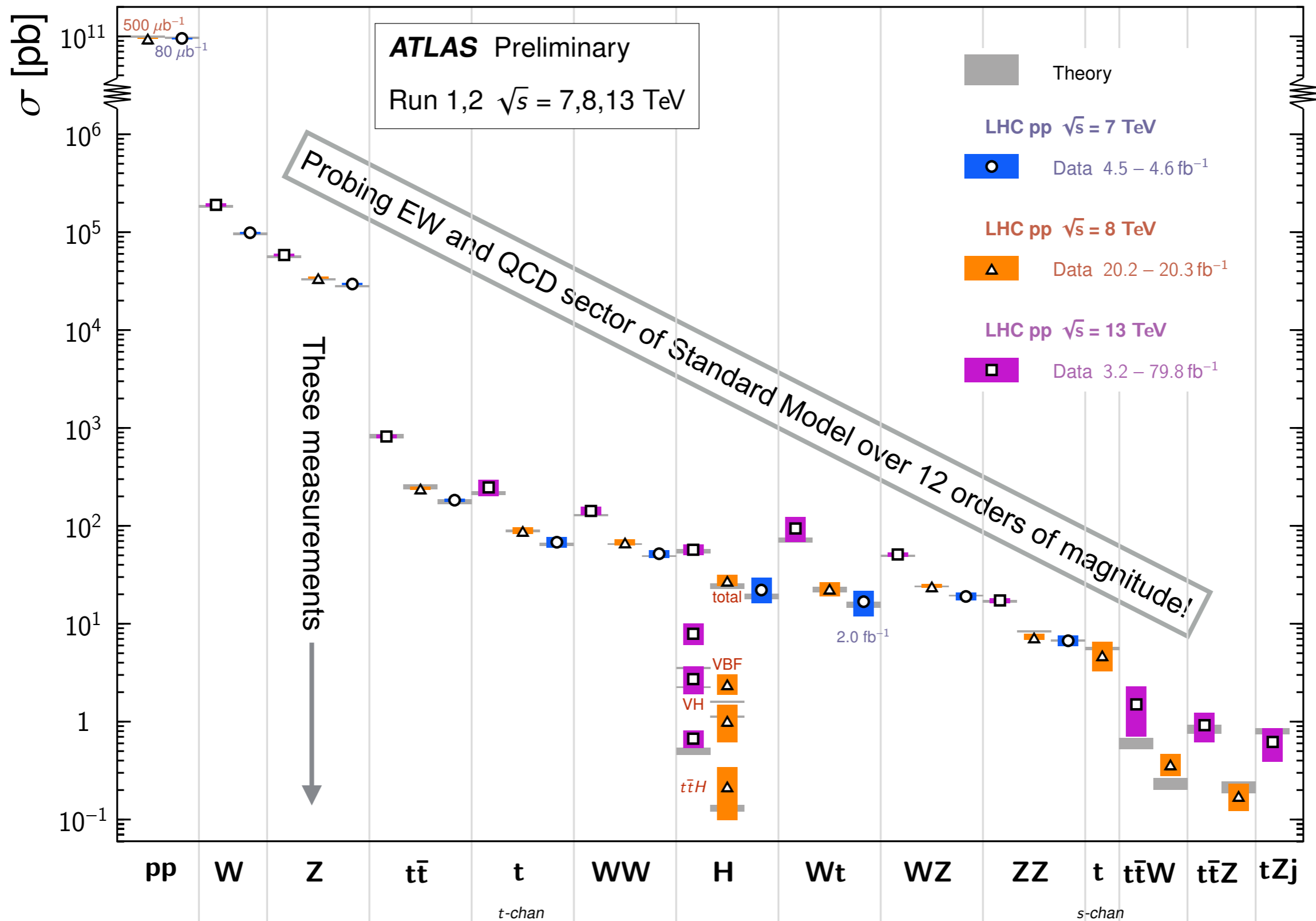
$$\sqrt{s} = 8, 13 \text{ TeV}$$



\*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).

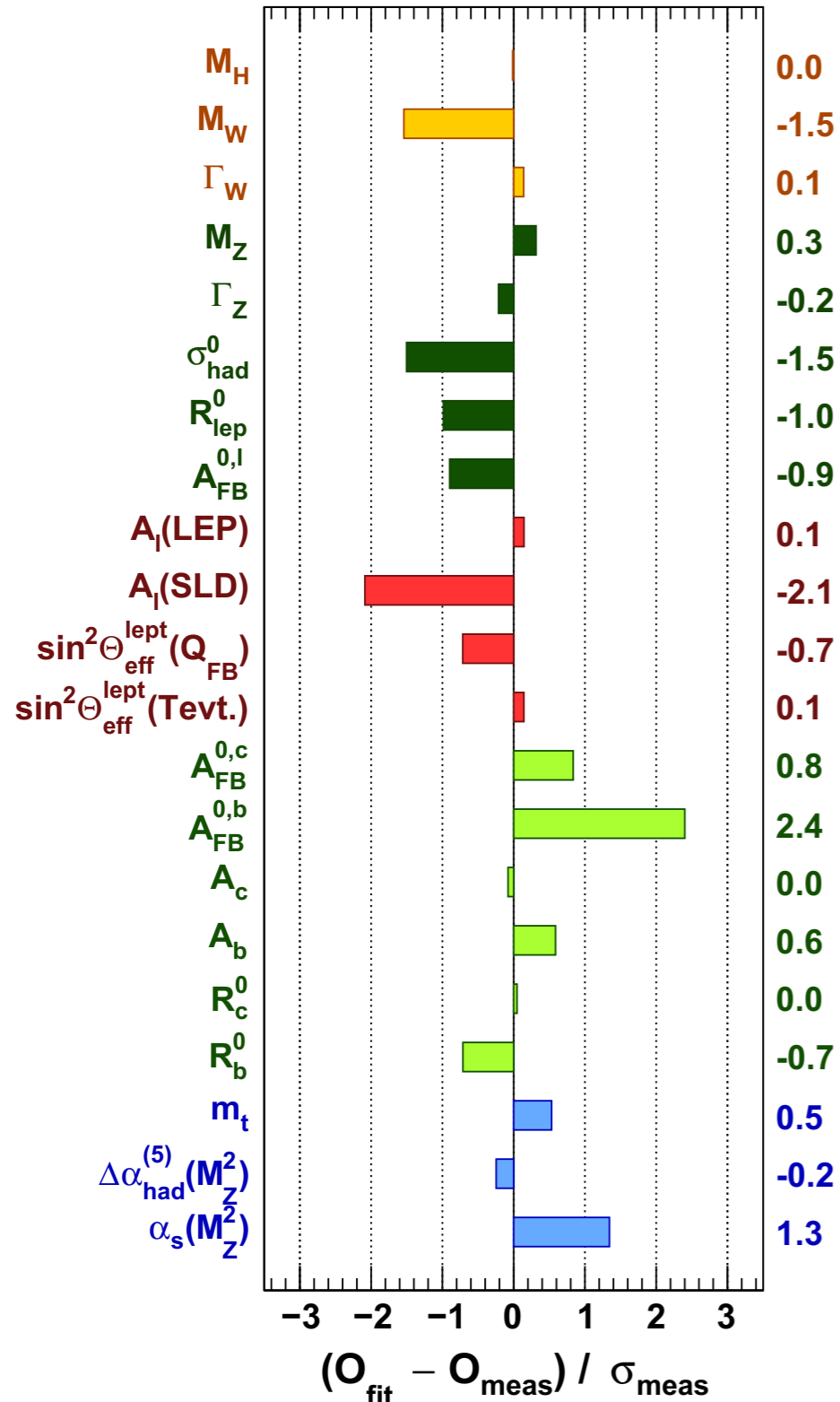
## Standard Model Total Production Cross Section Measurements Status: July 2018





Gfitter 2018

Global EW fit of all precision data



With known  $m_h$  EW sector of SM is over-constrained

- $m_Z = 91.1876 \text{ GeV}$
- $G_\mu = 1.16637 \times 10^{-5} \text{ GeV}^{-2}$
- $\alpha_{\text{QED}}(0) = 1/137.035$
- .... several others ....

$\sin^2\theta_W$  is a fundamental SM parameter of the SM  
 Specifies the mixing between EM and weak fields  
 Relates the Z and W couplings  $g_Z$  and  $g_W$  (and their masses)

$$\text{At leading order } \sin^2\theta_W = 1 - \frac{g_W^2}{g_Z^2} = 1 - \frac{m_W^2}{m_Z^2}$$

Higher order EW corrections modify this  
 to an effective mixing angle  
 dependent on fermion flavour  $f$

$$\sin^2\theta_{\text{eff}}^f = \left(1 - \frac{m_W^2}{m_Z^2}\right) \cdot (1 + \Delta r)$$



EW scheme dependent  
 corrections incorporated into  
 $\Delta r \rightarrow \Delta r(m_H, m_{\text{top}}, \dots)$

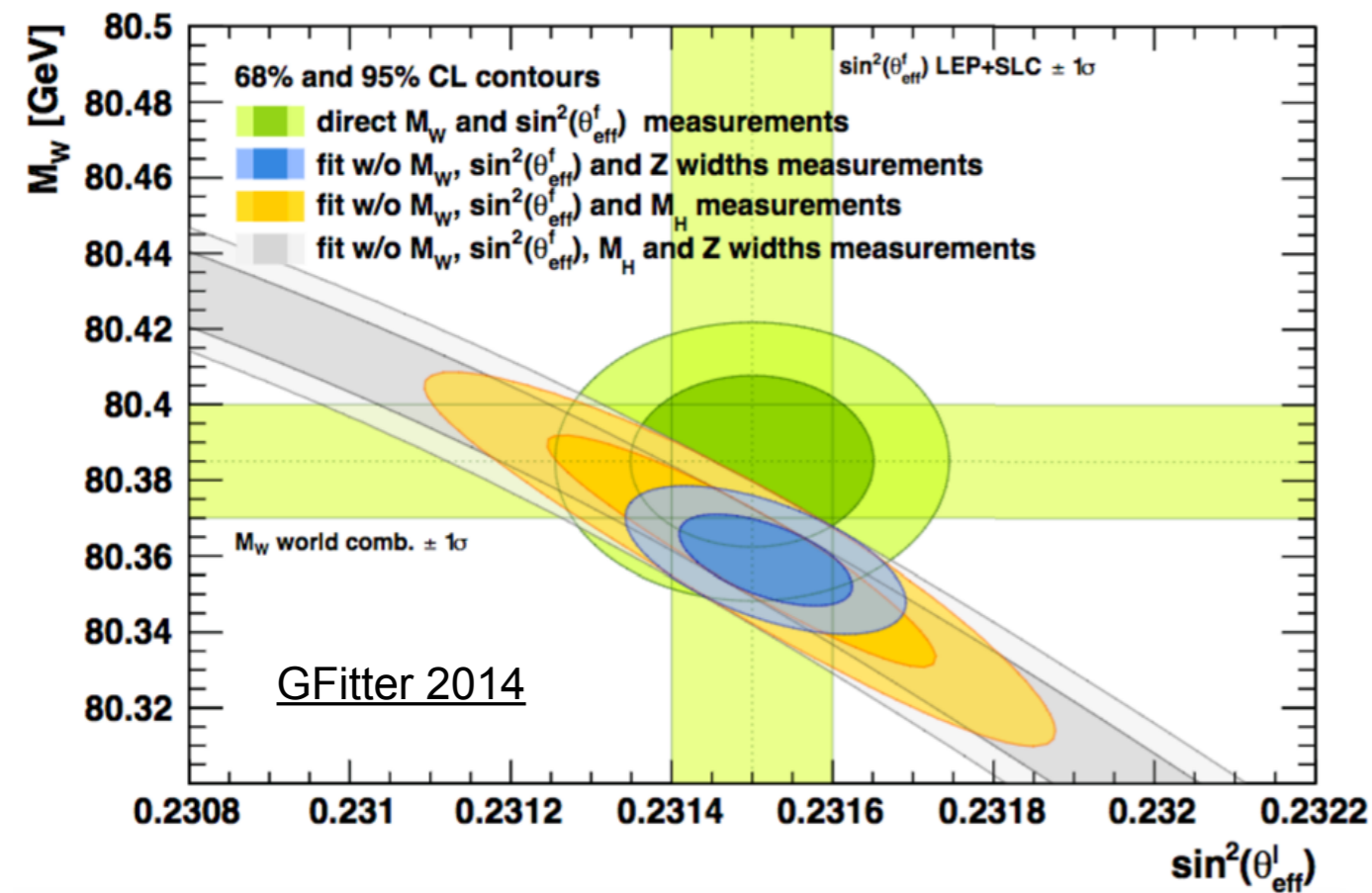


$$\sin^2\theta_{\text{eff}}^f = \left(1 - \frac{m_W^2}{m_Z^2}\right) \cdot (1 + \Delta r)$$

EW scheme dependent corrections incorporated into  $\Delta r \rightarrow \Delta r(m_H, m_{\text{top}}, \text{new physics})$



In context of EFT extension to SM  
EW oblique parameters S, T, U, Y, W  
incorporate new BSM dim-6 operators  
in self-energy terms



Measurement of one observable can predict the other  
 $m_W \Leftrightarrow \sin^2\theta_W$

$$m_W^2 = \frac{\pi\alpha(0)}{\sqrt{2}G_\mu \sin^2\theta_W} \frac{1}{1 - \Delta r}$$

$m_W$  and  $\sin^2\theta_{\text{eff}}$  allows self-consistency check of SM  
New physics hidden in the higher order corrections ??  
Valuable test in absence of direct BSM signals

### Final Precision on $\sin^2\theta_{\text{eff}}$

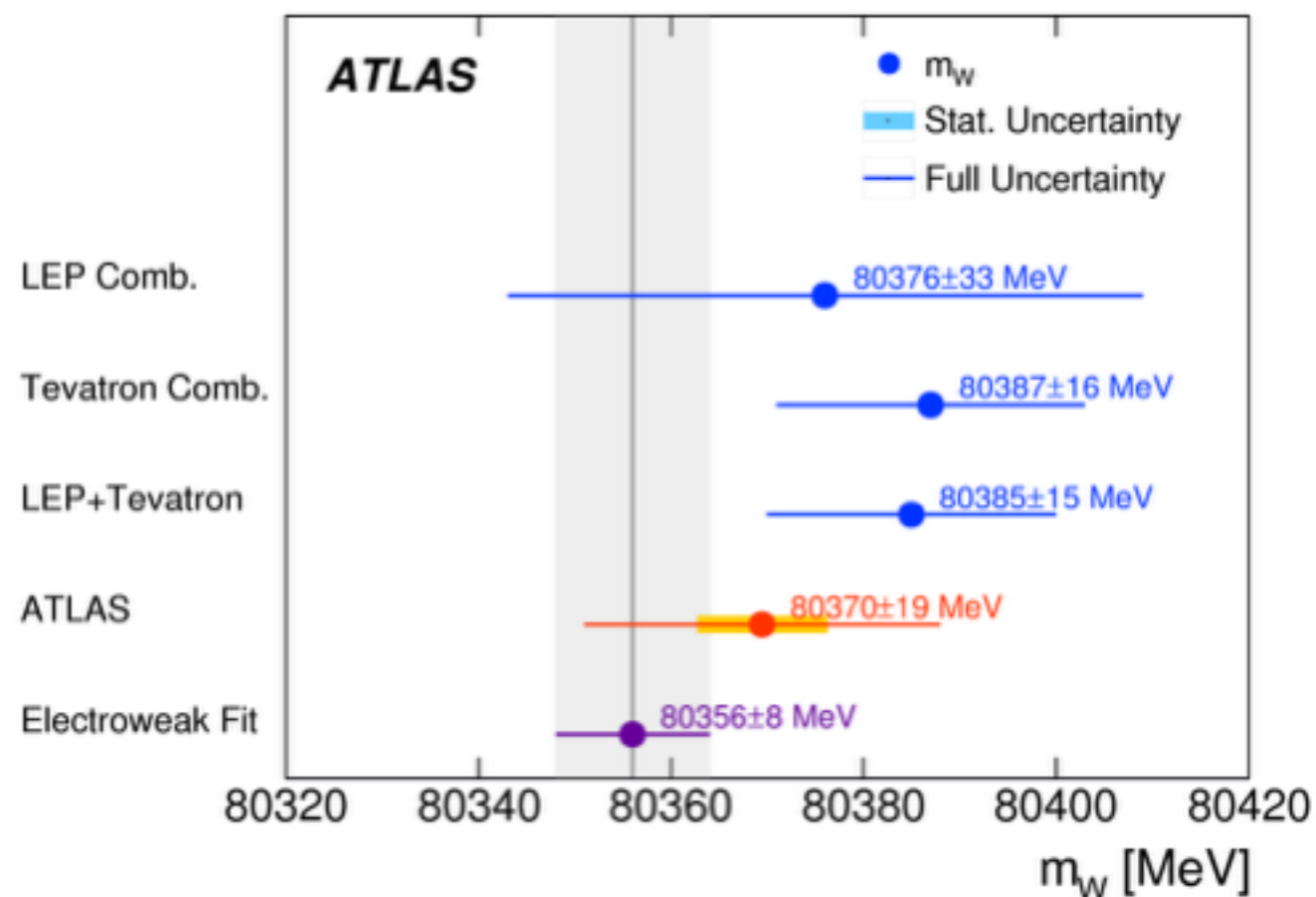
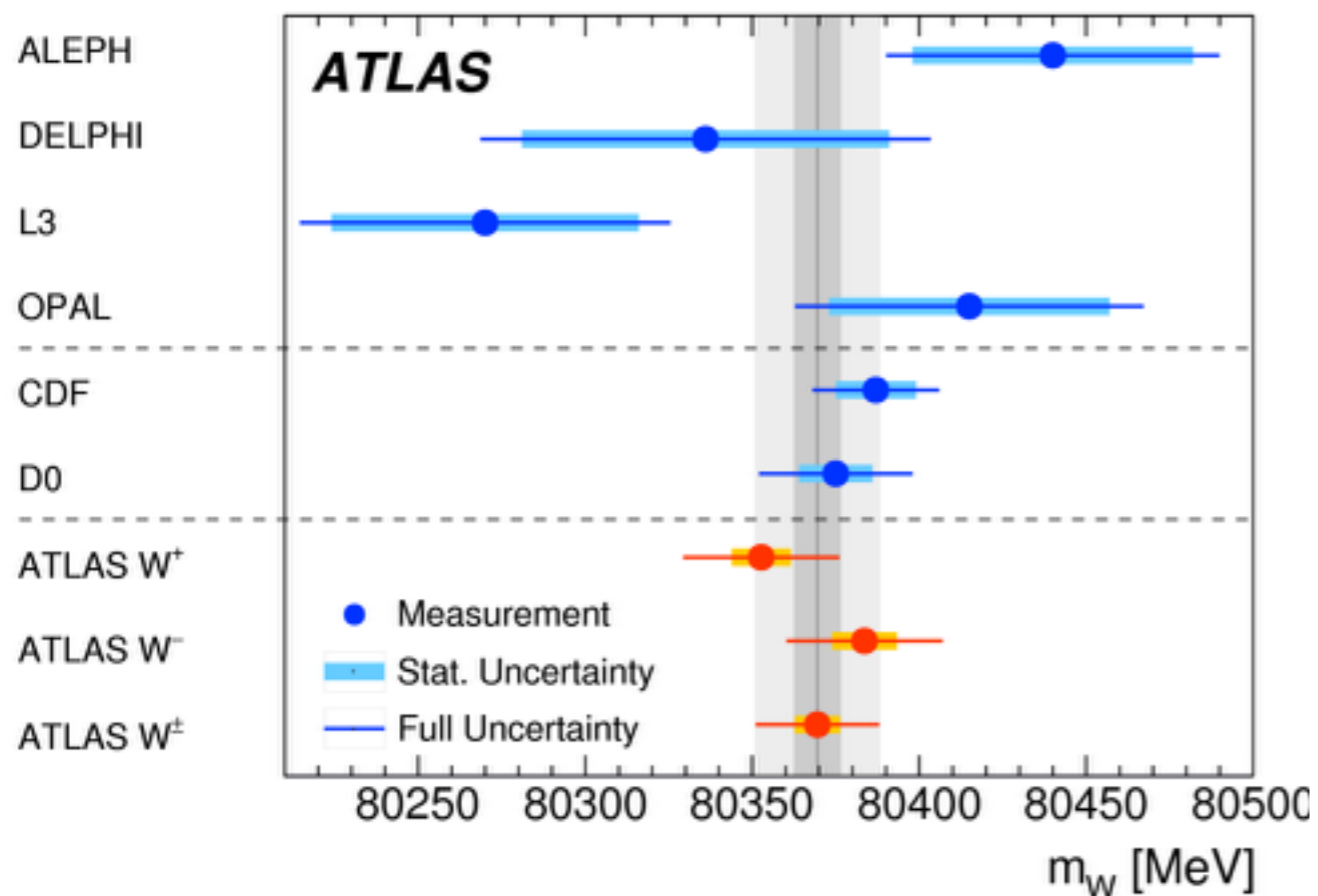
LEP:  $\pm 29 \times 10^{-5}$   
SLD:  $\pm 26 \times 10^{-5}$   
CDF+D0:  $\pm 35 \times 10^{-5}$

### First LHC results on $\sin^2\theta_{\text{eff}}$

CMS(7TeV):  $\pm 320 \times 10^{-5}$   
ATLAS(7TeV):  $\pm 120 \times 10^{-5}$

$\sin^2\theta_{\text{eff}}$  precision  $\pm 50 \times 10^{-5}$  equivalent to  $\pm 25$  MeV in  $m_W$

New ATLAS measurement of  $m_W$  reaches  $\pm 19$  MeV precision



ATLAS approaches precision of combined LEP + Tevatron measurement  
Theory prediction from EW fit has uncertainty  $\pm 8$  MeV

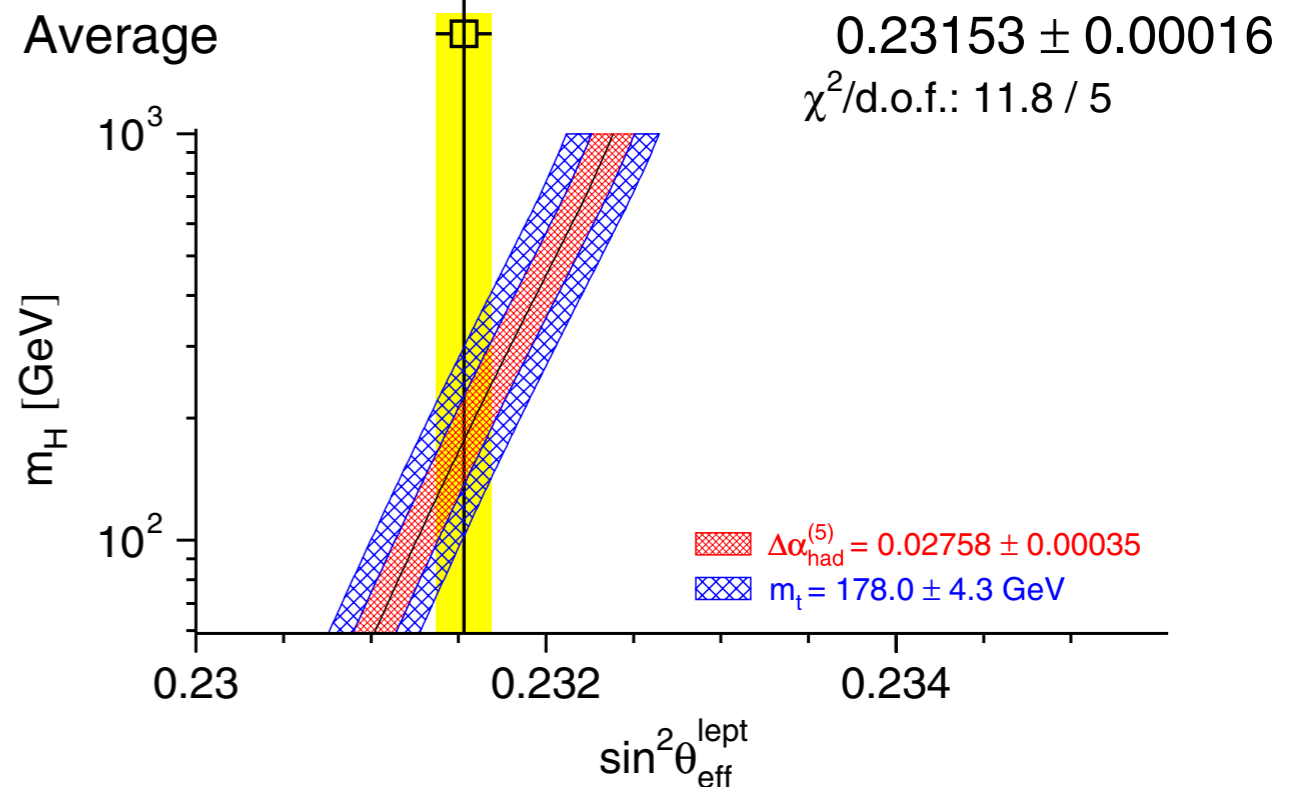
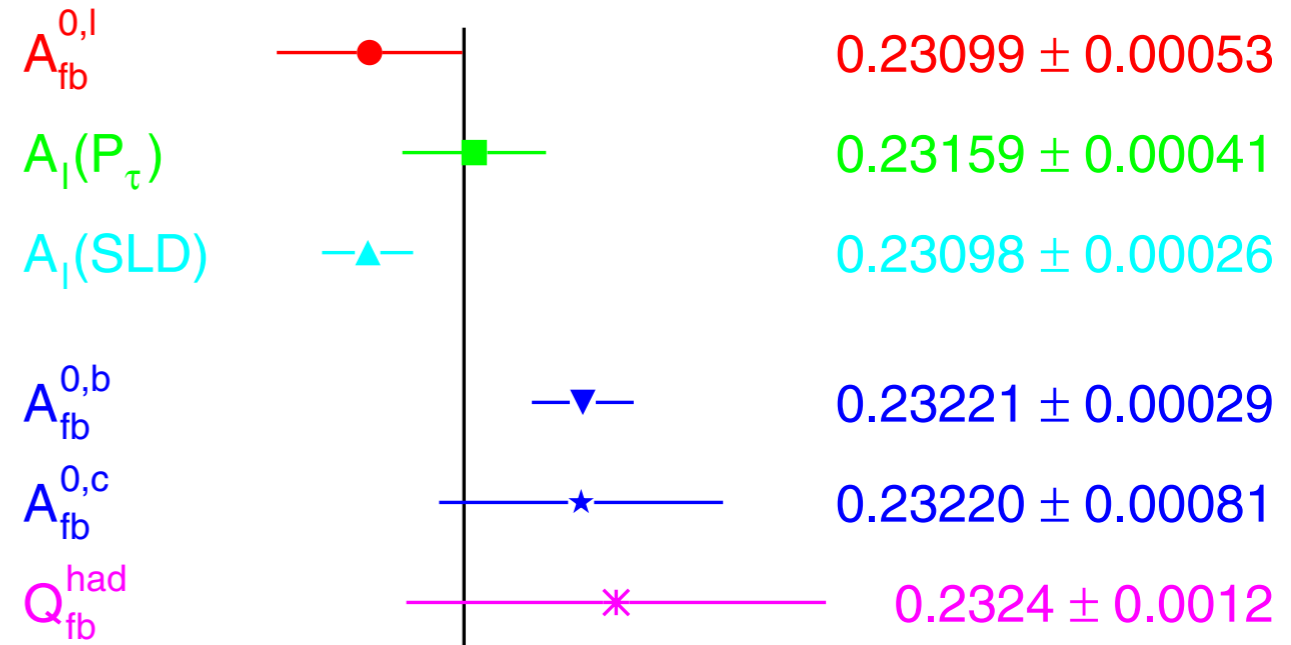


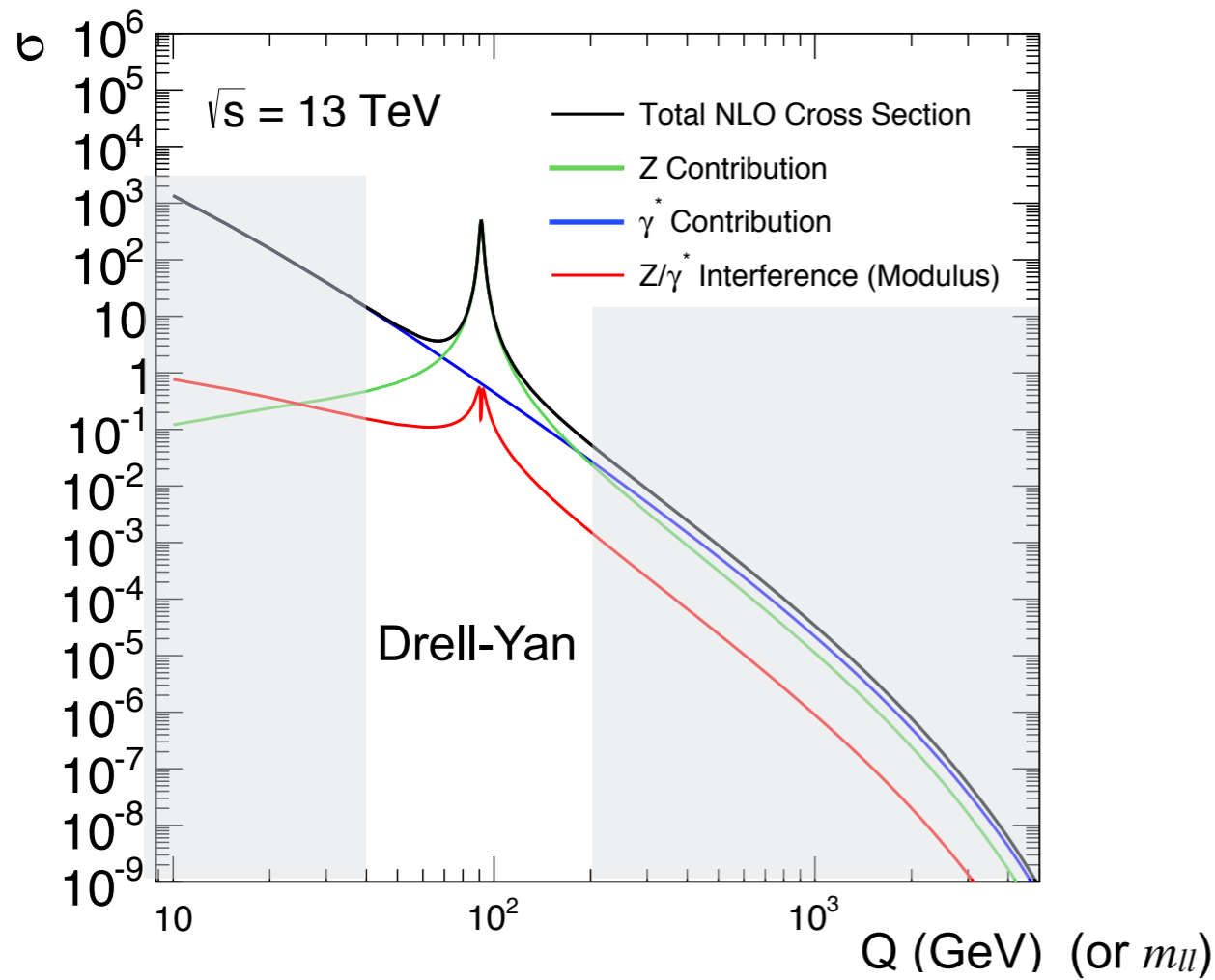
Previous generation of  $\sin^2\theta_W$  measurements LEP/SLD  
 Several different observables and asymmetries used

$A_l$  = polarisation L/R asymmetry at SLD

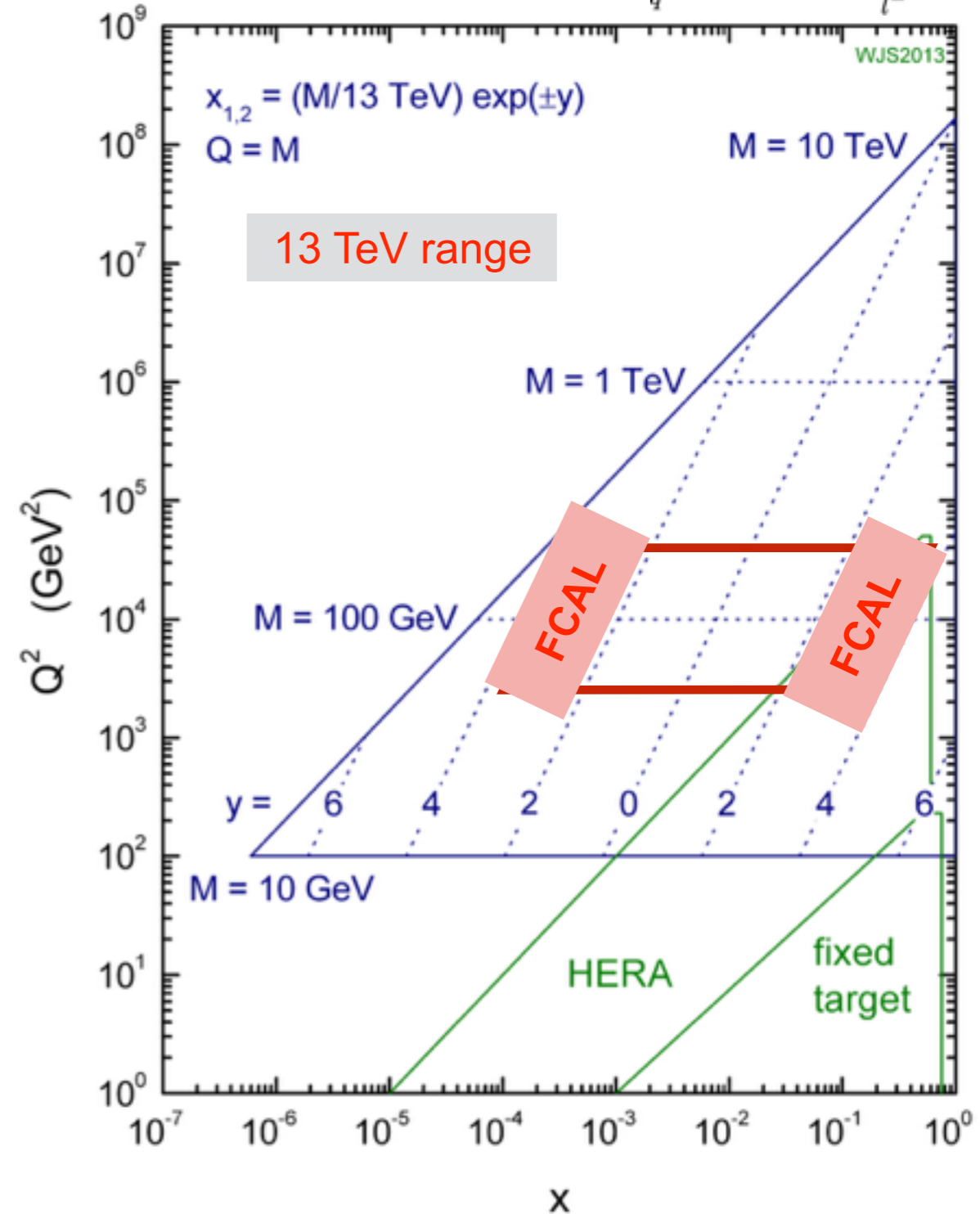
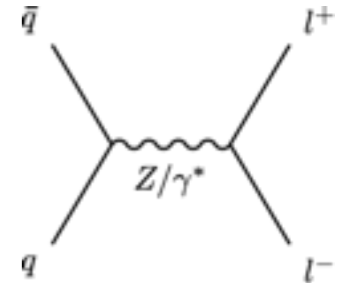
$A_{\text{FB}}^{0,b}$  = forward/backward asymmetry in  $Z \rightarrow b\bar{b}$

Long-standing  $3.2\sigma$  discrepancy between LEP and SLD





Drell—Yan cross section  
Large event rates and clean final-state  
allow high experimental precision



Measure triple differential cross sections:  $\frac{d^3\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|d\cos\theta^*}$

Can be integrated to derive  $\frac{d\sigma}{dm_{\ell\ell}}$   $\frac{d^2\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|}$

Measurements access range of

$$x > 4 \times 10^{-4}$$

$$0 < |y| < 3.6$$

$$46 \leq m \leq 200 \text{ GeV}$$





leptonic decay angle in Collins-Soper frame  $\cos \theta^* = \frac{p_{z,\ell\ell}}{m_{\ell\ell}|p_{z,\ell\ell}|} \frac{p_1^+ p_2^- - p_1^- p_2^+}{\sqrt{m_{\ell\ell}^2 + p_{T,\ell\ell}^2}}$

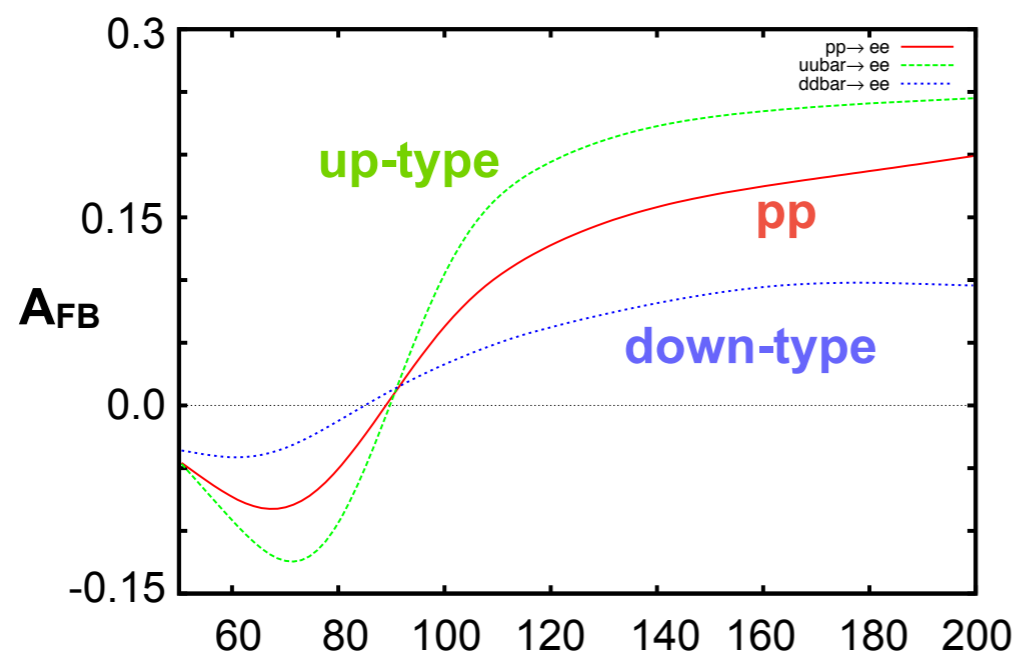
lepton and quark angle or anti-lepton / anti-quark angle

$$\frac{d^3\sigma}{dm_{\ell\ell} dy_{\ell\ell} d\cos\theta^*} = \frac{\pi\alpha^2}{3m_{\ell\ell}s} \sum_q P_q \left[ f_q(x_1, Q^2) f_{\bar{q}}(x_2, Q^2) + (q \leftrightarrow \bar{q}) \right]$$

**pure  $\gamma^*$**   $P_q = e_l^2 e_q^2 (1 + \cos^2 \theta^*)$   $f_q(x, Q^2) =$  parton density functions

**interference  $Z/\gamma^*$**   $+ e_l e_q \frac{2m_{\ell\ell}^2 (m_{\ell\ell}^2 - m_Z^2)}{\sin^2 \theta_W \cos^2 \theta_W [(m_{\ell\ell}^2 - m_Z^2)^2 + \Gamma_Z^2 m_Z^2]} [v_\ell v_q (1 + \cos^2 \theta^*) + 2a_\ell a_q \cos \theta^*]$

**pure Z**  $+ \frac{m_{\ell\ell}^4}{\sin^4 \theta_W \cos^4 \theta_W [(m_{\ell\ell}^2 - m_Z^2)^2 + \Gamma_Z^2 m_Z^2]} [(a_\ell^2 + v_\ell^2)(a_q^2 + v_q^2)(1 + \cos^2 \theta^*) + 8a_\ell v_\ell a_q v_q \cos \theta^*]$



forward =  $\cos \theta^* > 0$  Asymmetry  
backward =  $\cos \theta^* < 0$

$$A_{FB} = \frac{d^3\sigma(\cos \theta^* > 0) - d^3\sigma(\cos \theta^* < 0)}{d^3\sigma(\cos \theta^* > 0) + d^3\sigma(\cos \theta^* < 0)}$$

Sensitive to  $\sin^2 \theta_W$   
Sensitive to PDFs  $f(x, Q^2)$

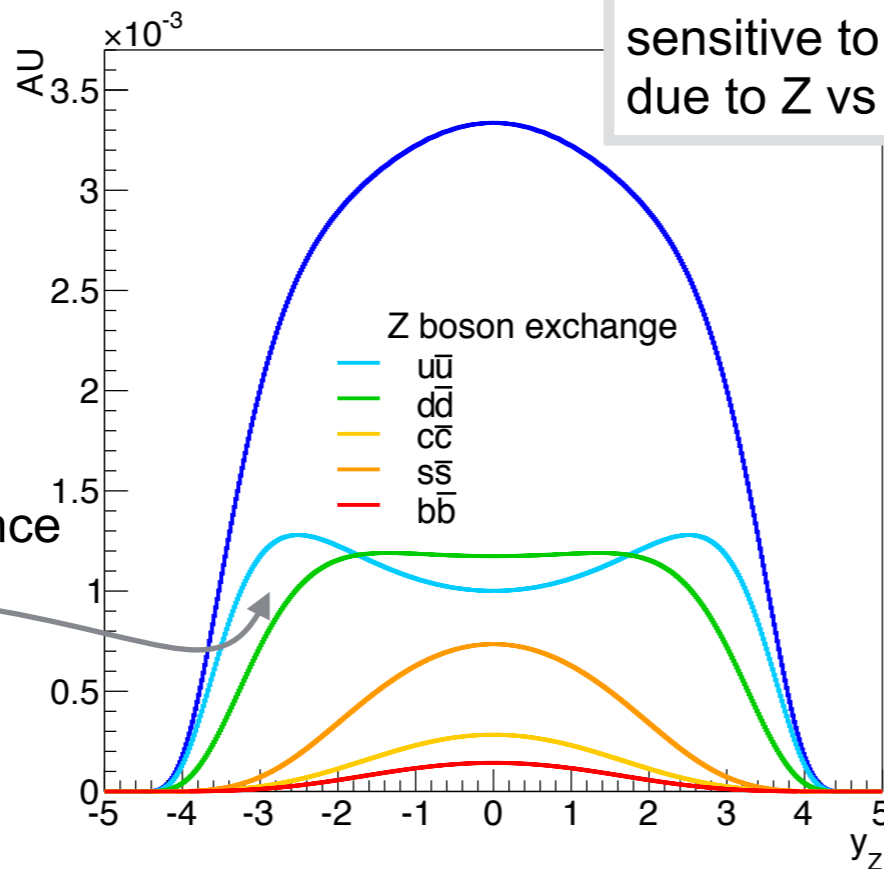
Zll vertex more sensitive than Zqq to  $\sin^2 \theta_W$   
(by factor  $\sim 5u - 20d,s$ )

# Triple-differential Z/ $\gamma^*$ Measurement Motivation

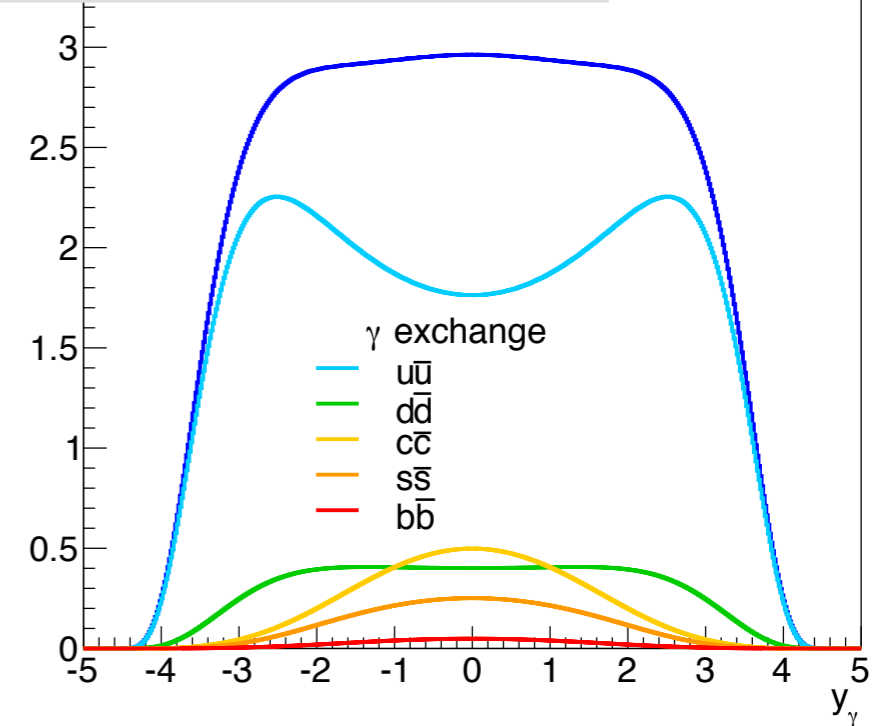


In different  $m$  regions  $y$  spectrum shape changes dramatically for  $m_{ll} \neq m_Z$

Sensitivity to  $u_v$  &  $d_v$  valence quarks at  $|y| > 1$



sensitive to difference in u-type and d-type due to Z vs  $\gamma^*$  couplings on- and off-shell

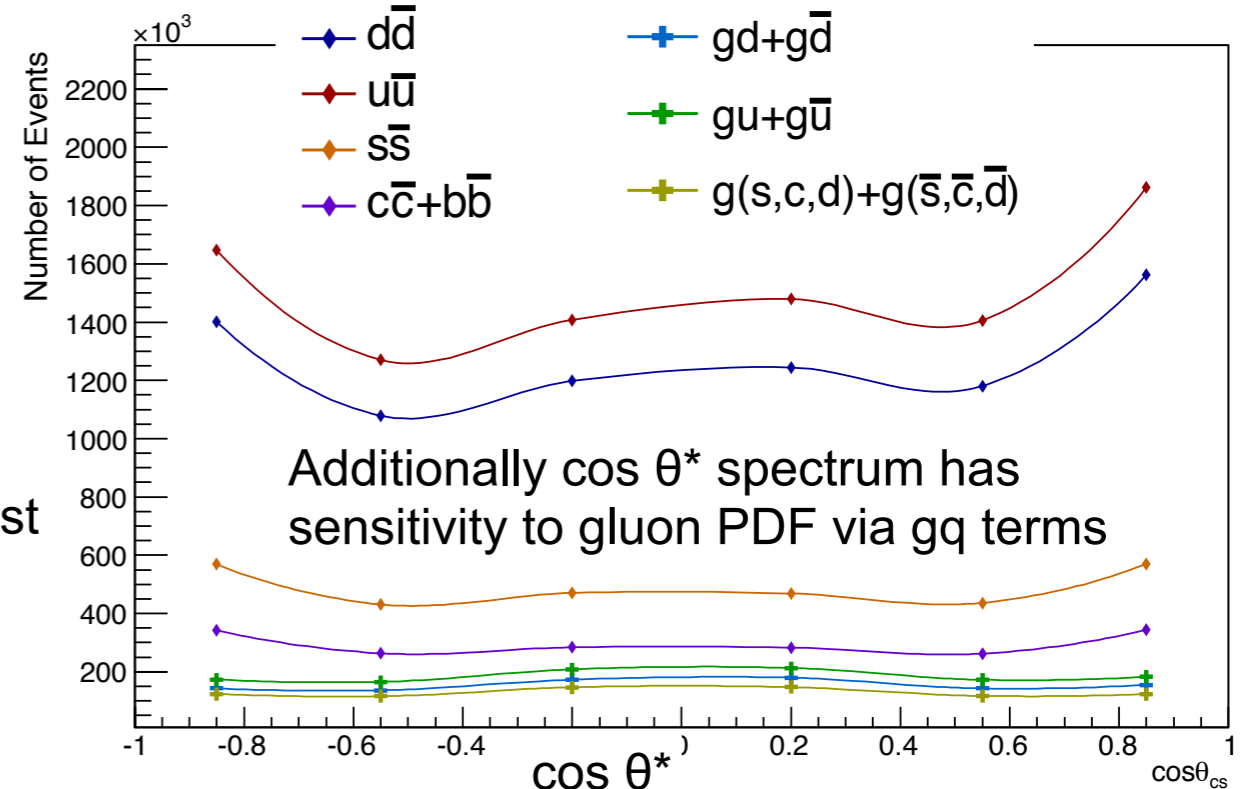


$y$  dependence measures  $x$  distribution of PDFs

$$x_{1,2} = \frac{m_{ll}}{\sqrt{s}} e^{\pm y}$$

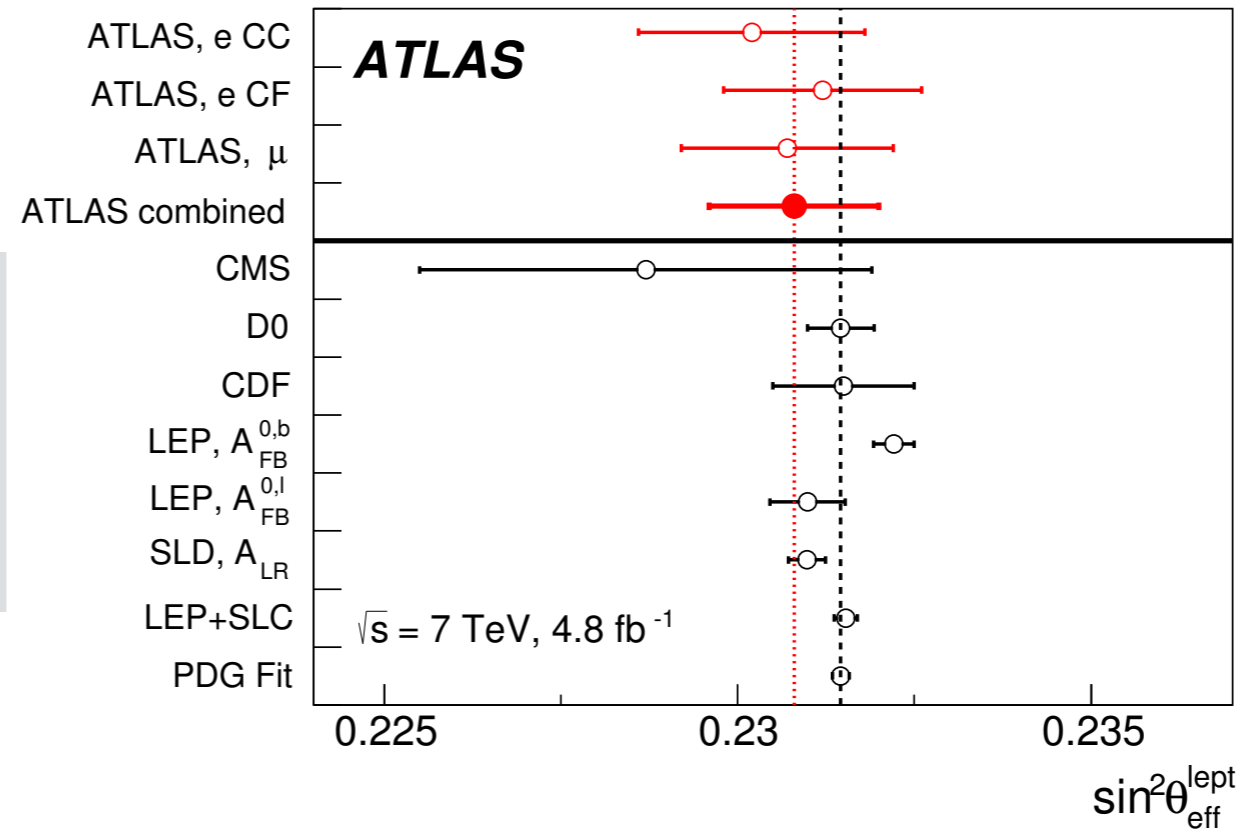
At LHC direction of incoming quark is unknown  
 Therefore there is ambiguity in defining  $\theta^*$   
 (not a problem at Tevatron)  
 Ambiguity dilutes  $A_{FB}$   
 Dilution is reduced at large  $|y|$  due to valence quark boost

$\Rightarrow$  greater sensitivity to  $\sin^2\theta_{eff}$  at larger  $y$   
 zero sensitivity at  $y=0$





Previous ATLAS measurement of  $\sin^2 \theta_W$   
JHEP09(2015)049  
 $5 \text{ fb}^{-1}$   
 $\sqrt{s} = 7 \text{ TeV}$



Uncertainty Source	CC electrons $\times 10^{-5}$	CF electrons $\times 10^{-5}$	Muons $\times 10^{-5}$	Combined $\times 10^{-5}$
PDF	100	100	90	90
MC statistics	50	20	50	20
Elec energy scale	40	60	—	30
Elec energy res.	40	50	—	20
Muon energy scale	—	—	50	20
higher order corrs	30	10	30	20
Other source	10	10	20	20

ATLAS measurement of  $\sin^2 \theta_{\text{eff}}$   
 limited by PDF uncertainty  
 Values are in units of  $10^{-5}$  of  $\sin^2 \theta_{\text{eff}}$

# Run-I Measurements from ATLAS

triple-differential cross sections  $d^3\sigma = \frac{d^3\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|d\cos\theta^*}$

On-shell DY 8 TeV  
 Neutral current - e &  $\mu$  channels  
 $46 < m < 200$  GeV  
 Extended to high y with FCAL analysis

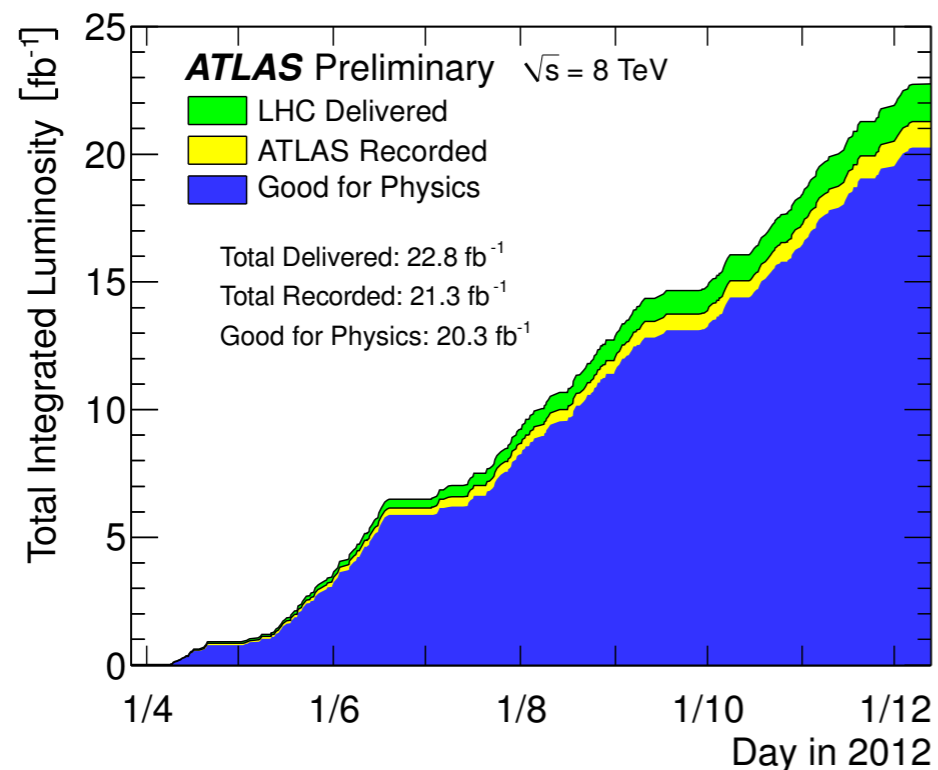
[arXiv:1710.05167](https://arxiv.org/abs/1710.05167)  
[Hepdata](#)

Use  $d^3\sigma$  to derive  
 ancillary measurements  
 for purely visual purposes

$A_{FB}(m,|y|)$

$$\frac{d\sigma}{dm_{\ell\ell}}$$

$$\frac{d^2\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|}$$



Complete 2012 data set analysed

Centre of mass energy  $\sqrt{s} = 8$  TeV

$$\int \mathcal{L} dt = 20.2 \text{ fb}^{-1}$$

7M di-electron events (CC)  
 9M di-muon events (CC)  
 1M forward di-electron events (CF)

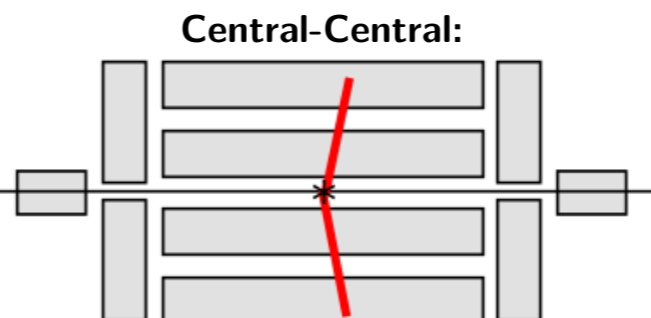
## Muon Selection

- $\geq 2$  isolated muons
- muon  $|\eta| < 2.4$
- muon  $p_T > 20$  GeV
- opposite charge

## Central Electron Selection

- $\geq 2$  good quality “medium” electrons
- electron  $|\eta| < 2.4$  excl.  $1.37 < |\eta| < 1.52$
- electron  $E_T > 20$  GeV

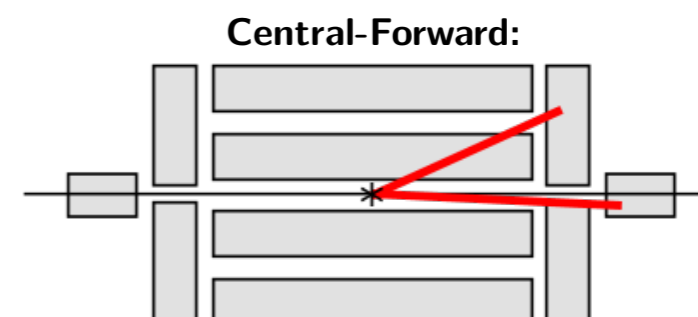
Central-central topology (CC)  
Two leptons with  $|\eta| < 2.4$



## Forward Electron Selection

- 1 good quality “tight” central electron
  - electron  $|\eta| < 2.47$  excl.  $1.37 < |\eta| < 1.52$
  - electron  $E_T > 25$  GeV
- **1 good quality “tight” forward electron**
  - **electron  $2.5 < |\eta| < 4.9$**  excl.  $3.0 < |\eta| < 3.4$
  - electron  $E_T > 20$  GeV

Central-forward topology (CF)  
One electron with  $|\eta| < 2.4$   
One electron with  $2.5 < |\eta| < 4.9$



Identical dataset and almost identical selection as ATLAS angular coefficients analysis (see later)

**Already good precision achieved for run-II !**

Need to ensure phase-space corners are well covered e.g. boosted Zs access high pT lepton efficiencies  
For run-I lepton pT  $\sim 200$  GeV  
(For run-II we should reach lepton pT  $\sim 400$  GeV)

## Electron Channel

Energy scale typically  $<1\%$  in peak region  
dominates error at large  $|\cos \theta^*|$   
 $\rightarrow \sim 3\%$

efficiency error typically  $<0.5\%$  in peak region  
larger at large  $\cos \theta^*$  (even at small  $|y|$ )  $\rightarrow \sim 2-3\%$

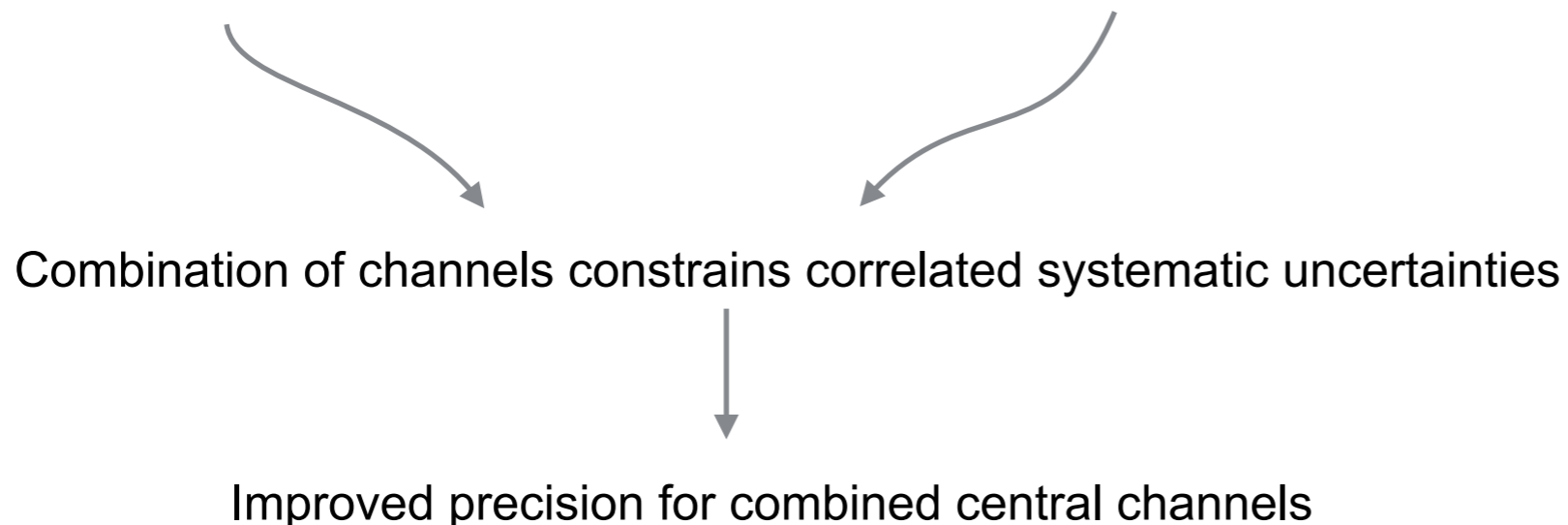
## Muon Channel

In peak region at  $m \sim m_Z$   
momentum scale dominates sys  
error  $\rightarrow \sim 0.6\%$   
compared to  $0.8\%$  stat error

Tracking misalignments  $<1\%$   
upto  $2\%$  at small  $\cos \theta^*$  or large  $y$

## High Rapidity Electron Channel

Energy scale / resolution dominates error at large  $|\cos \theta^*|$  &  $y$   
 $\rightarrow \sim 5\%$  compared to  $\sim 3\%$  stat error

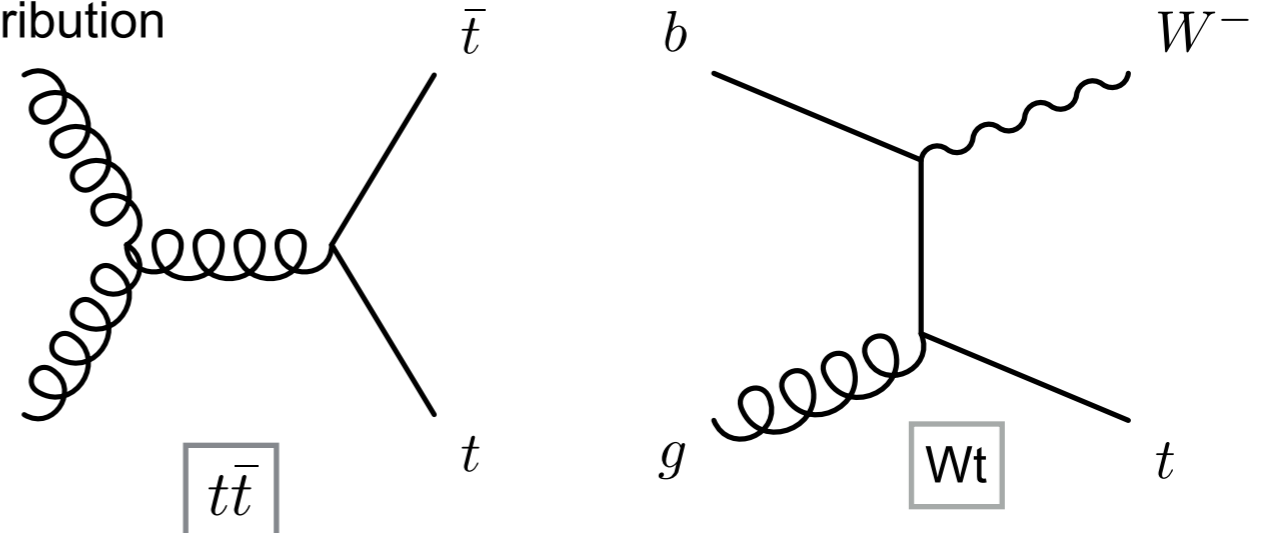


Several sources of so-called “electroweak” backgrounds yielding isolated leptons pairs:

DY → tau production modes found to be negligible contribution

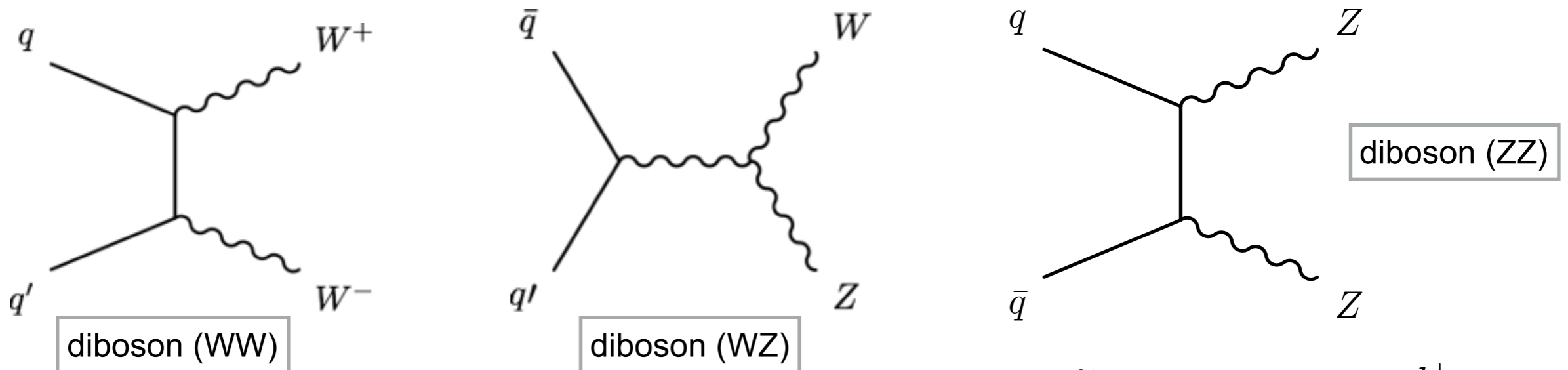
## top background

2-10% contribution top (largest at high  $\cos \theta^*$ )  
 below 5% for high  $y$  channel  
 background estimated from MC



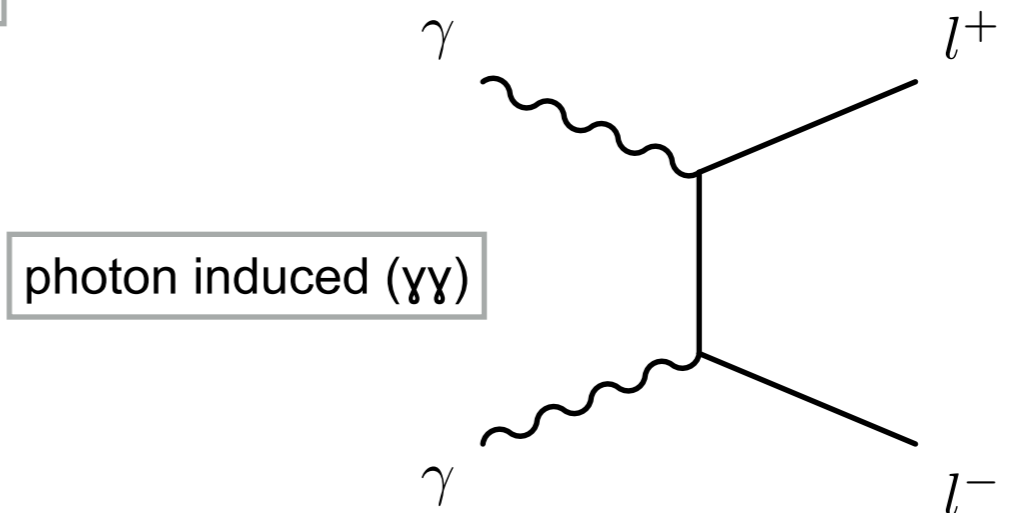
## diboson background

3-6% contribution  
 estimated from MC



## photon induced background

2-5% contribution (largest at large  $m$ )  
 background estimated from MC





## multijet background

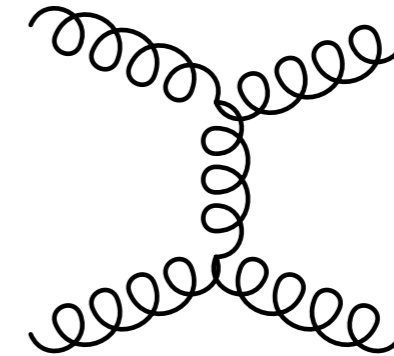
multijet production has large cross section at LHC  
 contributes to background via:

- b,c quark leptonic meson decays
- misidentification of hadron jet as calorimeter electron

soft leptons produced typically

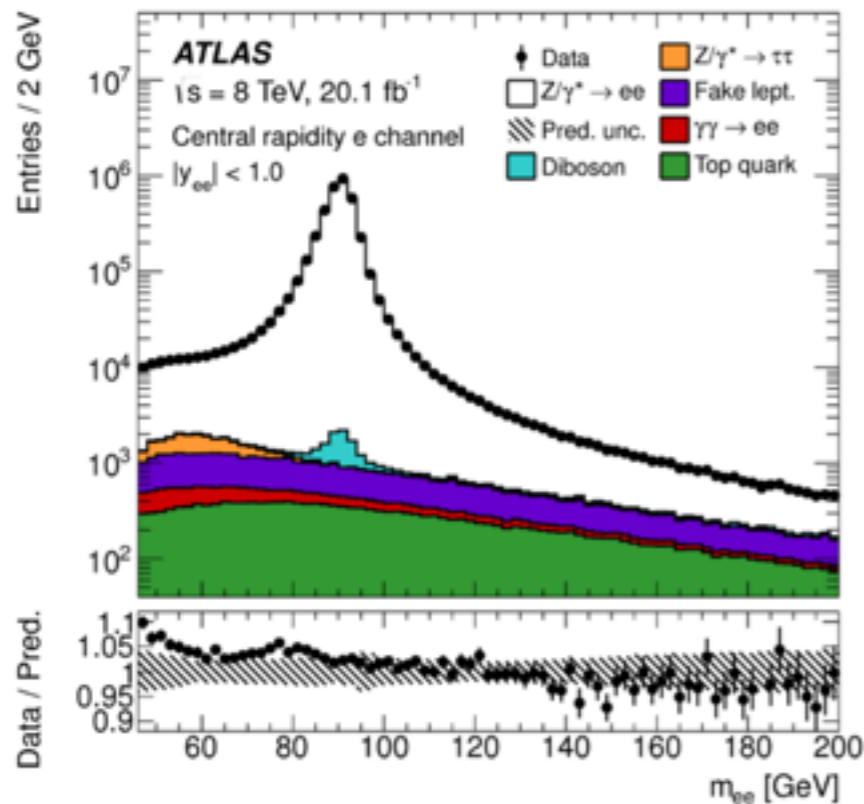
contributing processes involve complex hadronisation simulation

⇒ use data to estimate this background

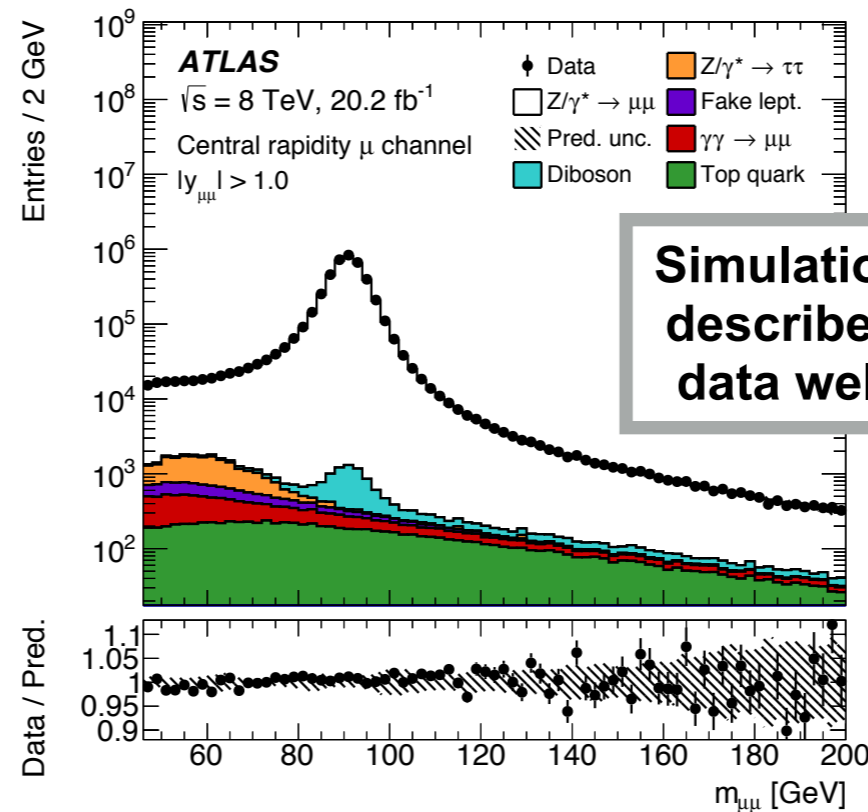


electron / muon channel at  $m \sim m_Z$  b/g is  $< 0.1\%$   
 significant off-peak upto to 15% low  $m_{ee}$   
 less than 5% everywhere in muon channel

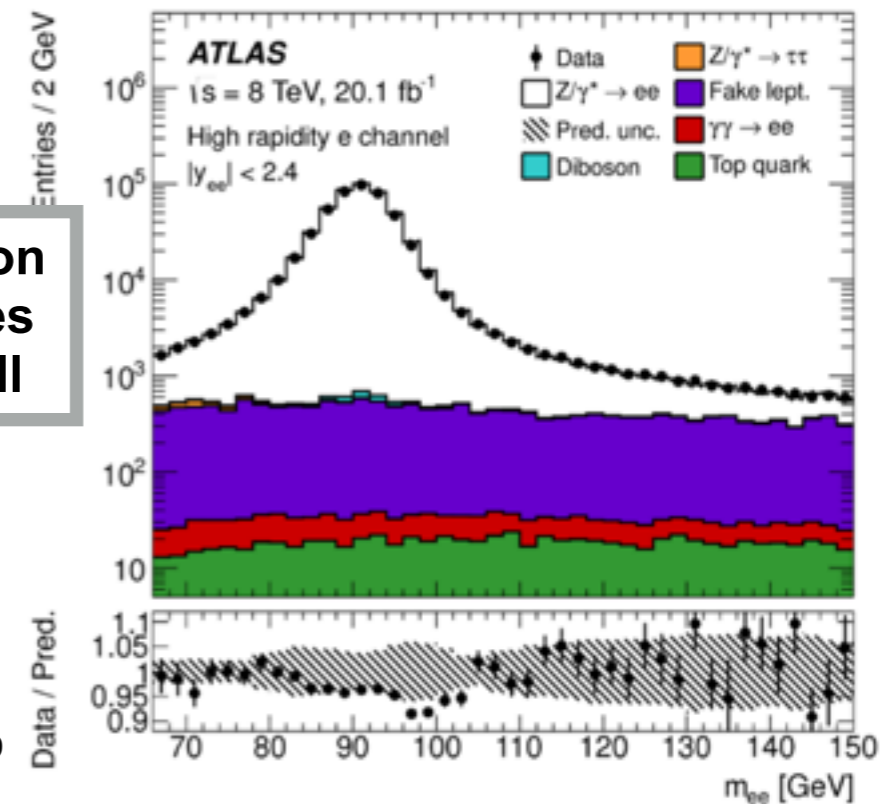
## Electron Channel



## Muon Channel



## High Rapidity Electron Channel





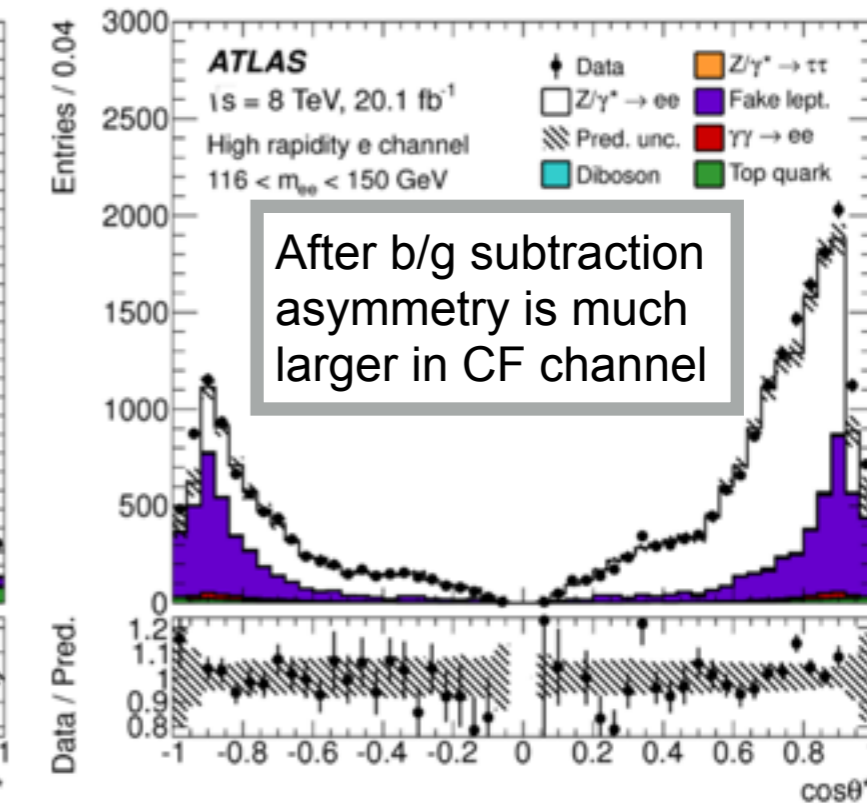
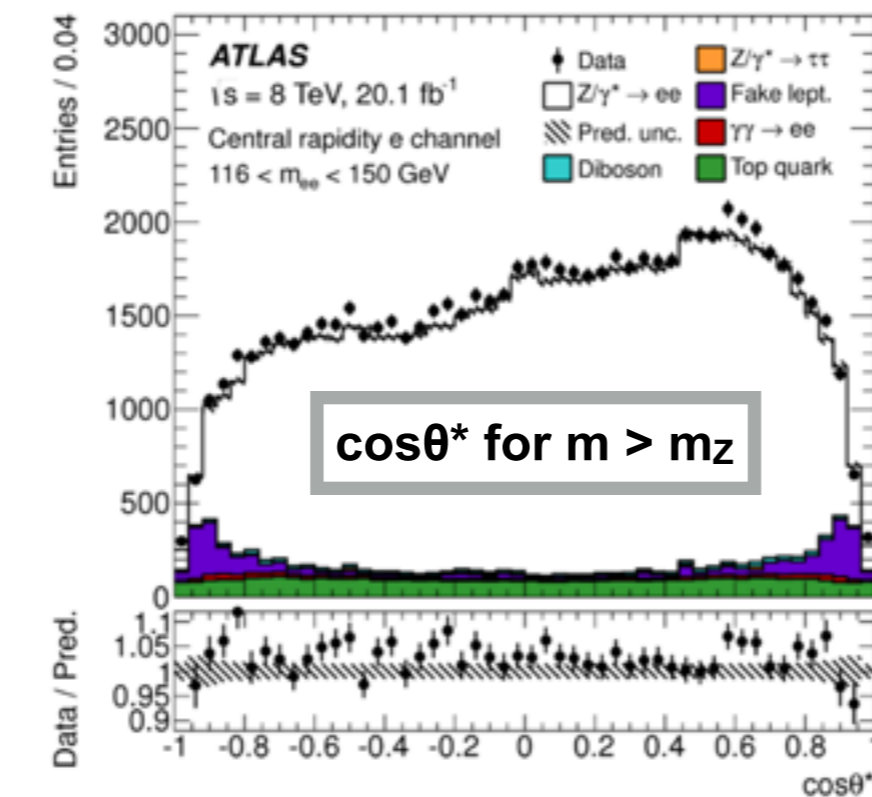
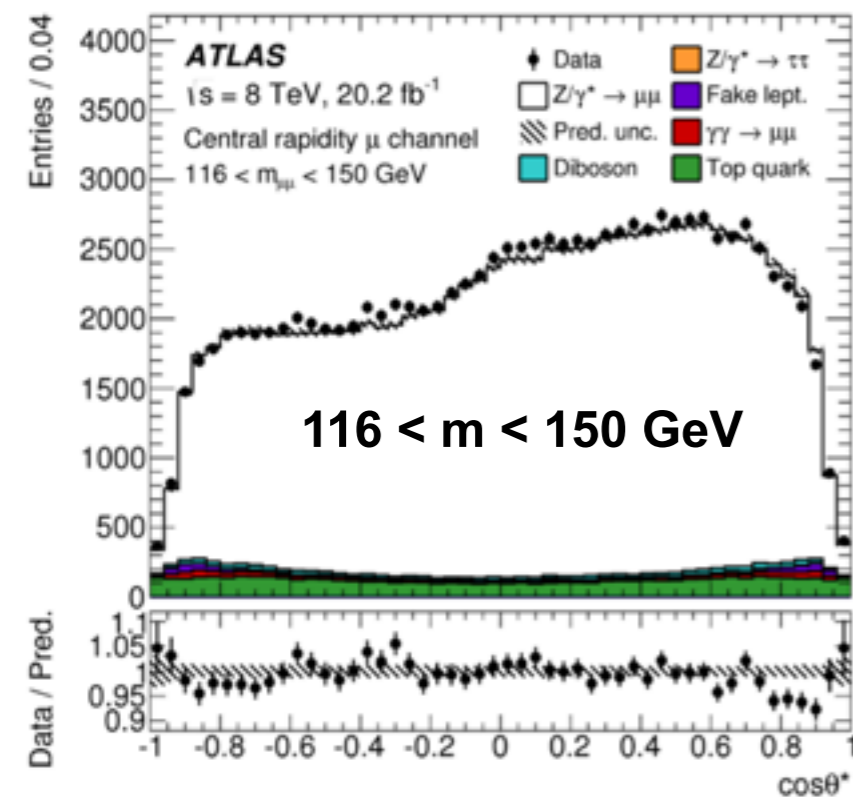
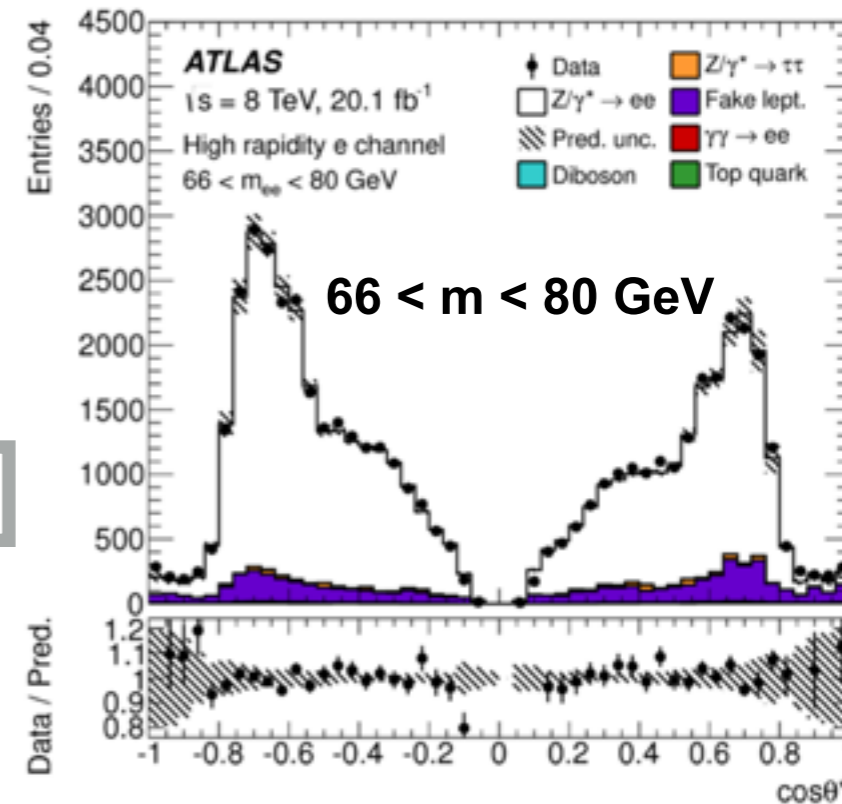
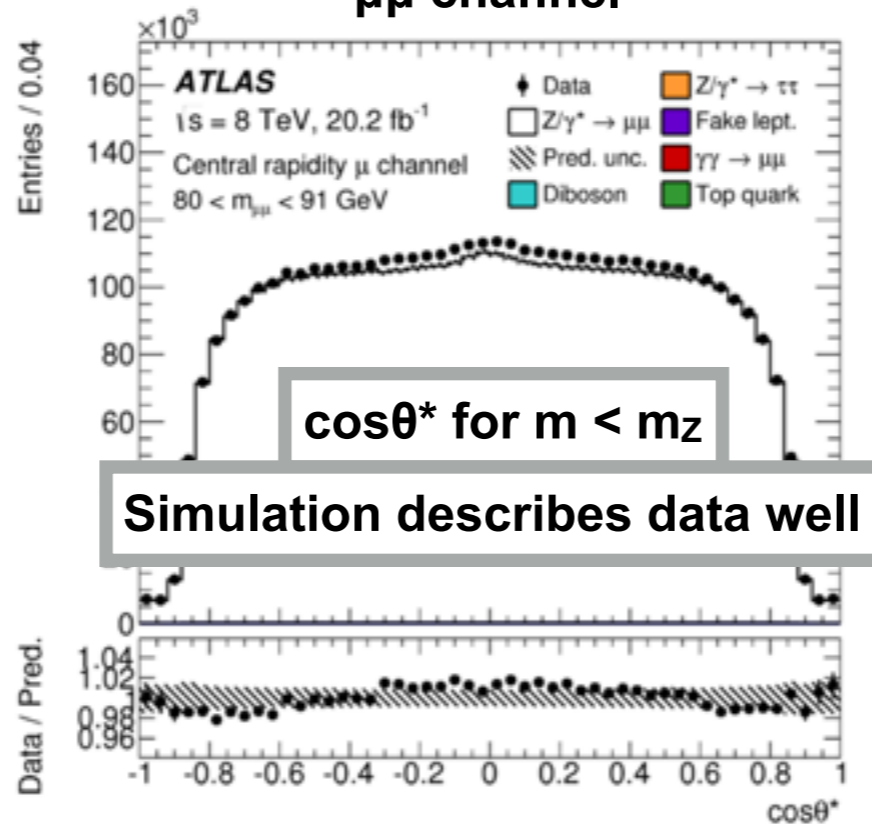
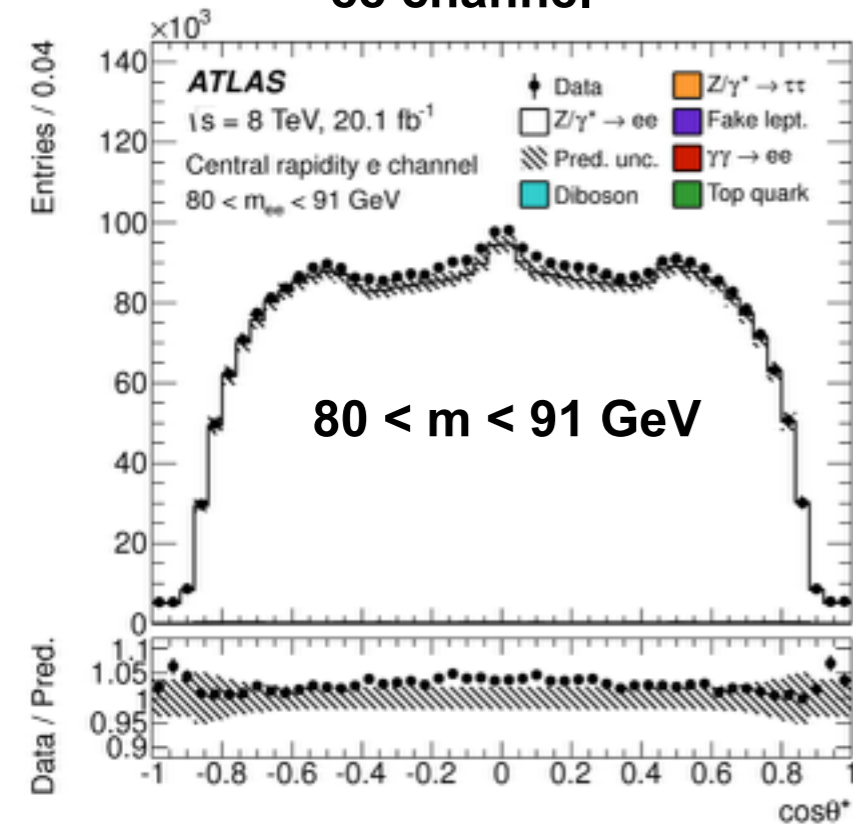
# ATLAS $Z/\gamma^*$ $d^3\sigma$ Cross Section $\sqrt{s} = 8$ TeV



ee channel

$\mu\mu$  channel

forward channel [JHEP12\(2017\)059](https://arxiv.org/abs/1705.059)



<https://link.springer.com/article/10.1007%2FJHEP12%282017%29059>



## Central Rapidity Channel

$m_{  } =$	[46, 66, 80, 91, 102, 116, 150, 200] GeV	7 bins
-----		
$ y_{  }  =$	[0.0, 0.2, 0.4, 0.6, 0.8, 1.0, 1.2, 1.4, 1.6, 1.8, 2.0, 2.2, 2.4]	12 bins
-----		
$\cos \theta^* =$	[-1.0, -0.7, -0.4, 0.0, 0.4, 0.7, 1.0]	6 bins
-----		
	Total bins =	504
	x2 channels	

measure in electron + muon channels  
check for consistency of channels  
combine both measurements  
account for **331** correlated systematic errors  
improved result for both statistical & systematic precision

- Binning choice optimised for
- control experimental bin migrations
  - statistical precision
  - physics sensitivity

## High Rapidity Channel

$m_{  } =$	[66, 80, 91, 102, 116, 150] GeV	5 bins
-----		
$ y_{  }  =$	[1.2, 1.6, 2.0, 2.4, 2.8, 3.6]	6 bins
-----		
$\cos \theta^* =$	[-1.0, -0.7, -0.4, 0.0, 0.4, 0.7, 1.0]	6 bins
-----		
	Total bins =	150



## Unfolding

Remove influence of ATLAS detector by unfolding  
 Use ATLAS detector simulation to quantify event resolution migrations and efficiency losses  
 Define the particle-level phase space of the final quoted result

### CC fiducial cross section definition

- lepton  $p_T > 20$  GeV
- lepton  $|\eta| < 2.5$
- $46 < m_{ll} < 200$  GeV
- Unfolding to Born level lepton kinematics (dressed level available as a correction factor)

### CF fiducial cross section definition

- lepton  $p_T > 25$  GeV & lepton  $|\eta| < 2.5$
- lepton  $p_T > 20$  GeV & lepton  $2.5 < |\eta| < 4.9$
- $66 < m_{ll} < 150$  GeV
- Unfolding to Born level lepton kinematics (dressed level available as a correction factor)

Cross sections unfolded using iterative Bayesian unfolding

$$\left. \frac{d^3\sigma}{dm_{\ell\ell} d|y_{\ell\ell}| d\cos\theta^*} \right|_{l,m,n} = \mathcal{M}_{ijk}^{lmn} \cdot \frac{N_{ijk}^{\text{data}} - N_{ijk}^{\text{bkg}}}{\mathcal{L}_{\text{int}}} \frac{1}{\Delta_{m_{\ell\ell}} \cdot 2\Delta_{|y_{\ell\ell}|} \cdot \Delta_{\cos\theta^*}}$$

$i, j, k$  = reco bin indices  
 $l, m, n$  = Born bin indices  
 $\mathcal{M}$  = inverted response matrix  
 $\Delta$  = bin widths in each variable



## Combination

Combine CC electron & muon channel measurements in averaging procedure  
 Minimise difference between two measurements  
 Taking correlated uncertainties into account

$i$  data points  
 $j$  systematic error sources

$$\chi_{tot}^2(\mathbf{m}, \mathbf{b}) = \sum_i \frac{[\mu^i - m^i(1 - \sum_j \gamma_j^i b_j)]^2}{\delta_{i,stat}^2 \mu^i m^i (1 - \sum_j \gamma_j^i b_j) + (\delta_{i,unc} m^i)^2} + \sum_j b_j^2$$

- bin-to-bin correlated error sources  $j$  including
- lepton trigger, ID, isolation efficiencies
  - lepton scale and resolution uncertainties
  - background contributions
  - etc....

- $\mu^i$  = measurement
- $m^i$  = averaged value
- $b_j$  = systematic error source strength  
 nuisance parameter left free in fit but constrained  
 no extra degrees of freedom due to additional constraint
- $\gamma_j^i$  = correlated sys uncertainty on point  $i$  from error source  $j$

Method allows cross-calibration of systematics between e and  $\mu$  channels  
 Improves statistical and systematic precision



Integrated single differential cross section

$$\frac{d\sigma}{dm_{\ell\ell}}$$

electron & muon CC channels combined

electron/muon combination gives  
 $\chi^2/\text{ndf} = 12.8/7$

Prediction from Powheg with CT10 PDFs

Partial NNLO (QCD) + NLO (EW) k-factors included:

→ 1-dimensional in  $m_{\ell\ell}$

Calculated with FEWZ in  $G_\mu$  EW scheme

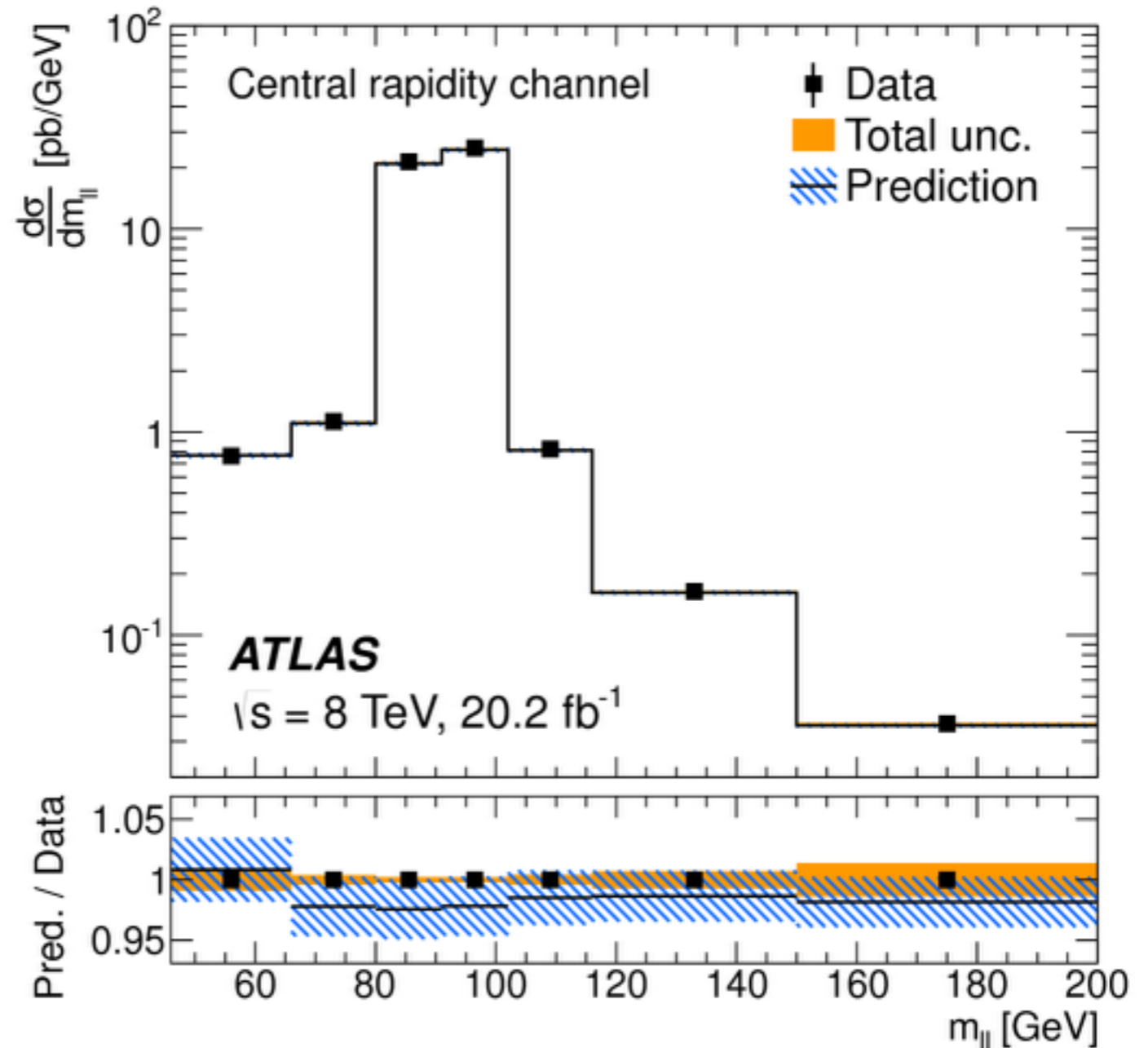
→ k-factor  $\sim 1.03$

Powheg has known mismodelling of  $A_0$  angular polarisation coefficient (goes negative)

→ reweighted vs  $p_{T,Z}$  and  $y_{\ell\ell}$

Computed with DYNNLO

$$\frac{d^2\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|} \quad \text{2d cross sections in back-up}$$



orange band: data uncertainty (excl. lumi  $\pm 1.9\%$ )

blue band: MC stat + PDF uncertainty

(CT10 68% eigenvectors)

# Triple-differential $Z/\gamma^*$ Cross Sections $\sqrt{s} = 8 \text{ TeV}$



$$\frac{d^3\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|d\cos\theta^*}$$

Central rapidity electron & muon combined result

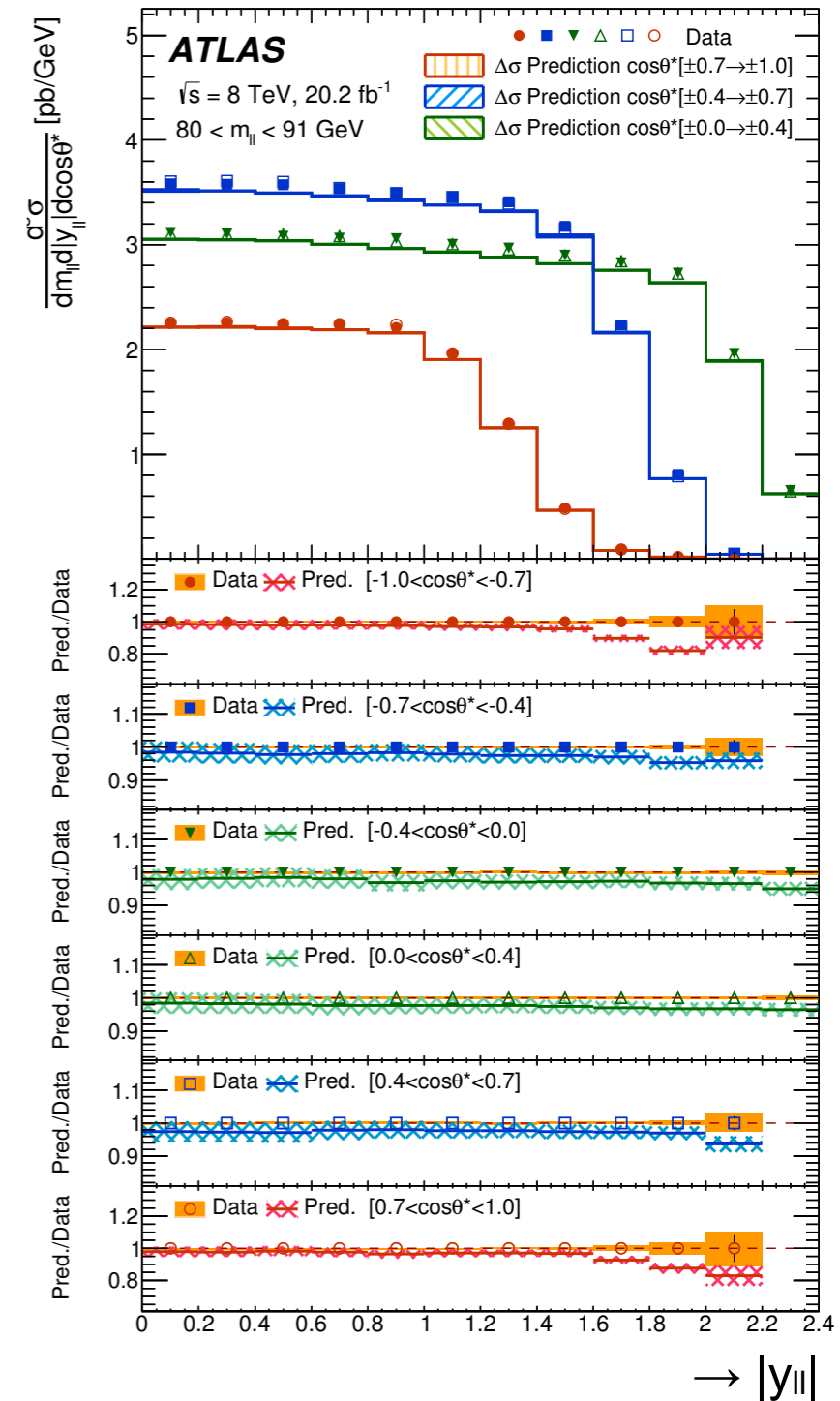
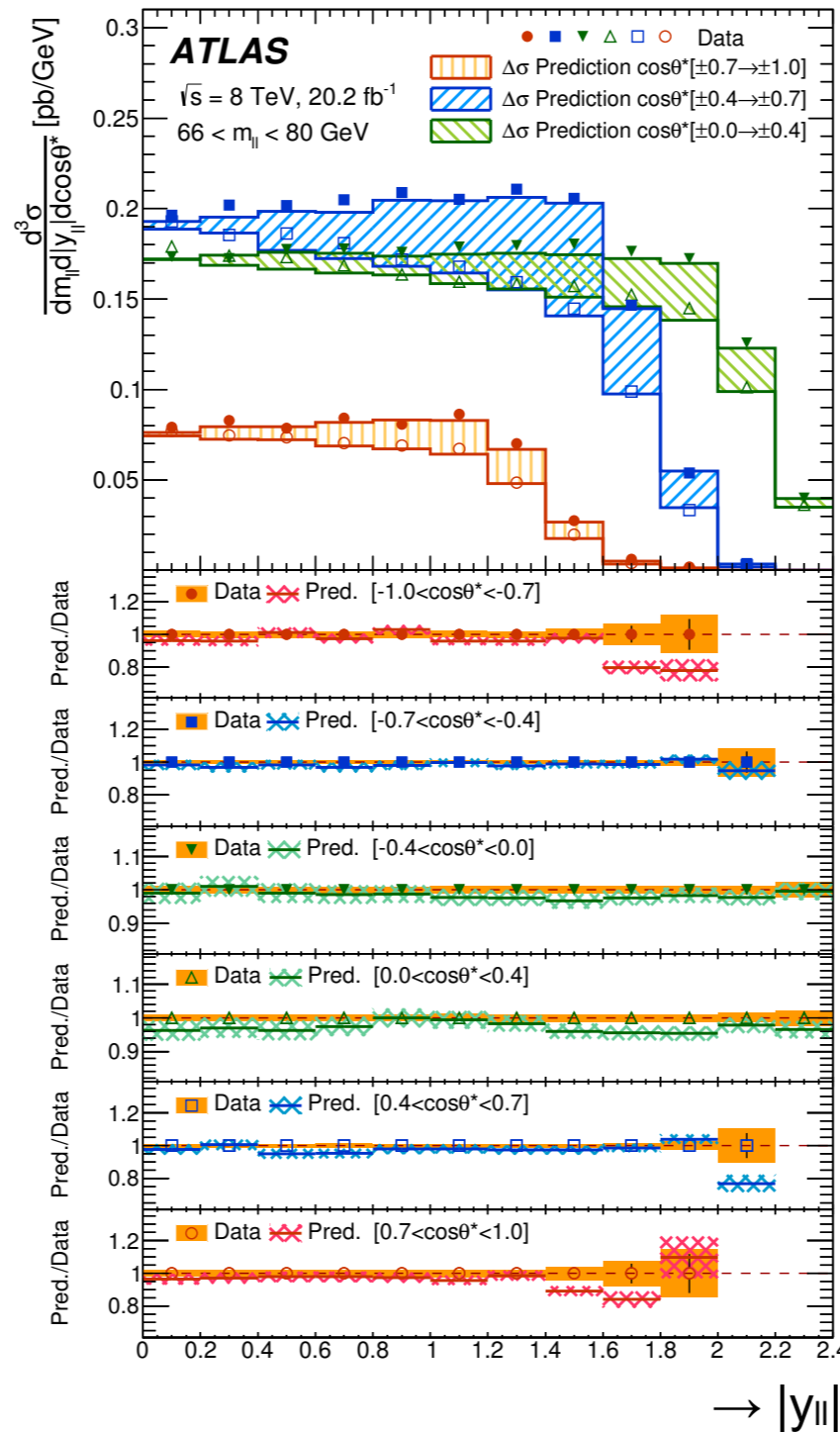
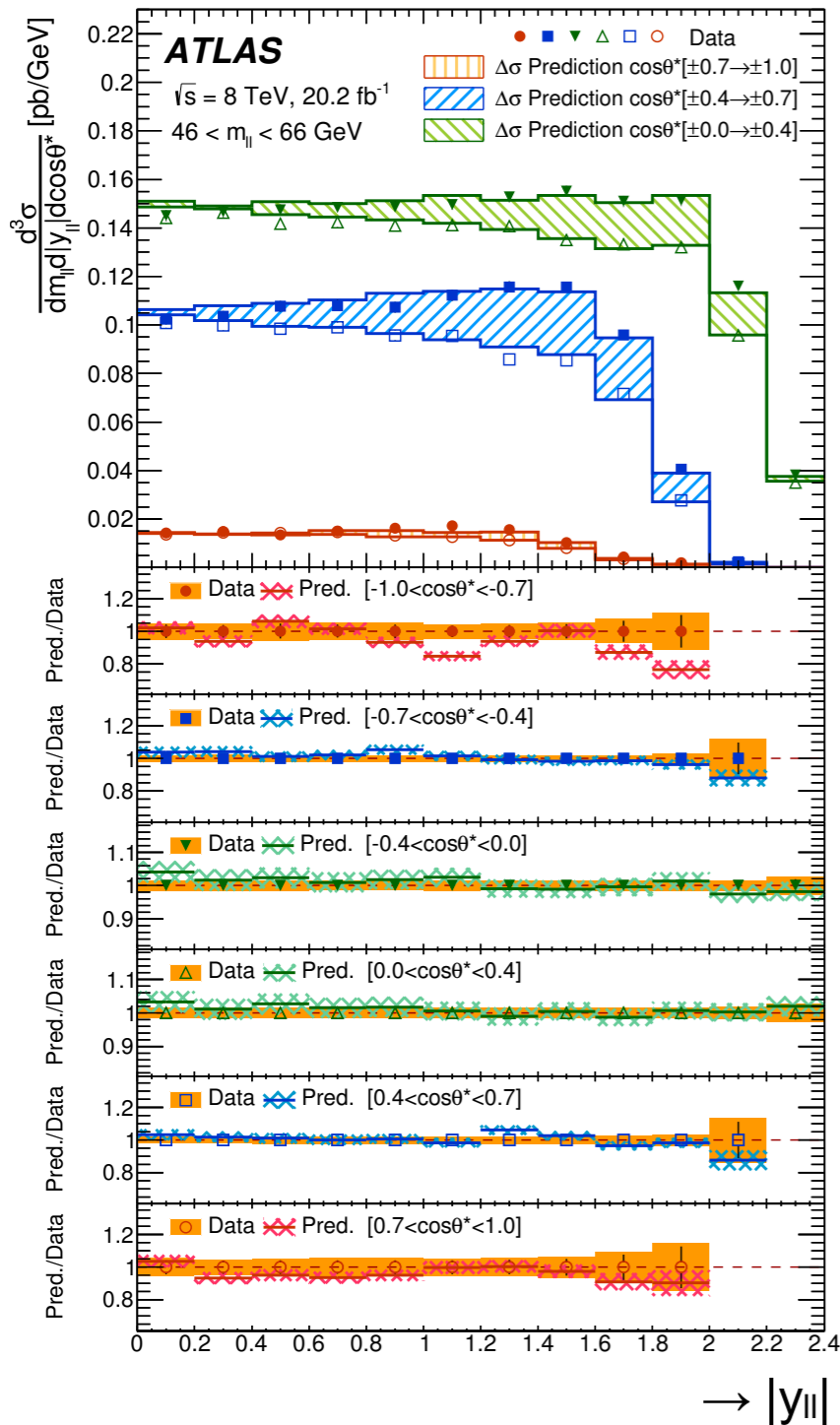
Large forward-backward asymmetry at low mass, decreasing to ~zero at  $m_{\parallel} \sim m_Z$

Upper plots: shaded regions highlight equal  $|\cos\theta^*|$

46 < m < 66 GeV

66 < m < 80 GeV

80 < m < 91 GeV

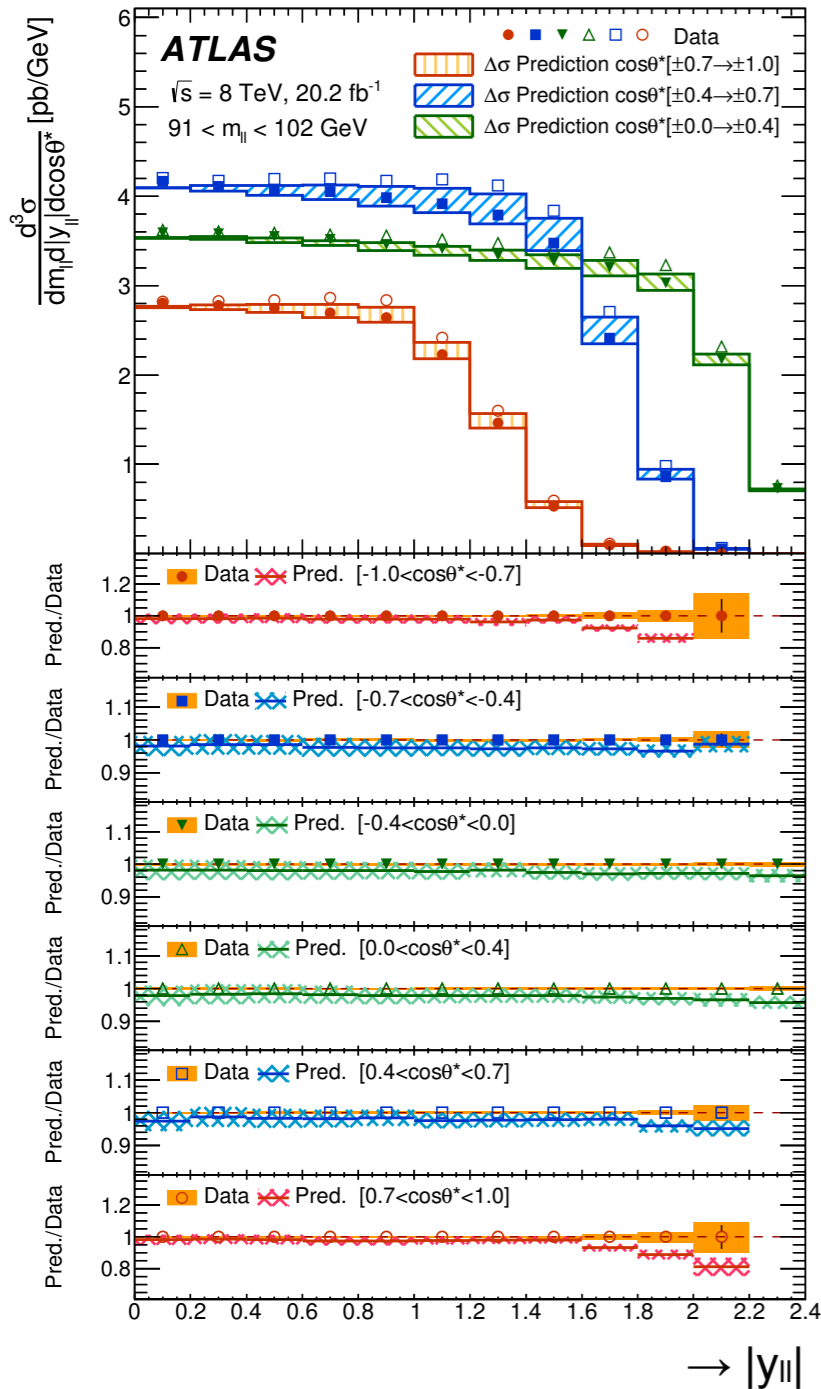


# Triple-differential Z/ $\gamma^*$ Cross Sections $\sqrt{s} = 8$ TeV

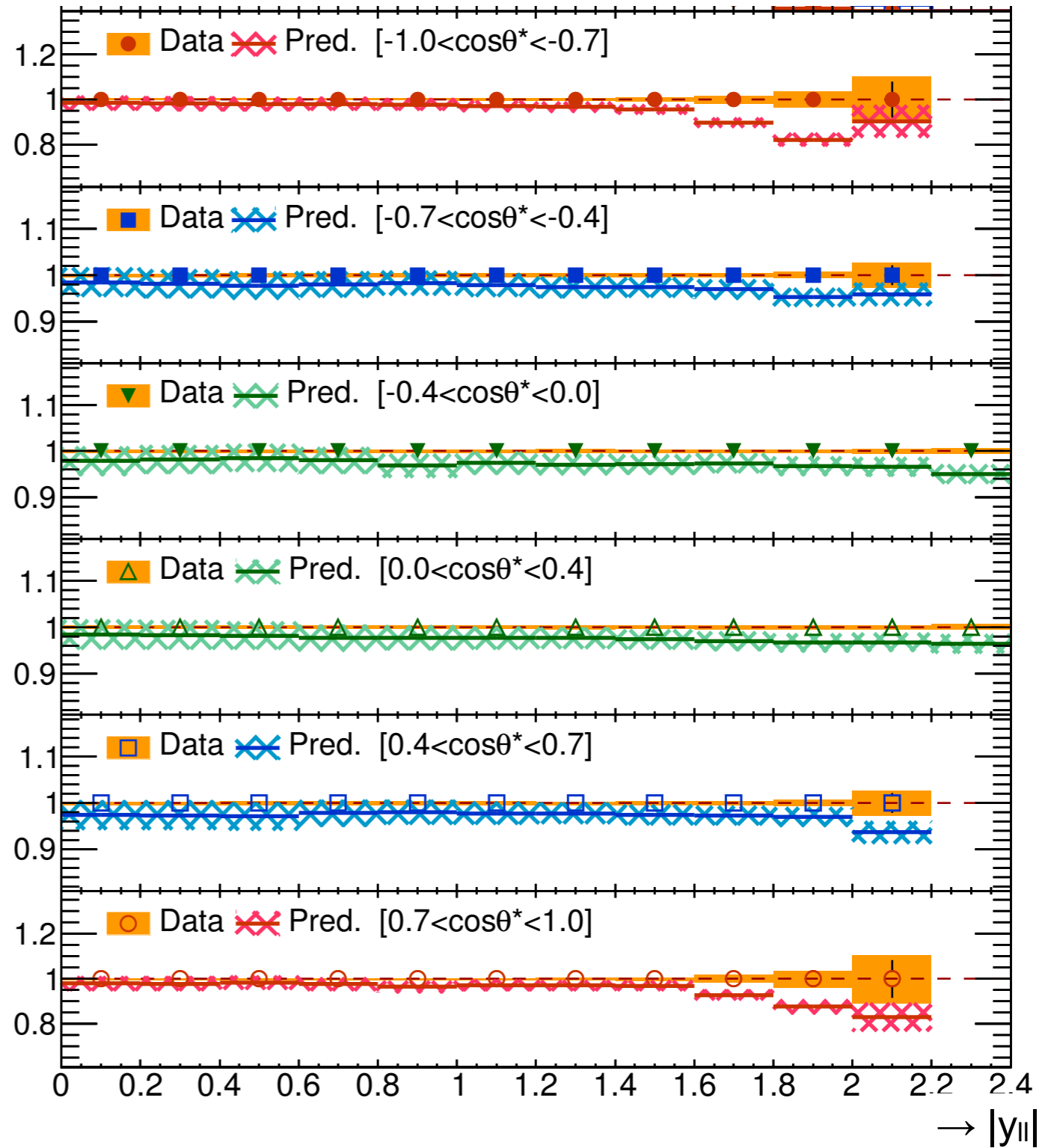


electron / muon combination gives  $\chi^2/\text{ndf} = 489.4 / 451$

91 < m < 102 GeV



Pred./Data



Data precision reaches  $\sim 0.5\%$  for  $m_{||} \sim m_Z$

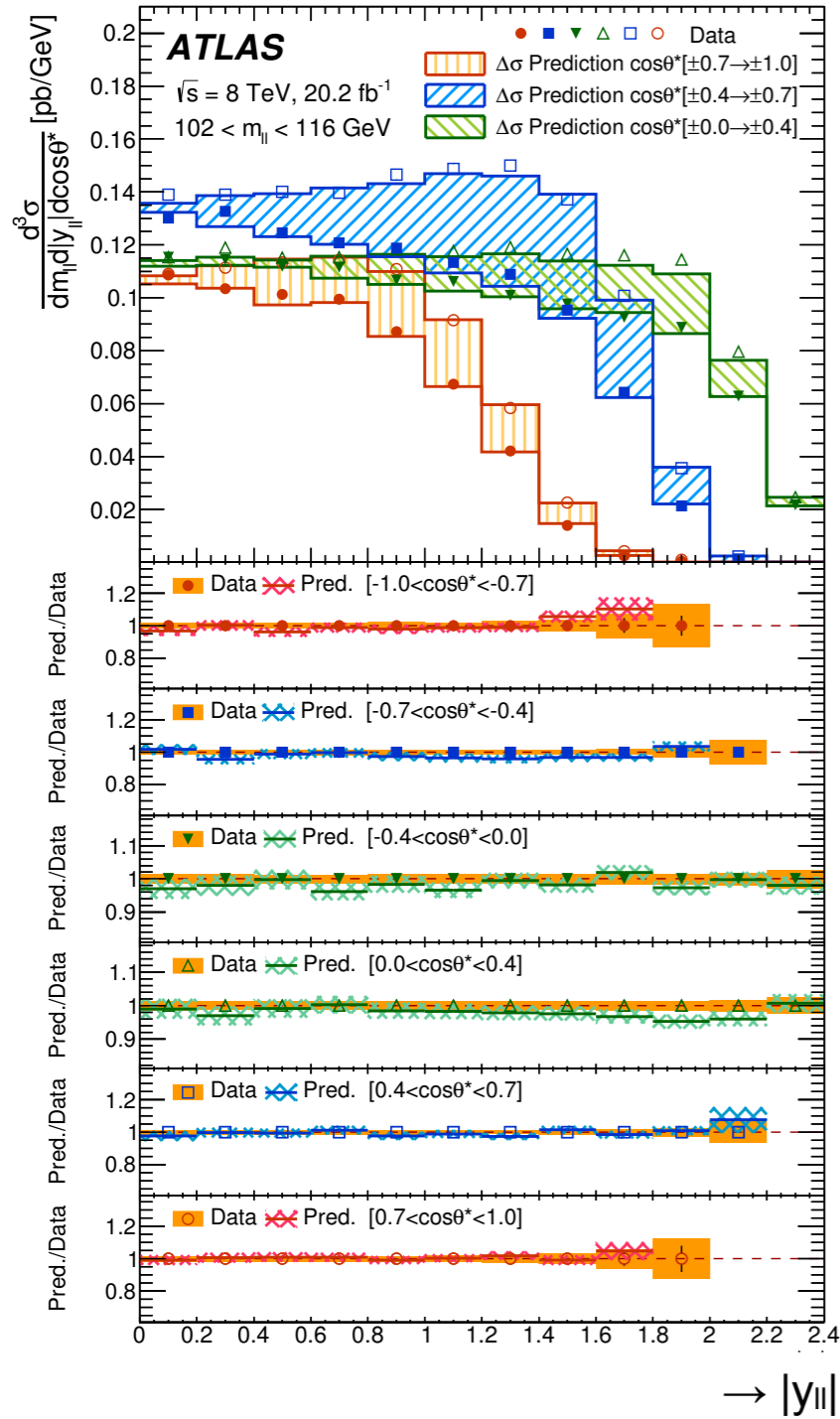
Good agreement with Powheg based predictions incl. NNLO/NLO k-factor (and  $A_0$  polarisation correction)

Interesting features at high  $|y|$  + large  $|\cos \theta^*| \approx 1$  ....

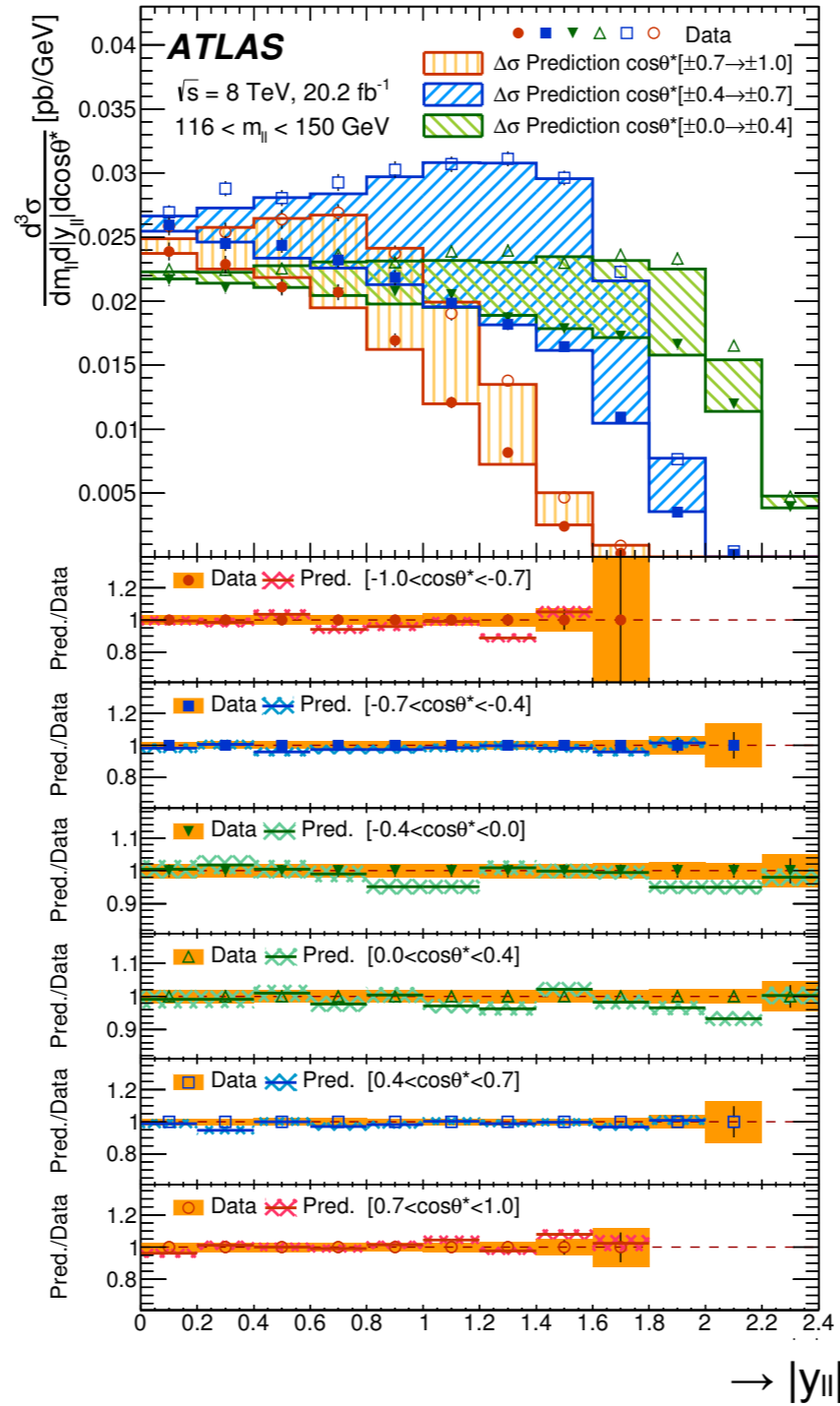
# Triple-differential $Z/\gamma^*$ Cross Sections $\sqrt{s} = 8 \text{ TeV}$



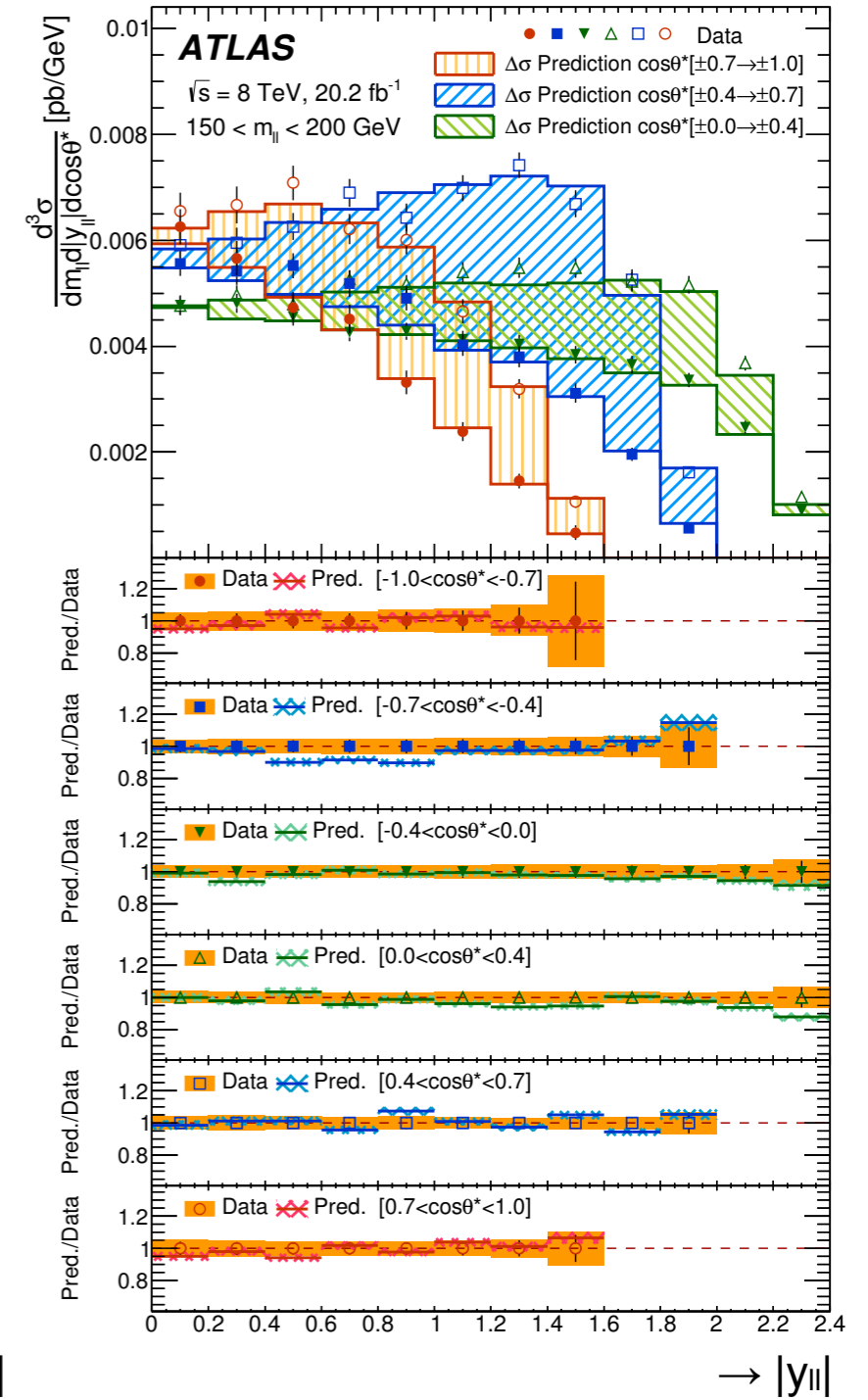
## 102 < m < 116 GeV



## 116 < m < 150 GeV



## 150 < m < 200 GeV

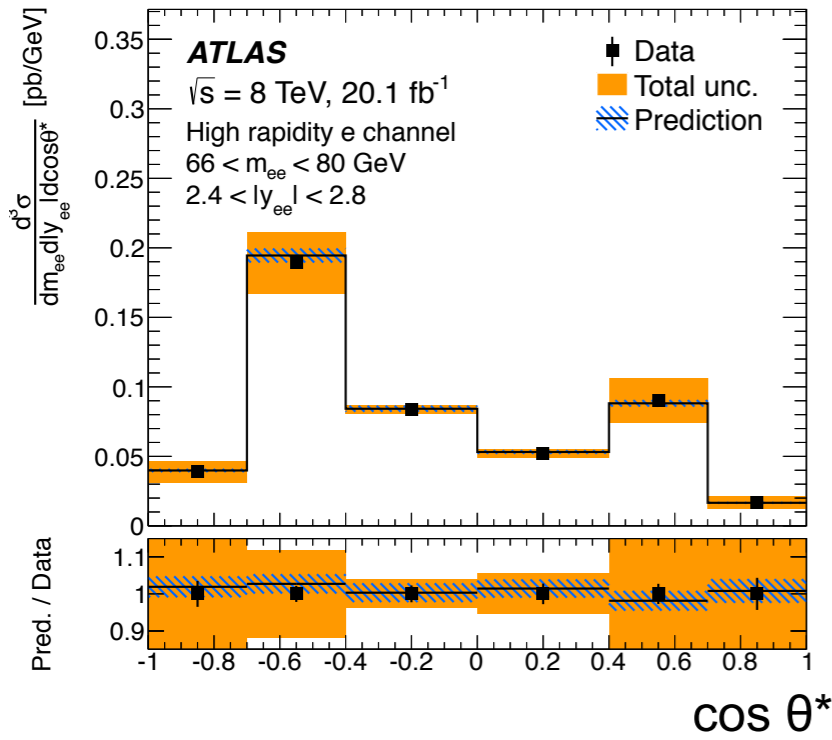




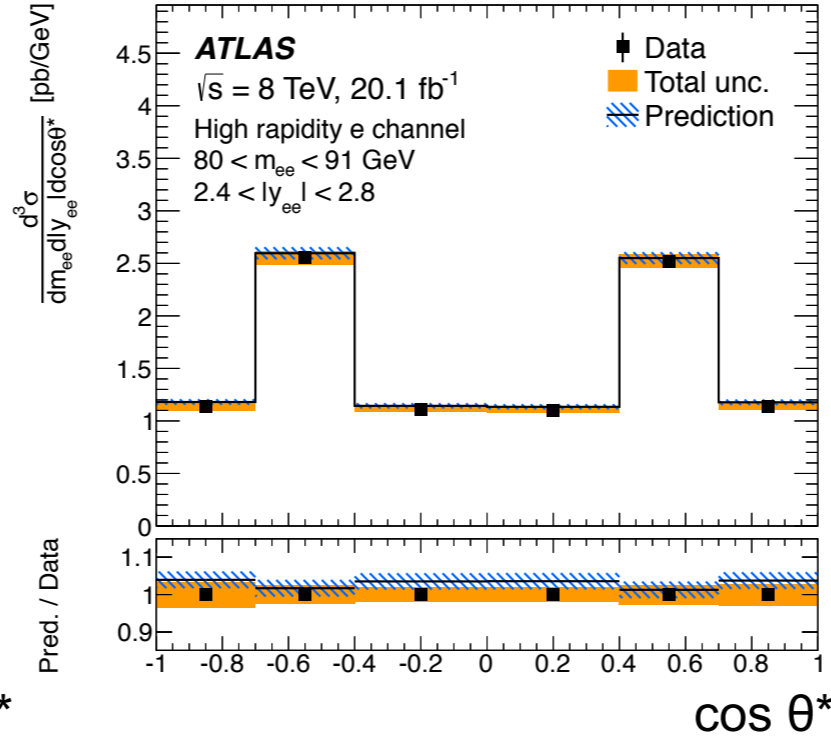
# Triple-differential $Z/\gamma^*$ Cross Sections $\sqrt{s} = 8 \text{ TeV}$



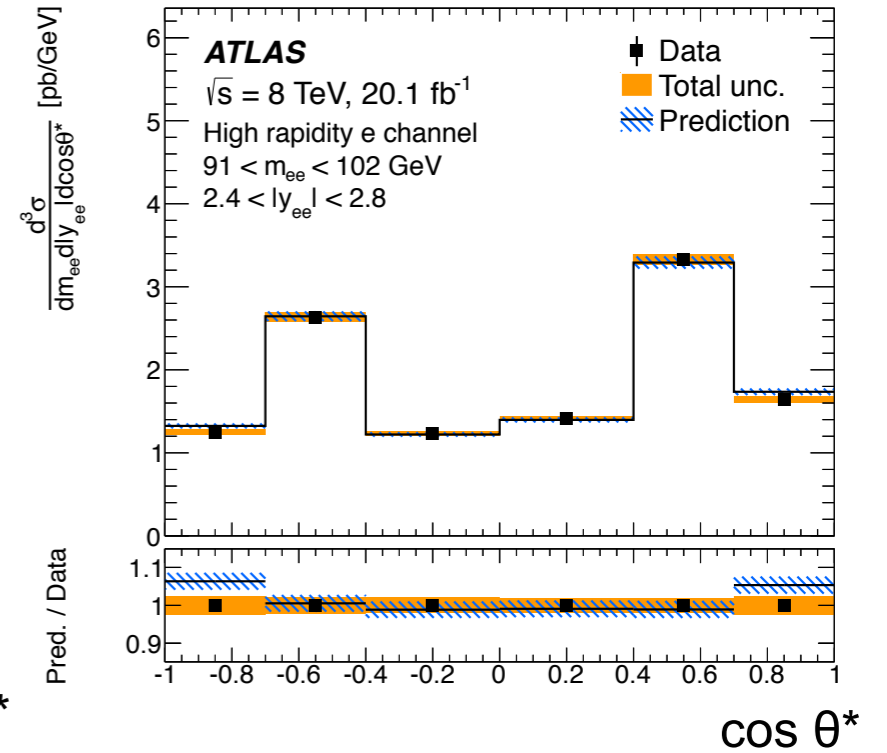
### 66 < m < 80 GeV



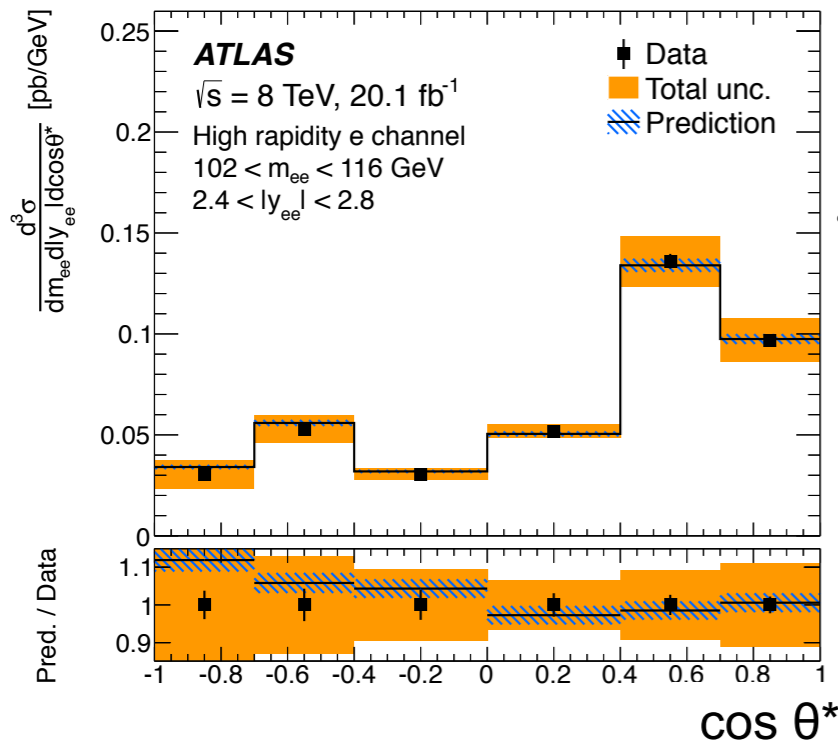
### 80 < m < 91 GeV



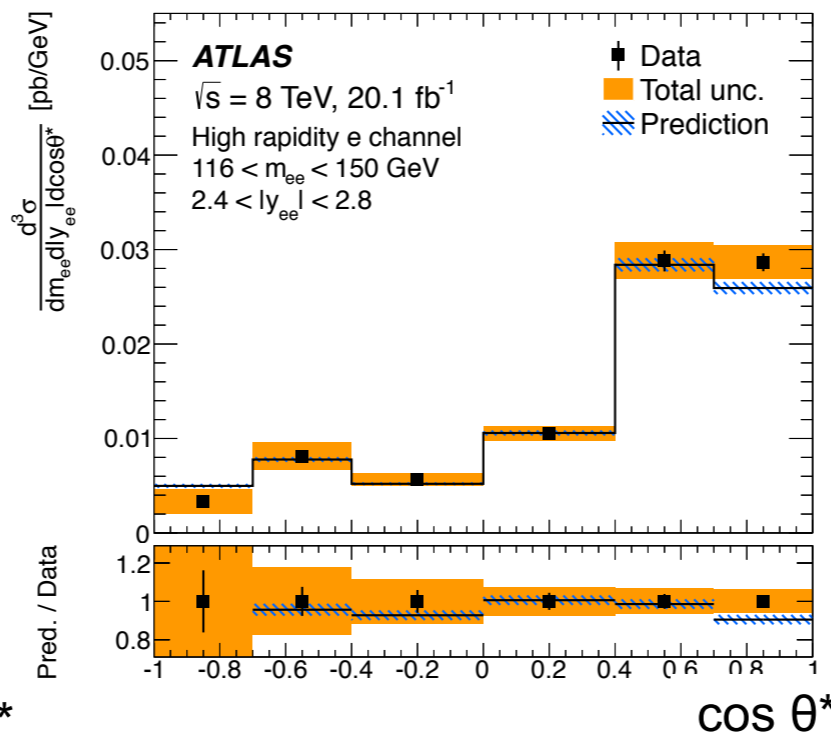
### 91 < m < 102 GeV



### 102 < m < 116 GeV



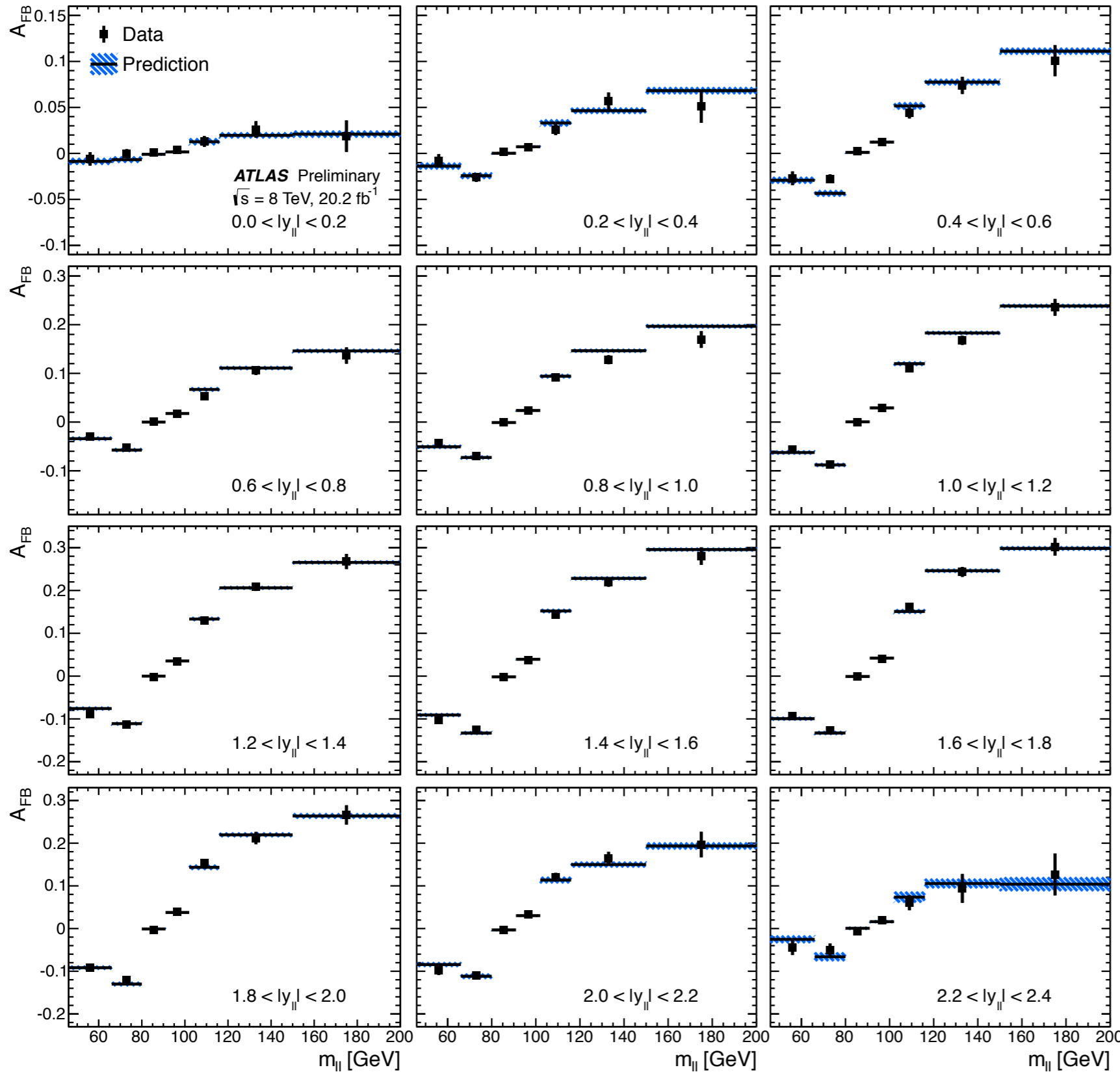
### 116 < m < 150 GeV



$$\frac{d^3 \sigma}{dm_{ee} d|y_{ee}| d \cos \theta^*}$$

**High rapidity channel**  
 Showing selected bins  
 $2.4 < |y| < 2.8$

High  $y$  region has greatest sensitivity to  $\sin^2 \theta_W$  and PDFs  
 High  $y$  analysis shows much larger asymmetry



## Central rapidity channel

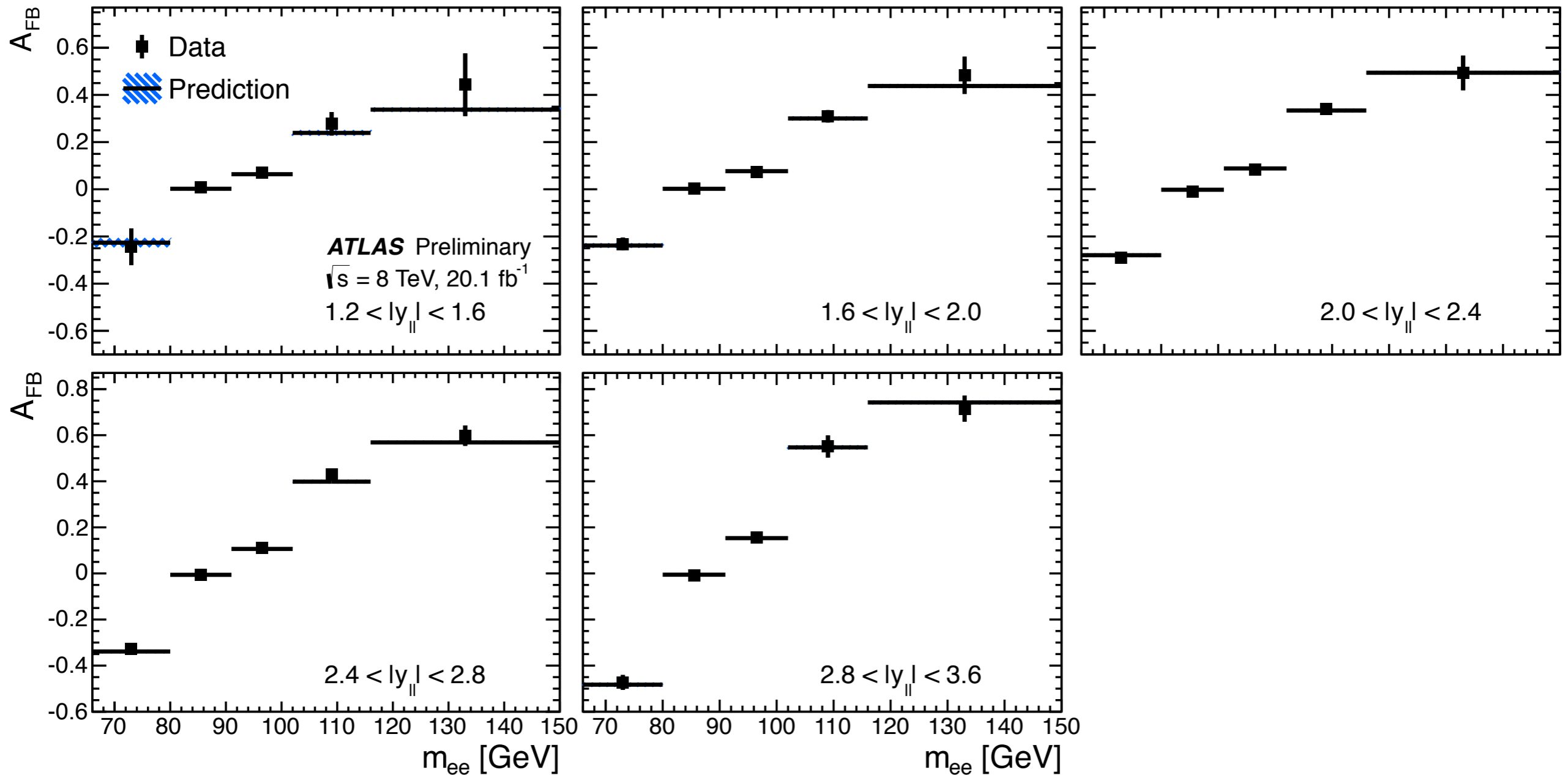
$$A_{\text{FB}} = \frac{d^3\sigma(\cos\theta^* > 0) - d^3\sigma(\cos\theta^* < 0)}{d^3\sigma(\cos\theta^* > 0) + d^3\sigma(\cos\theta^* < 0)}$$

Note:  $A_{\text{FB}}$  derived from unfolded cross section measurements

symmetric uncertainties cancel in  $A_{\text{FB}}$   
Scale, resolution, backgrounds

Asymmetry increases with  $|y|$   
Due to better determination of initial quark direction (less dilution)

(high  $|y|$  access higher  $x$  valence PDF)



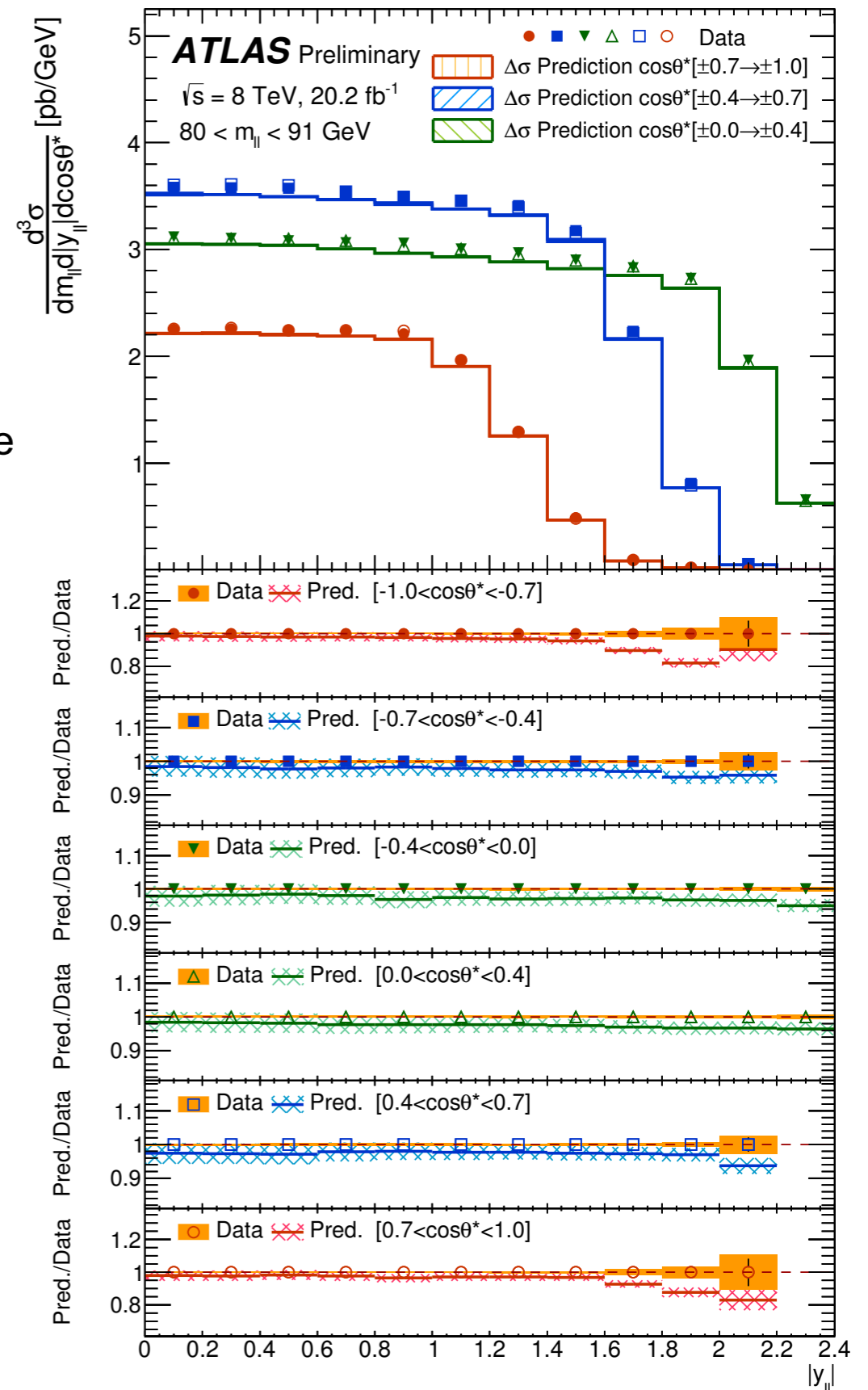
## High rapidity channel

For  $A_{FB}$  measurements uncorrelated sources dominate:  
 data stats are factor 2 larger than MC stat / multijet unc / bg MC stats  
 correlated sources ~ factor 10 smaller

- New  $d^3\sigma$  measurement of DY cross section at  $\sqrt{s} = 8$  TeV available
- on-shell analysis covers phase space  $46 < m < 200$  GeV
- Precision of 0.5% attained at  $m = m_z$
- Data compatible with NNLO pQCD  $\otimes$  NLO EW
- Data available on HepData with full systematic breakdown

**Now extract  $\sin^2\theta_{\text{eff}}$  using this data**

Method of using unfolded  $d^3\sigma$  cross sections never used before



Typically experiments measure  $A_{\text{FB}}$

- unfold detector effects / dilution → fit for  $\sin^2\theta_{\text{eff}}$
- or, perform detector level template fits to  $A_{\text{FB}}$
- estimate PDF uncertainties on extraction

## D0 + CDF combination 2017

$$\begin{aligned} \sin^2 \theta_{\text{eff}}^{\text{lept}} &= 0.23148 \pm 0.00027 \text{ (stat.)} \\ &\pm 0.00005 \text{ (syst.)} \\ &\pm 0.00018 \text{ (PDF)} \end{aligned}$$

At LHC / Tevatron largest uncertainty ~ PDFs  
worse at LHC due to pp collisions  
worse at larger  $\sqrt{s}$  due to lower x (more dilution)

## ATLAS 7 TeV

$$\sin^2 \theta_{\text{eff}}^{\text{lept}} = 0.2308 \pm 0.0005(\text{stat.}) \pm 0.0006(\text{syst.}) \pm 0.0009(\text{PDF}) = 0.2308 \pm 0.0012(\text{tot.})$$

## CMS 8 TeV

$$\begin{aligned} \sin^2 \theta_{\text{eff}}^{\text{lept}} &= 0.23101 \pm 0.00036(\text{stat}) \pm 0.00018(\text{syst}) \pm 0.00016(\text{theory}) \pm 0.00030(\text{pdf}) \\ \sin^2 \theta_{\text{eff}}^{\text{lept}} &= 0.23101 \pm 0.00052. \end{aligned}$$

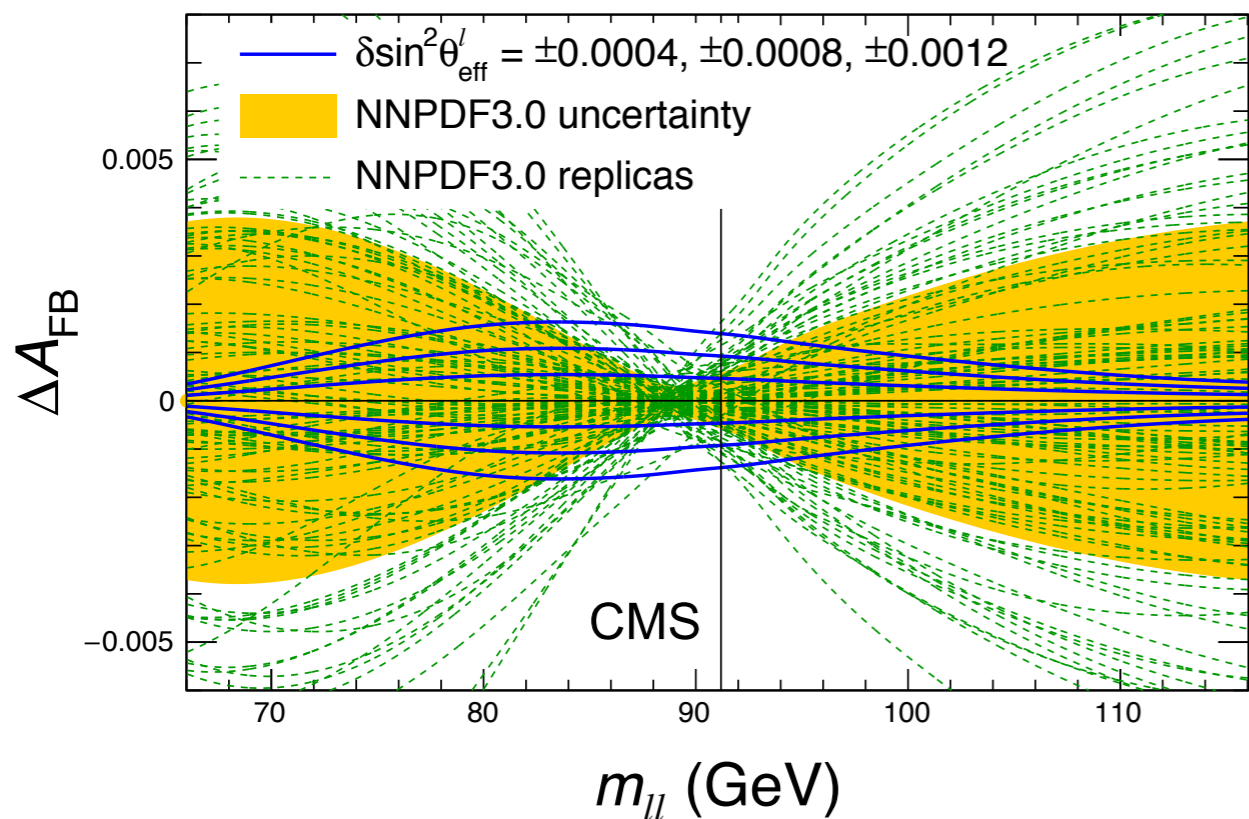
## LHCb 7 & 8 TeV

$$\sin^2 \theta_{\text{W}}^{\text{eff}} = 0.23142 \pm 0.00073(\text{stat}) \pm 0.00052(\text{sys}) \pm 0.00056(\text{theo})$$

dominated by PDF



Variation of  $A_{\text{FB}}$  from PDF replicas and  $\sin^2\theta_w$



$\sin^2\theta_w$  variations correlated across  $m$  spectrum  
 PDF variations anti-correlated about  $m=91$

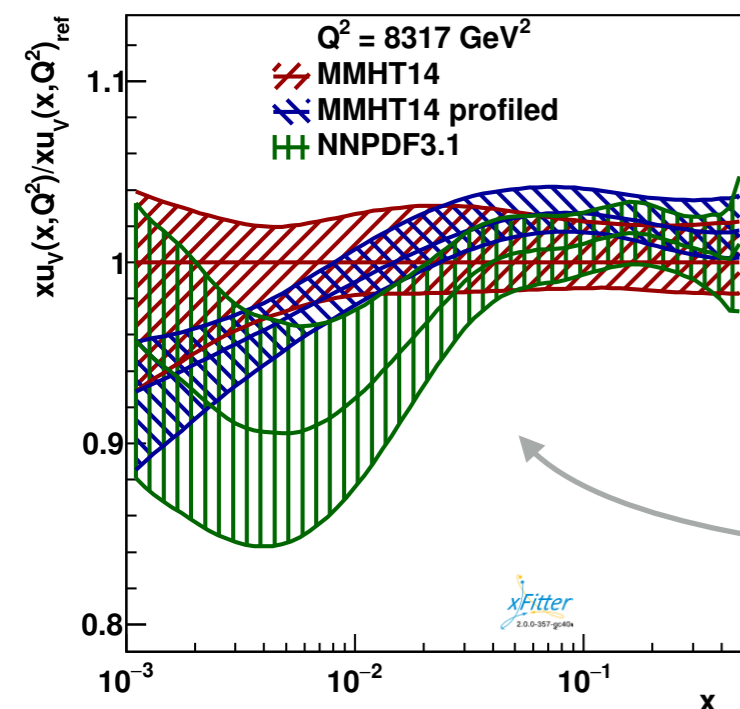
These correlations can be exploited  
 Use data to constrain PDFs  $\rightarrow$  reduce uncertainty

For NNPDF incompatible replicas rejected by data

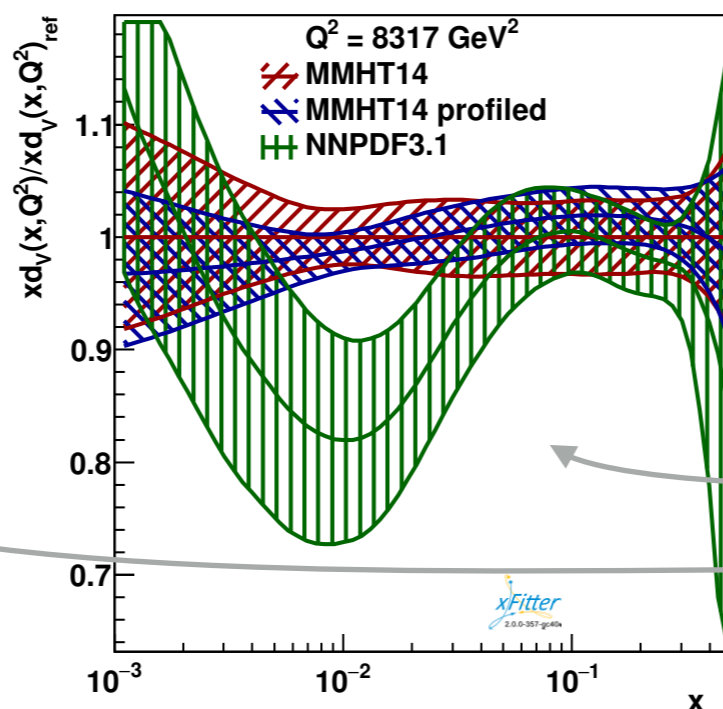
Other PDF sets: uncertainties given as eigenvector variations  
 Introduce nuisance parameters for each PDF eigenvector  
 Fit data + PDF nuisance parameters to constrain PDFs

Approximation to performing full PDF fit to data

$u_v$  compared to MMHT reference



$d_v$  compared to MMHT reference



Example of profiling using  $d^3\sigma$  pseudo-data  
 Pseudo-data produced with NNPDF set  
 Predictions generated using MMHT  
 Pseudo-data are profiled using predictions  
 Profiled PDFs move towards MMHT

Profiling works for  $u_v$  but fails for  $d_v$  where  
 PDF set has insufficient flexibility  
 $\rightarrow$  use several PDF sets

ATLAS uses 2 methods (same data set / similar selections):

Triple Differential cross section analysis  
Fit to unfolded  $d^3\sigma$  cross sections  
differential in  $m, |y|, \cos\theta^*$

Ai - Angular coefficient analysis  
methodology used here [arXiv:1606.00689](https://arxiv.org/abs/1606.00689)

Use QCD predictions from NNLOJet  
Use higher order EW corrections  
event weights to vary  $\sin^2\theta_{\text{eff}}$   
Perform PDF profiled fit to  $d^3\sigma$  data

common NLO EW  
corrections

Use QCD predictions from DYTurbo  
Use higher order EW corrections  
event weights to vary  $\sin^2\theta_{\text{eff}}$   
Perform PDF profiled fit to data

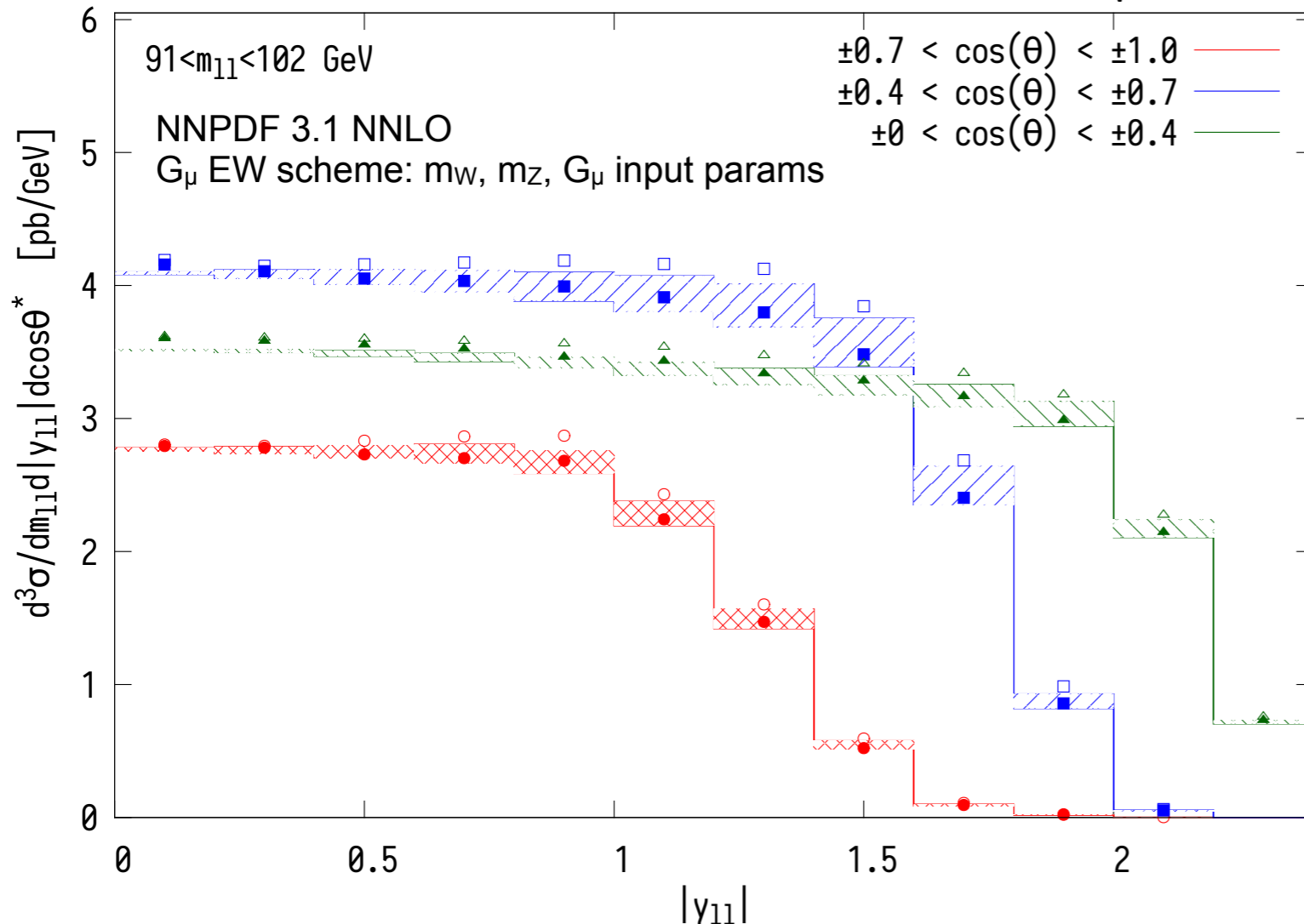
Both analyses were blinded to value of  $\sin^2\theta_{\text{eff}}$

Compare sensitivities and results

NNLO  $|y_{11}|$  Triple Differential [Central Region]

**NNLOJET**

$\sqrt{s} = 8 \text{ TeV}$



Parton level event generator at NN(N)LO  
 QCD using antenna subtraction

Processes available:

- $pp \rightarrow H, H+J,$
- $Z, Z+J,$
- $W^\pm, W^\pm J,$
- $VH,$
- dijets
- $ep \rightarrow 1,2J$
- $e^+e^- \rightarrow 3J \dots$

Comparisons to NNLOjet - collaboration with IPPP (Nigel Glover & Duncan Walker)

Provide **fiducial** NNLO QCD predictions for varying  $\sin^2\theta_W$

Full set of NNLO predictions = ~3-4 days grid time

Applfast interface under development (for PDF uncertainties)

QCD scale uncertainties  $\mu_R$  &  $\mu_F$  ~0.5% ...

...but larger dependence observed in some kinematic regions...

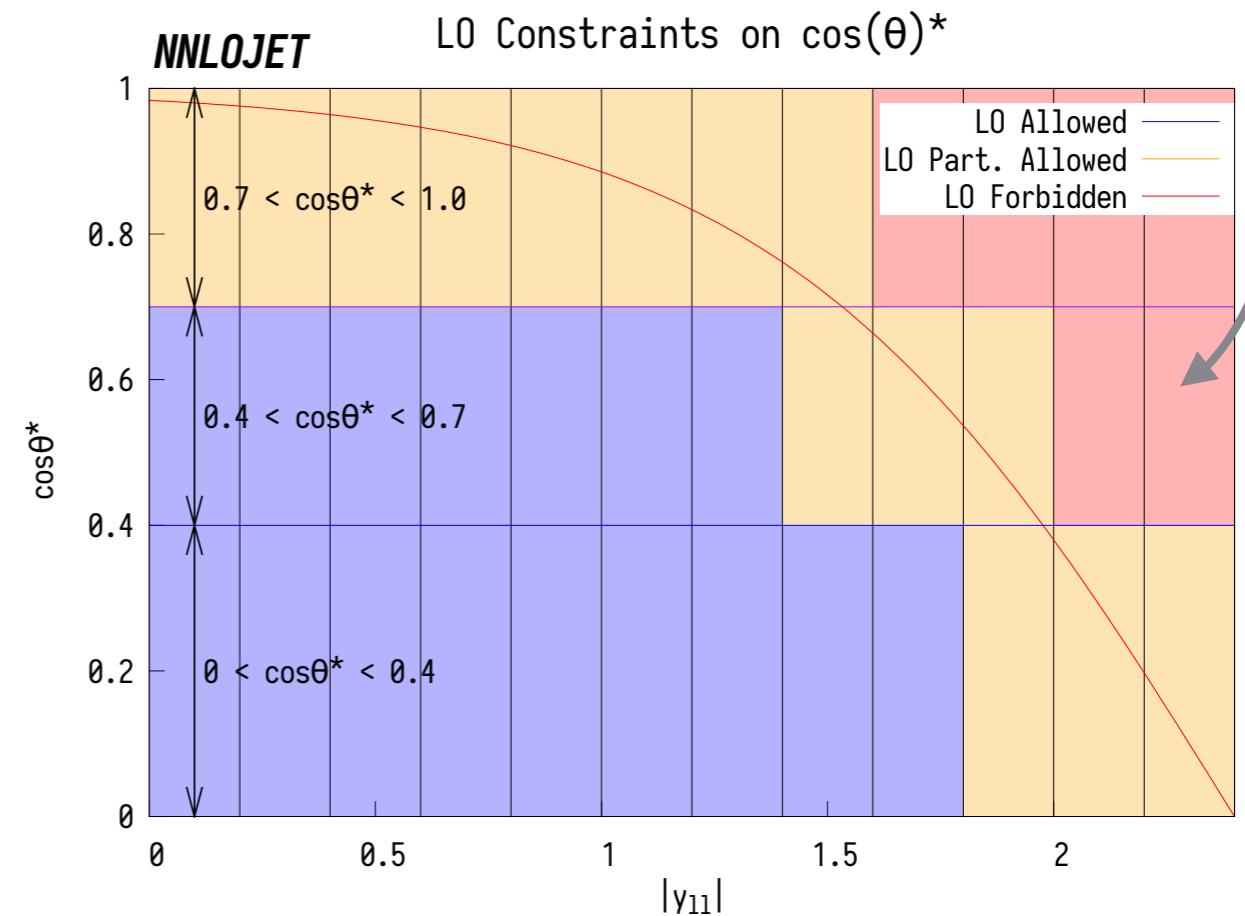
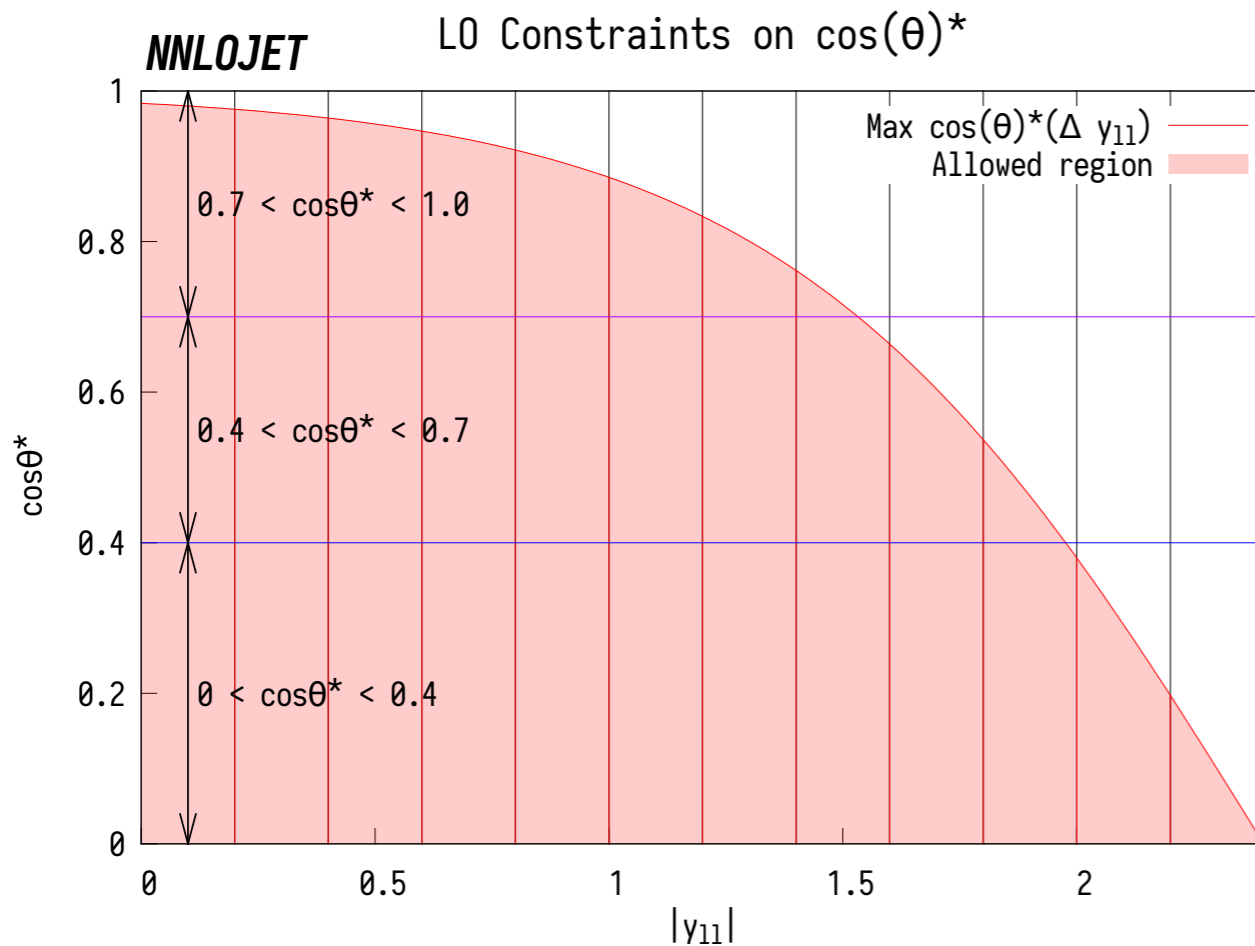


Using LO kinematics, we can write  $\cos \theta^*$  a function of the difference in rapidities of the leptons:

$$\cos \theta^* = \frac{\sinh(\Delta y_{ll})}{1 + \cosh(\Delta y_{ll})} \rightarrow \cos \theta^* \leq \frac{\sinh(2(y_l^{\text{max}} - |y_{ll}|))}{1 + \cosh(2(y_l^{\text{max}} - |y_{ll}|))}$$

Constraints on  $\Delta y_{ll}$  from the cuts give constraints on  $\cos \theta^*$ .  
Only NLO in these bins at  $\mathcal{O}(\alpha_S^2)$   $\rightarrow$  use NNLO ZJ calculation?

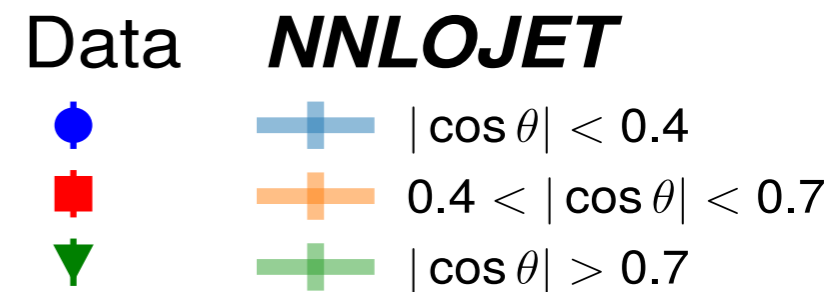
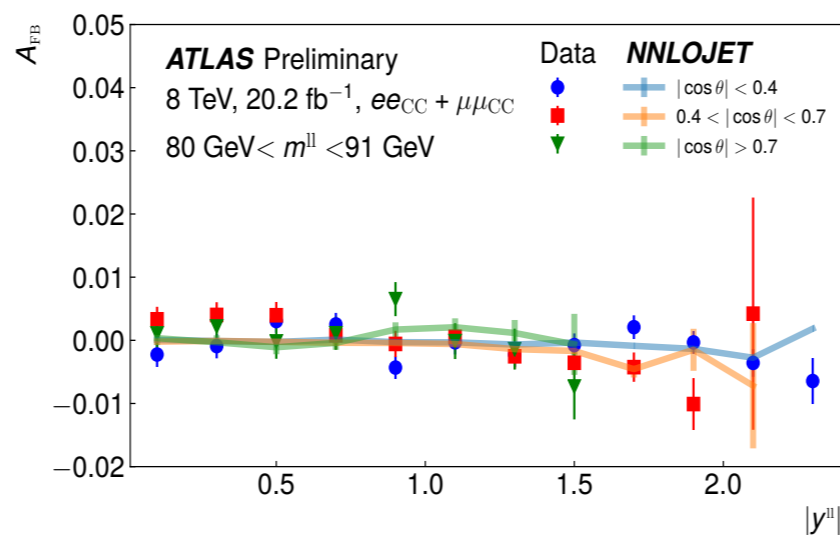
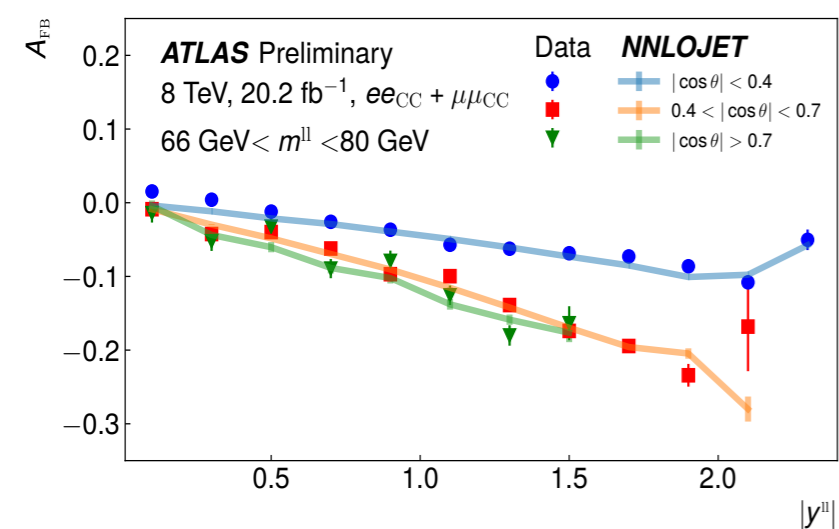
Region corresponds to Z recoiling against jet



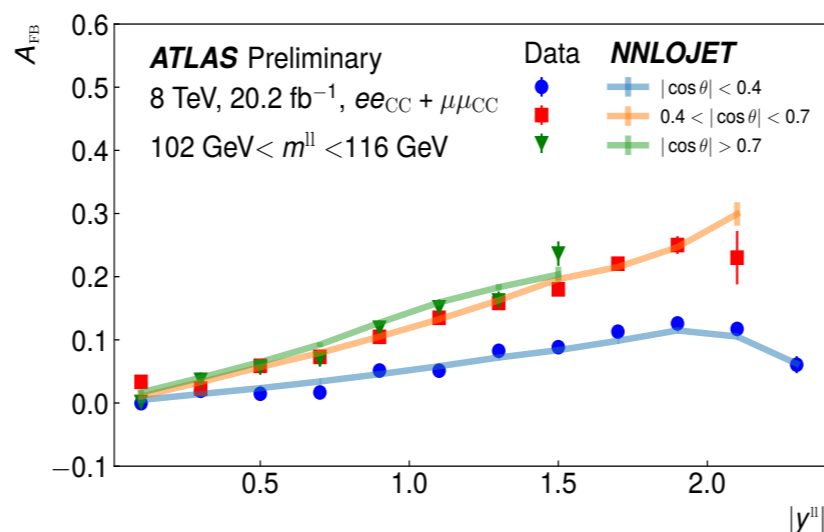
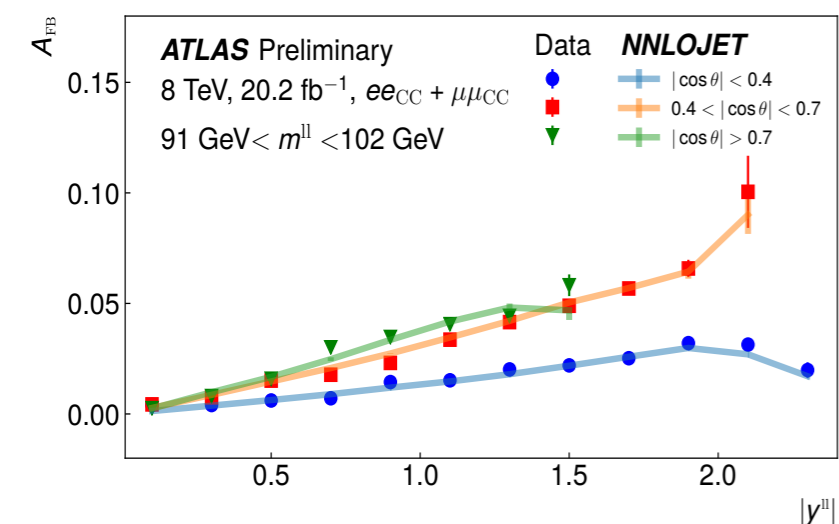
Observe large theory stat & scale errors in “forbidden region” predictions

$\Rightarrow$  Use differential  $A_{\text{FB}}$  in “forbidden region”  
Scale uncertainty cancels in  $A_{\text{FB}}$   
All data points can be used in fit

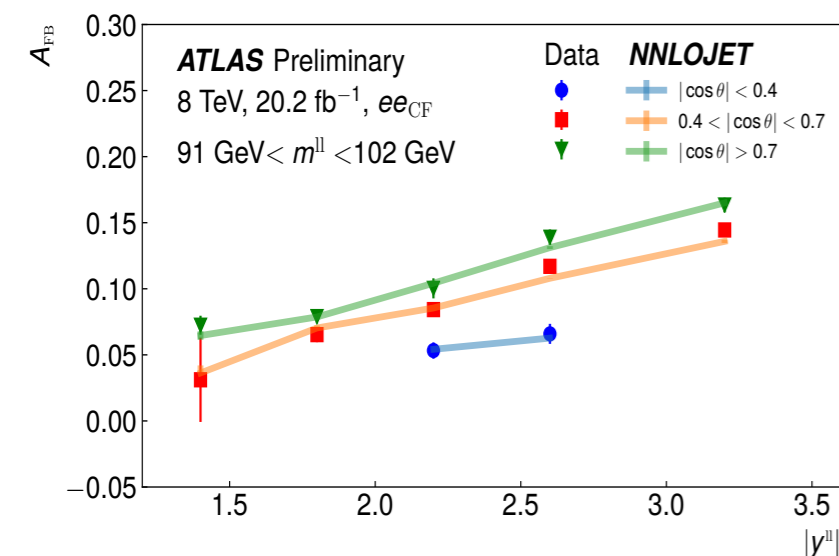
scale choice  $\mu^2 = m^2 + p_{\text{T},ll}^2$   
Equivalent to  $m^2$  at LO  
Apt choice for recoil jet topology



Triple differential  $A_{\text{FB}}(m, |y|, \cos\theta^*)$



$|A_{\text{FB}}|$  increases with  $y$   
 $A_{\text{FB}}$  negative  $m < m_z$   
Smallest for  $\cos\theta^* \sim 0$



Use predictions of differential  $A_{\text{FB}}(m, |y|, \cos\theta^*)$  from NNLOjet  
i.e. defined in slices of equal  $|\cos\theta^*|$

Apply identical event reweighting to vary  $\sin^2\theta_{\text{eff}}$   
for NLO EW effects in Improved Born Approximation (IBA)

ATLAS uses 2 methods (same data set / similar selections):

- Perform fit to unfolded  $A_{FB}$  from  $d^3\sigma$  cross sections differential in  $m, |y|, \cos\theta^*$
- $A_i$  - Angular coefficient analysis (methodology used here [arXiv:1606.00689](https://arxiv.org/abs/1606.00689))

## Angular Coefficients

Full 5d cross section decomposed into  
9 polynomials & 9 coefficients  $A_i(m, y, p_T)$   
Description is complete to all orders in QCD  
- only in full phase space of decay leptons

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ \begin{aligned} & (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \\ & + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \\ & + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \end{aligned} \right\}$$

factorised  
production dynamics  
from decay kinematics

$$A_{FB} = \frac{8}{3} A_4$$

in full phase space

$A_3$  and  $A_4$  related to  $\sin^2\theta_{\text{eff}}$   
( $A_3$  contributes for  
 $p_{T,Z} > 100$  GeV)

Using  $y$  and  $m$  binned data allows PDFs to be profiled Bin data in  $m, |y|$

CC (x2 channels):

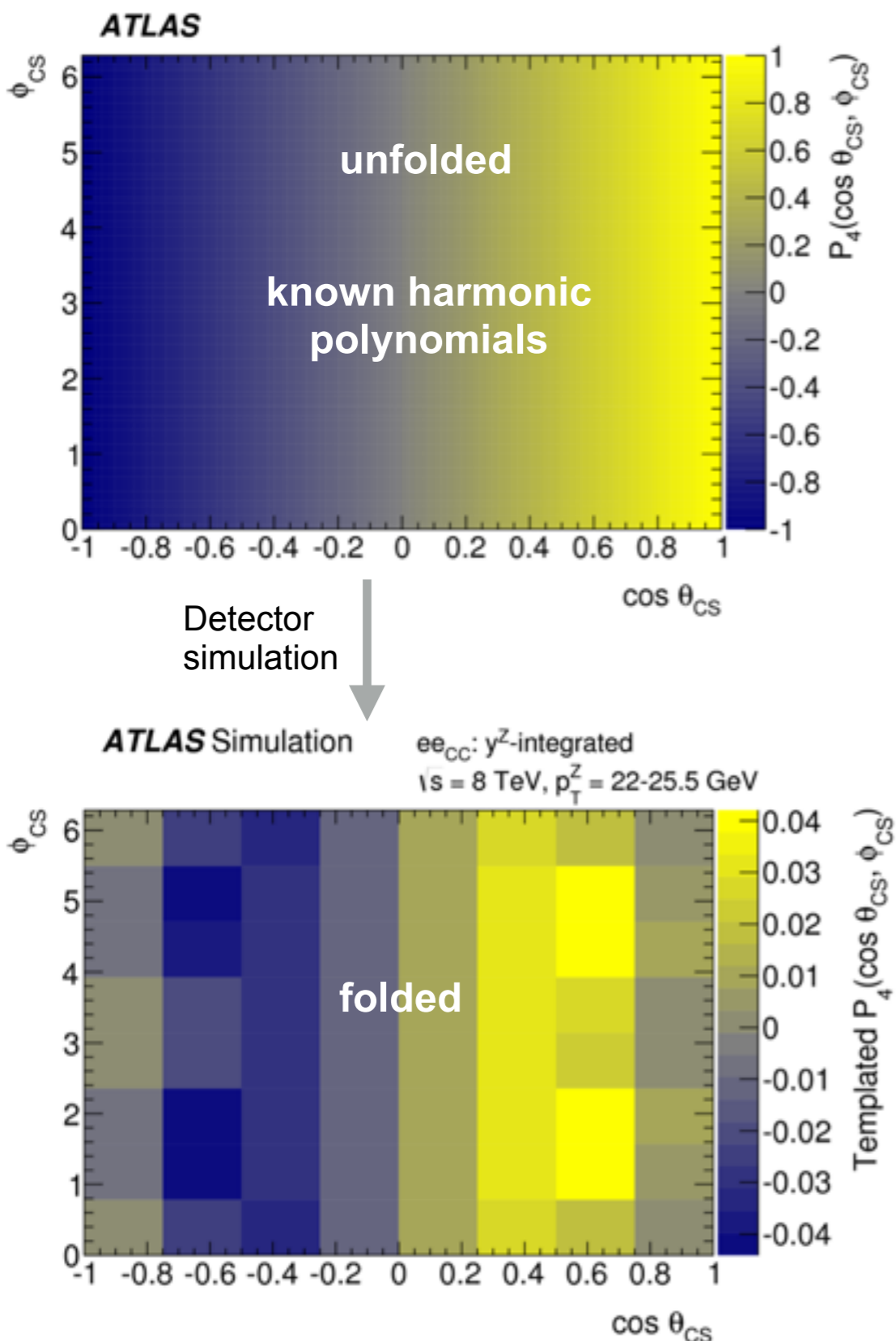
$m$ :- {70, 80, 100, 125} GeV  
 $|y|$ :- {0.0, 0.8, 1.6, 2.5}

CF

$m$ :- {80, 100} GeV  
 $|y|$ :- {1.6, 2.5, 3.6}

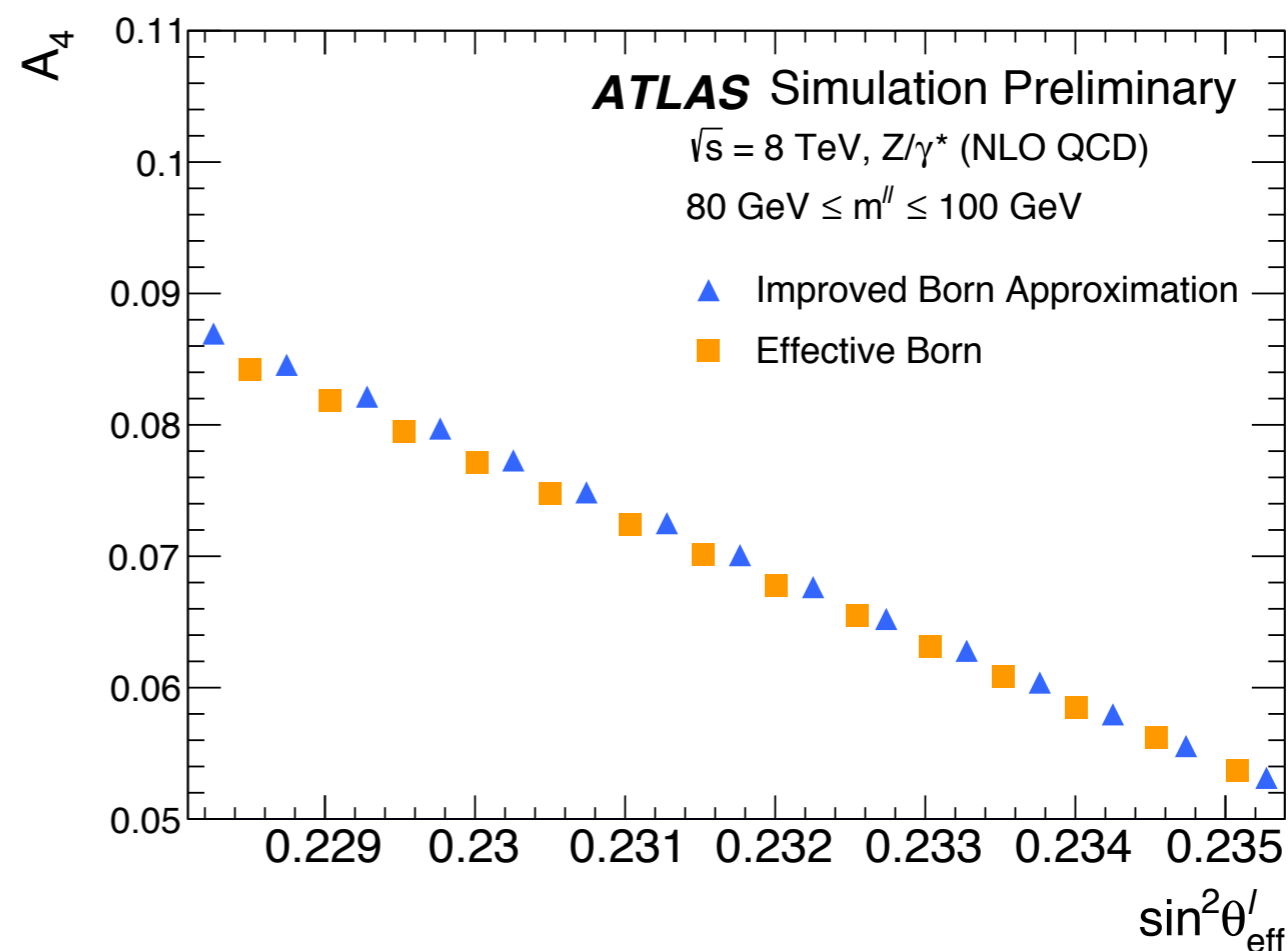
Analysis method uses folded MC templates in full phase-space

ATLAS-CONF-2018-037/

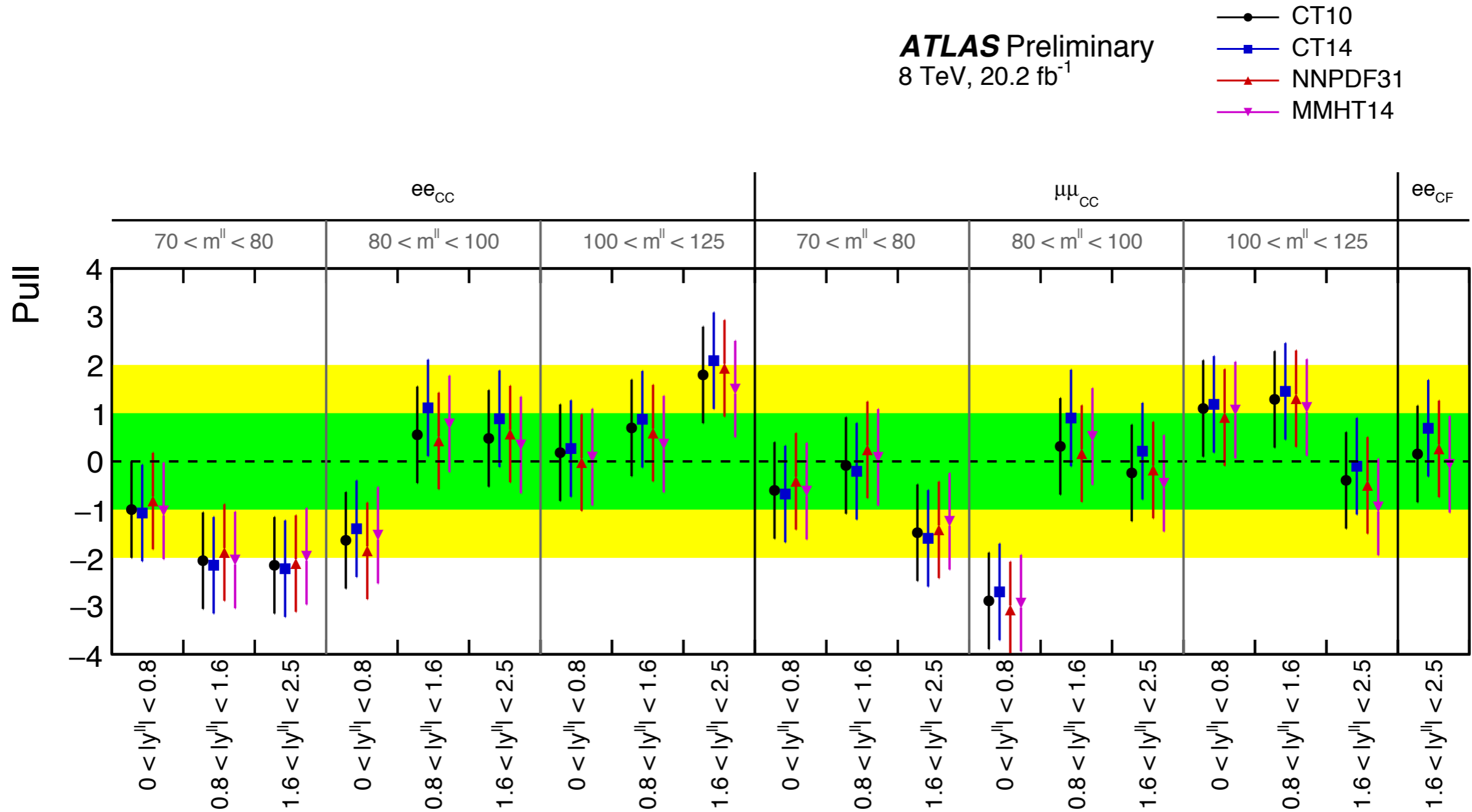


Perform likelihood fits to folded templates on  $m$  &  $|y|$  bins  
 Use event-wise reweighting to vary  $\sin^2\theta_{\text{eff}}$  in templates  
 Like performing analytic interpolation:  
 - known harmonic polynomials fitted to data  
 - reduces PDF sensitivity

Use linear interpolation model to extract  $\sin^2\theta_{\text{eff}}$



Consistency checks: pull of  $\sin^2\theta_{\text{eff}}$  for different data sub-sets



Uncertainties on  $\sin^2\theta_{\text{eff}} \times 10^{-5}$

Channel	$ee_{CC}$	$\mu\mu_{CC}$	$ee_{CF}$	$ee_{CC} + \mu\mu_{CC}$	$ee_{CC} + \mu\mu_{CC} + ee_{CF}$
Central value	0.23148	0.23123	0.23166	0.23119	0.23140
Uncertainties					
Total	68	59	43	49	36
Stat.	48	40	29	31	21
Syst.	48	44	32	38	29
Uncertainties in measurements					
PDF (meas.)	8	9	7	6	4
$p_T^Z$ modelling	0	0	7	0	5
Lepton scale	4	4	4	4	3
Lepton resolution	6	1	2	2	1
Lepton efficiency	11	3	3	2	4
Electron charge misidentification	2	0	1	1	< 1
Muon sagitta bias	0	5	0	1	2
Background	1	2	1	1	2
MC. stat.	25	22	18	16	12
Uncertainties in predictions					
PDF (predictions)	37	35	22	33	24
QCD scales	6	8	9	5	6
EW corrections	3	3	3	3	3

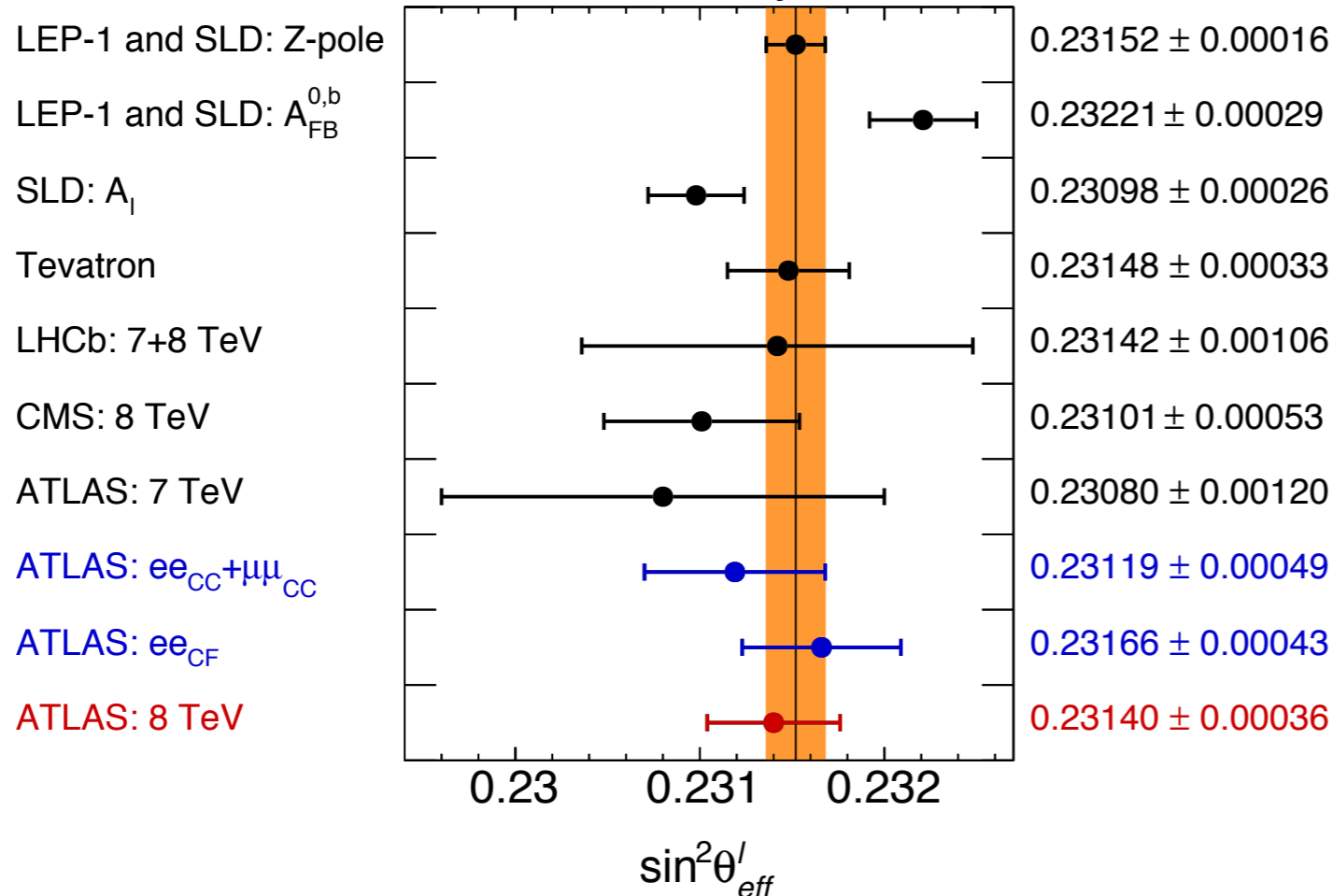
Extracted value / uncertainties of  $\sin^2\theta_{\text{eff}}$  from  $d^3\sigma$  agrees with angular analysis  
 Better precision from CF channel than CC (higher sensitivity / less dilution)  
 Dominated by PDF uncertainty  
 Sizeable uncertainty from data statistics

$$\sin^2\theta_{\text{eff}}^{\ell} = 0.23140 \pm 0.00021 \text{ (stat.)} \pm 0.00024 \text{ (PDF)} \pm 0.00016 \text{ (syst.)}$$

ATLAS reaches precision of single LEP/SLD experiments and combined CDF/D0 precision

	CT10	CT14	MMHT14	NNPDF31	
$\sin^2\theta_{\text{eff}}^{\ell}$	0.23118	0.23141	0.23140	0.23146	
Uncertainties in measurements					
Total	39	37	36	38	
Stat.	21	21	21	21	$\times 10^{-5}$
Syst.	32	31	29	31	

**ATLAS Preliminary**



Preliminary result  
Released at ICHEP 2018

ATLAS determination of  $\sin^2\theta_{\text{eff}}$  is nearing completion  
Timescale - aim for final publication spring 2019  
More detailed validation of DYTurbo vs NNLOjet

cross sections have larger  
PDF sensitivity allowing  
in-situ PDF constraints

angular coefficients reduce  
PDF sensitivity through  
known harmonic polynomials

## Triple Differential cross-section method:

- Use NNLO Z+j predictions in “forbidden region” ?
- use mixed method:
  - fit  $A_{\text{FB}}(m, |y|, |\cos\theta^*|)$  for  $|\cos\theta^*| < 0.4$  &  $m < 66$  GeV
  - fit  $d^3\sigma$  for  $m > 66$  GeV &  $|\cos\theta^*| < 0.4$
  - (we already did this and find PDF uncertainty is reduced!)
- using full  $d^3\sigma$  in fit yields smallest PDF uncertainty
- perform complete NNLO QCD fit (not PDF profiling)

## Angular coefficients method:

- adjust to  $d^3\sigma$  experimental selection
- evaluate statistical uncertainty with bootstraps
- a few experimental checks to complete (SFs etc)
- PDF profiling tests
- PDF reweighting tests
- include  $A_3$  ?

**Project is actively pursued in LPCC Electroweak Working Group: ATLAS / CMS / LHCb / Theory**

Much to be gained from LHC combination

- LHCb has higher  $y$  acceptance (but lower luminosity)
- CMS measurement has no ‘forward’ acceptance but complementary central channel

Now have  $150 \text{ fb}^{-1}$  of data at  $\sqrt{s}=13$  TeV  $\rightarrow$  factor 15 higher statistical sample (incl factor 2 from cross section)  
...but larger  $\sqrt{s}$  means lower  $x$   $\rightarrow$  worse dilution





# Double-differential $Z/\gamma^*$ Cross Sections $\sqrt{s} = 8$ TeV

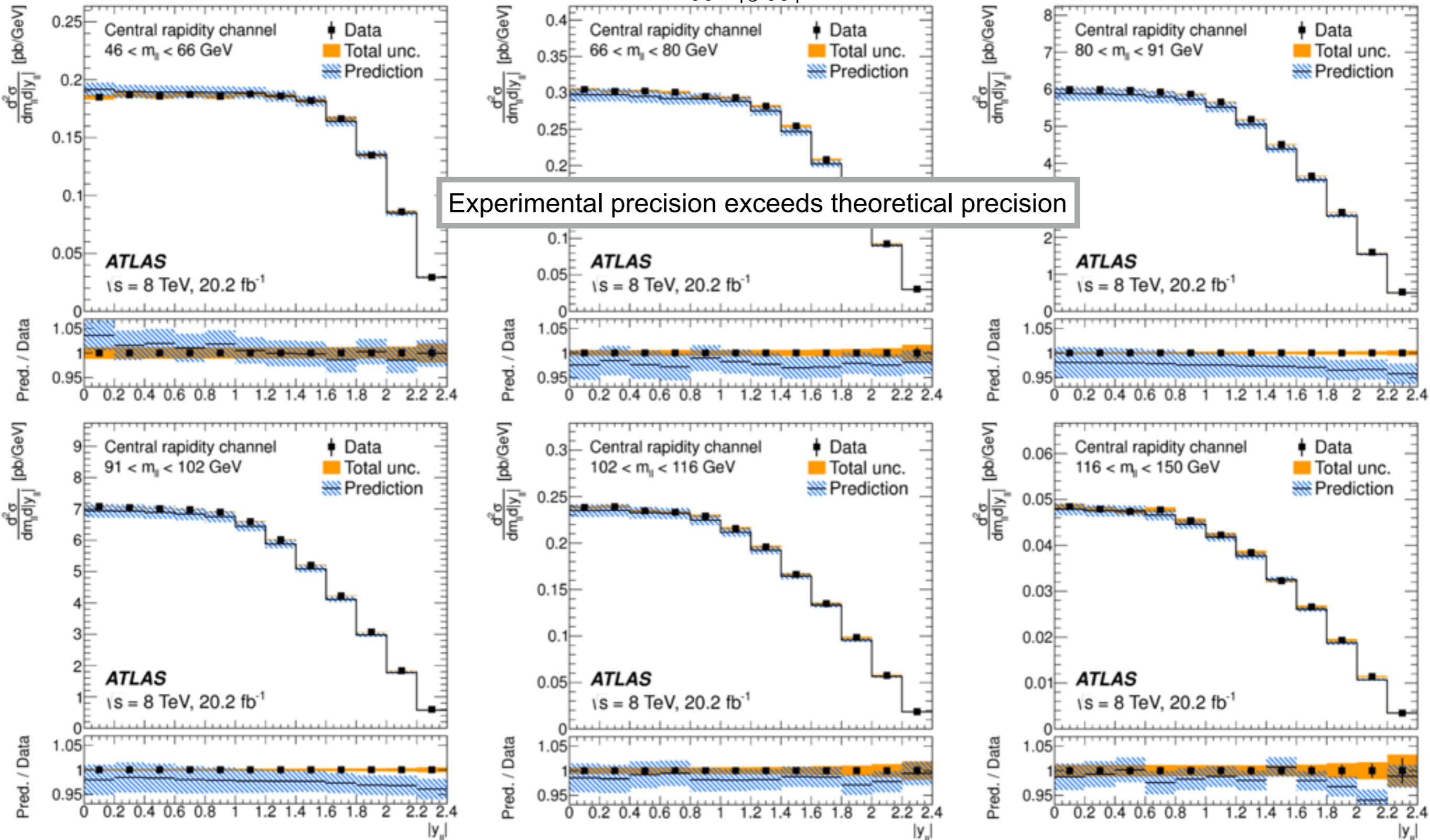


Integrated double-differential cross section

$$\frac{d^2\sigma}{dm_{\ell\ell}d|y_{\ell\ell}|}$$

electron & muon channel combination

$\chi^2/\text{ndf} = 103.4/84$



# Double-differential $Z/\gamma^*$ Cross Sections $\sqrt{s} = 8$ TeV

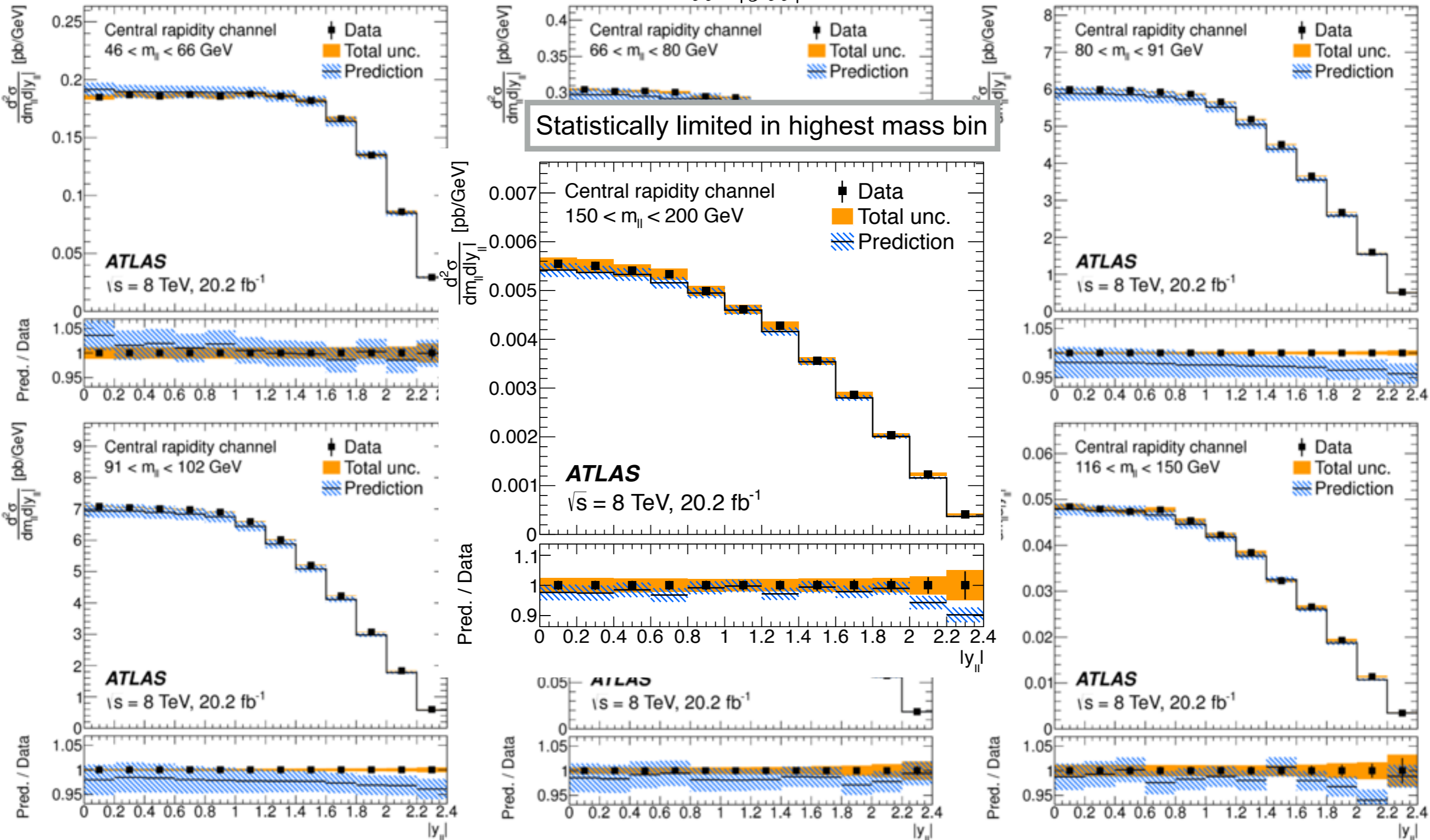


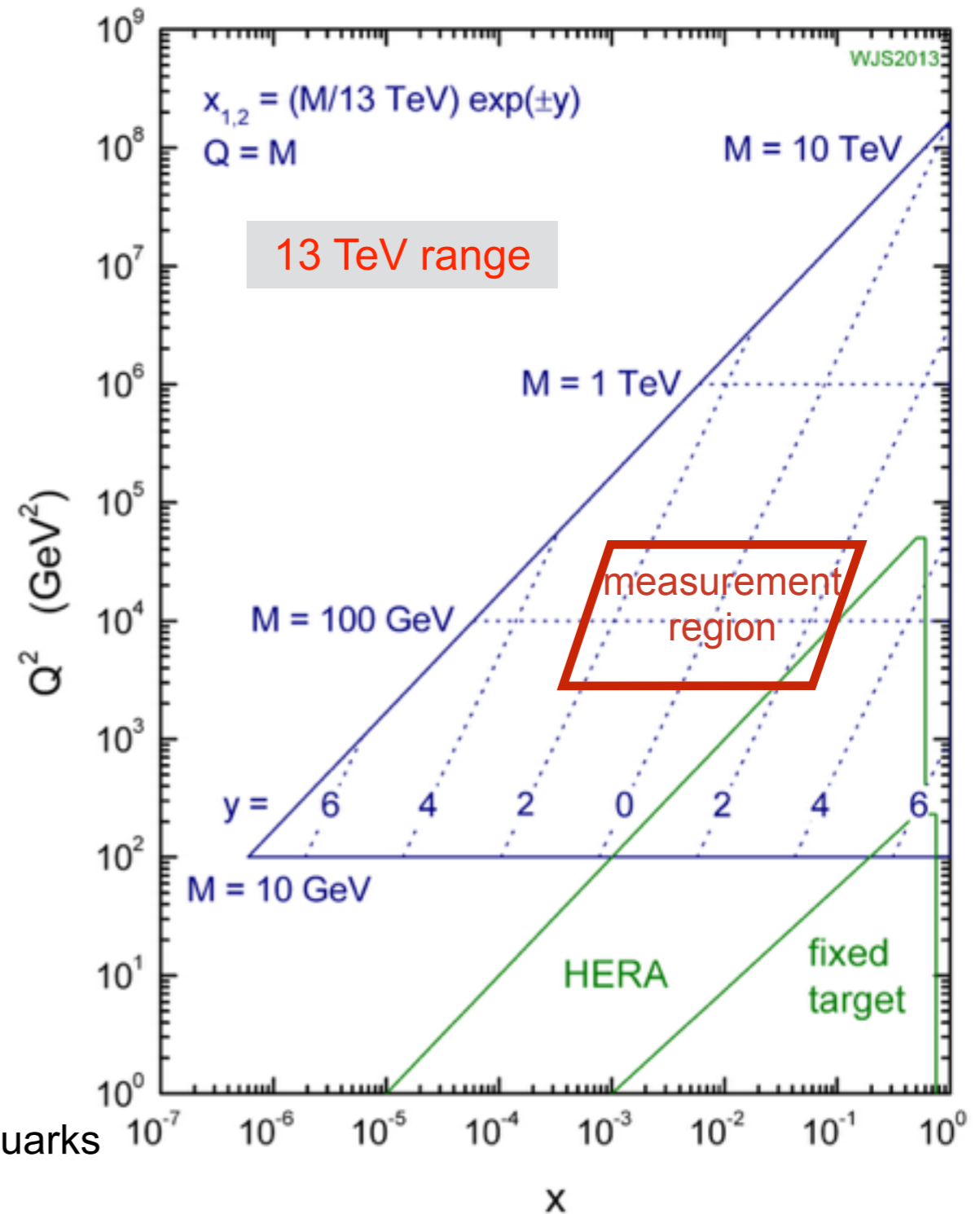
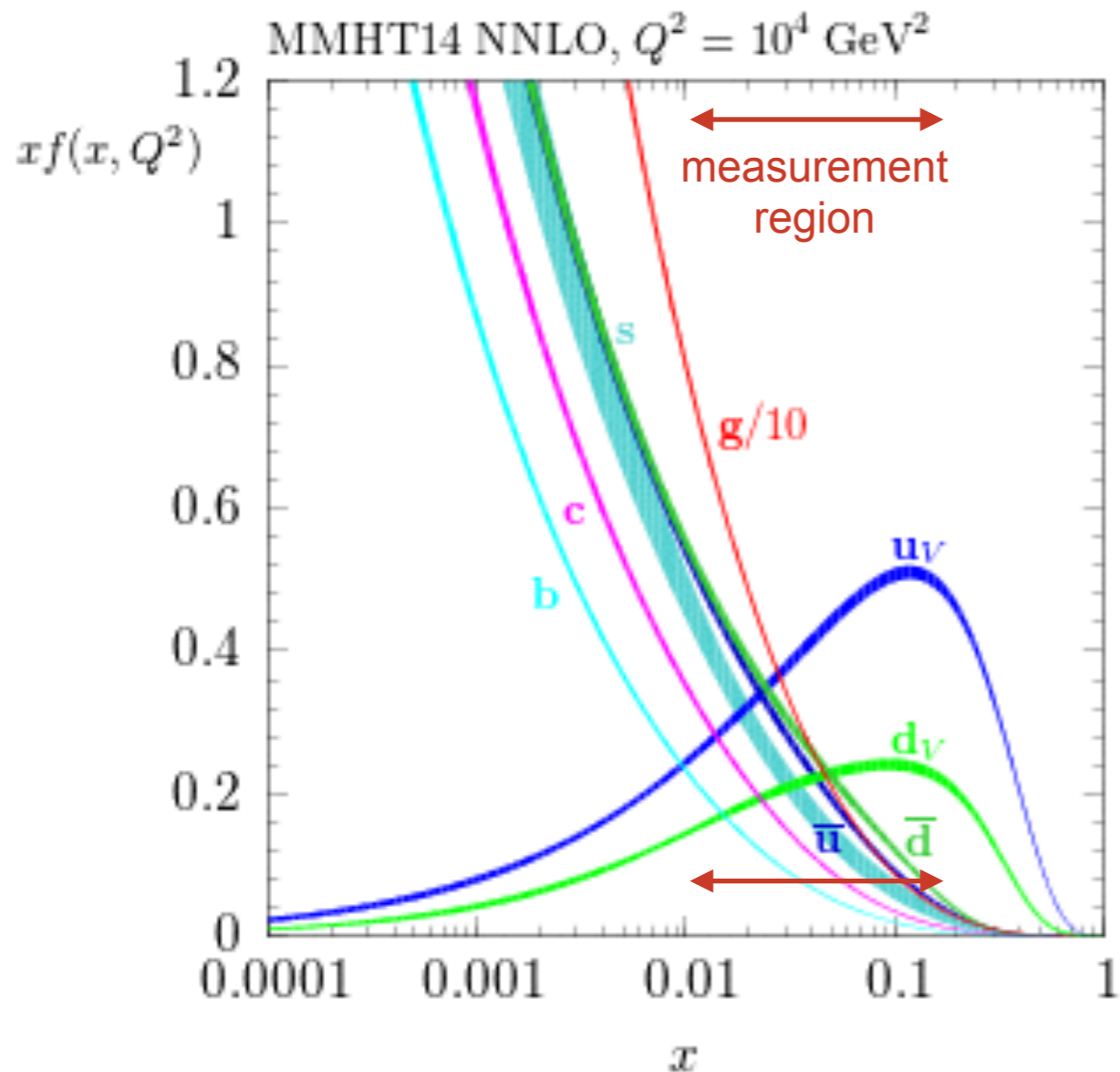
Integrated double-differential cross section

$$\frac{d^2\sigma}{dm_{ee}d|y_{ee}|}$$

electron & muon channel combination

$\chi^2/\text{ndf} = 103.4/84$





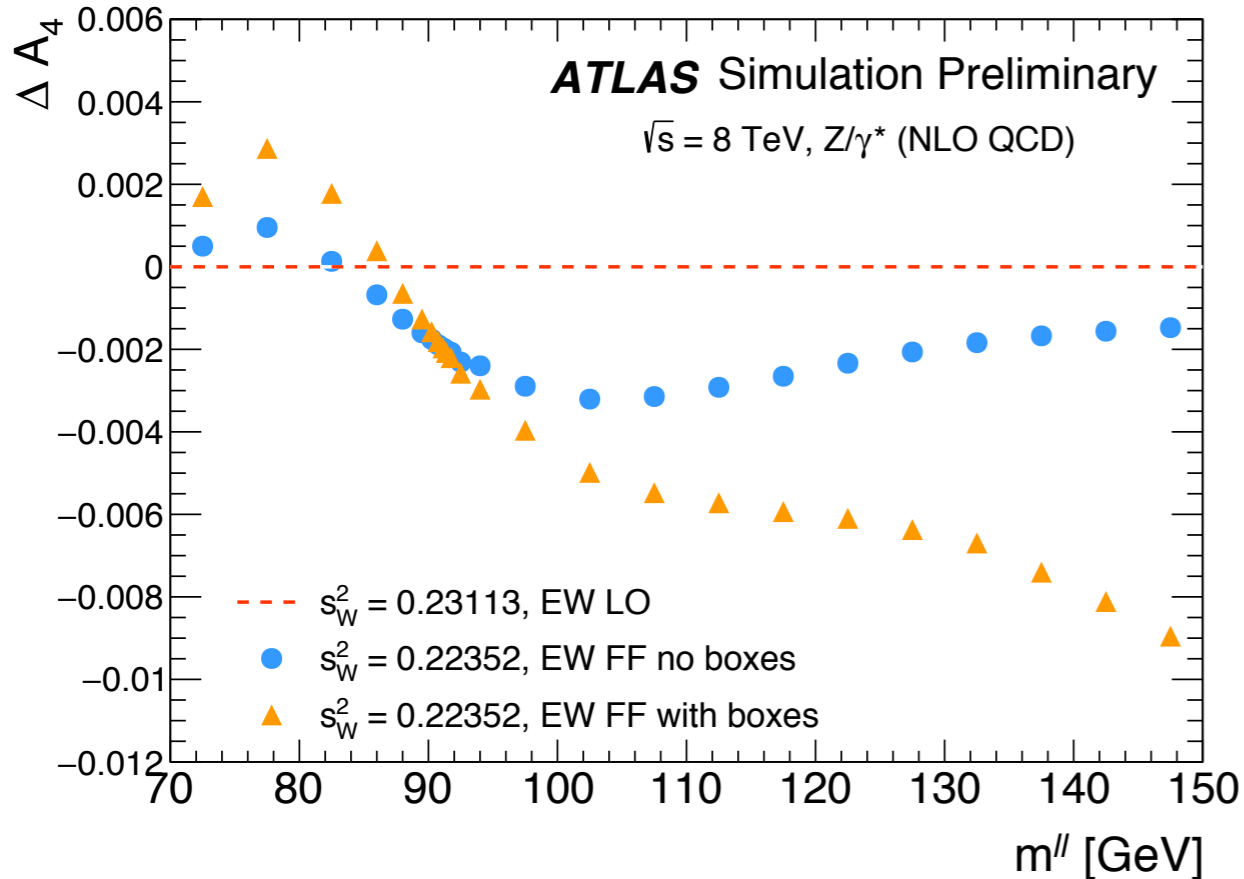
In pp Drell-Yan collisions we do not know direction of incoming quark  $\rightarrow$  ambiguity in defining  $\cos\theta^*$

Rely on valence quarks! At high x quarks dominate not anti-quarks

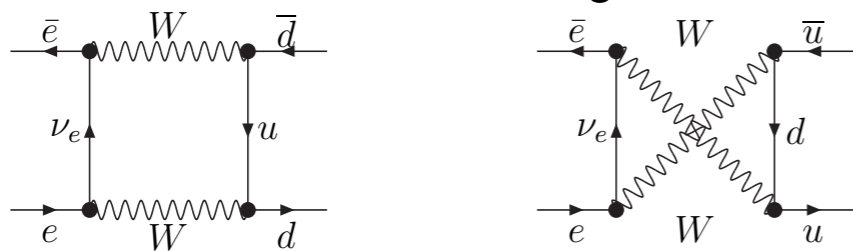
$$x_{1,2} = \frac{m}{\sqrt{s}} \exp^{\pm y} \quad \text{Large } |y| \rightarrow \text{large } x \text{ less dilution}$$

parton momentum fraction of proton

## Change to $A_4$ using NLO corrections



## weak boson box diagrams



NNLO QCD predictions determined using LO EW theory

Close to Z pole:

QED corrections can be factorised from higher order EW corrections

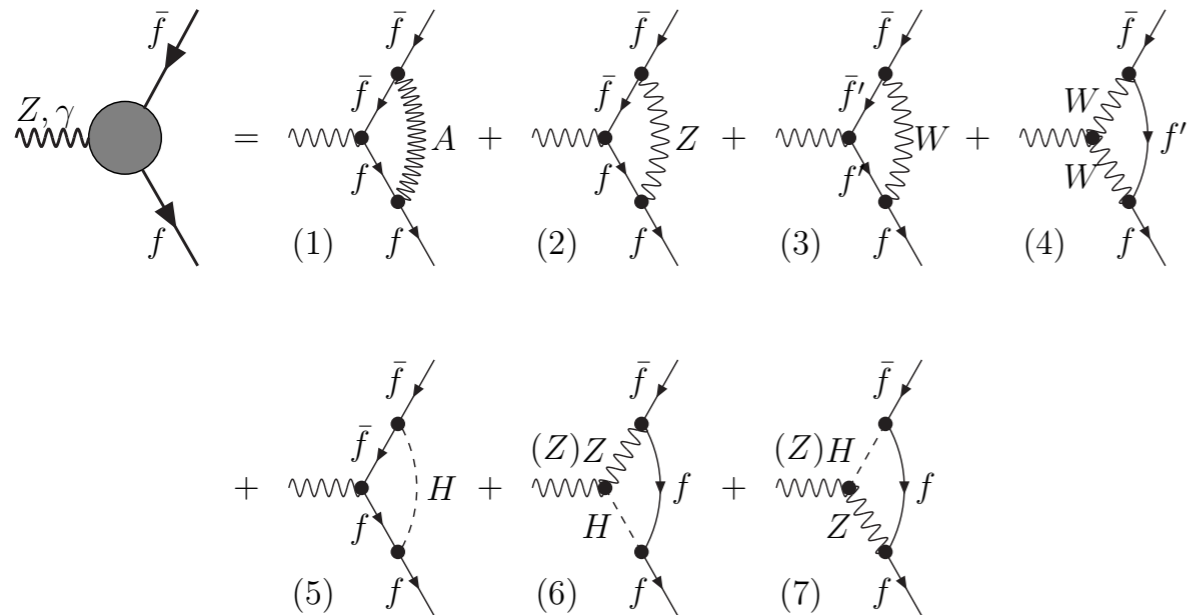
Improved Born Approximation absorbs NLO EW effects into form factors

Initial / final state QED/QCD radiative effects are factorised

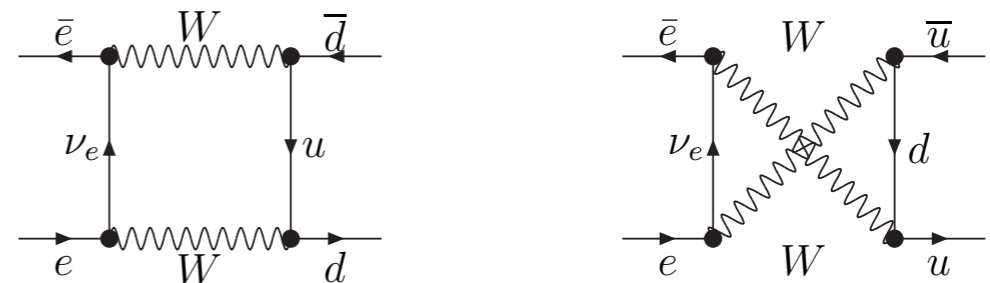
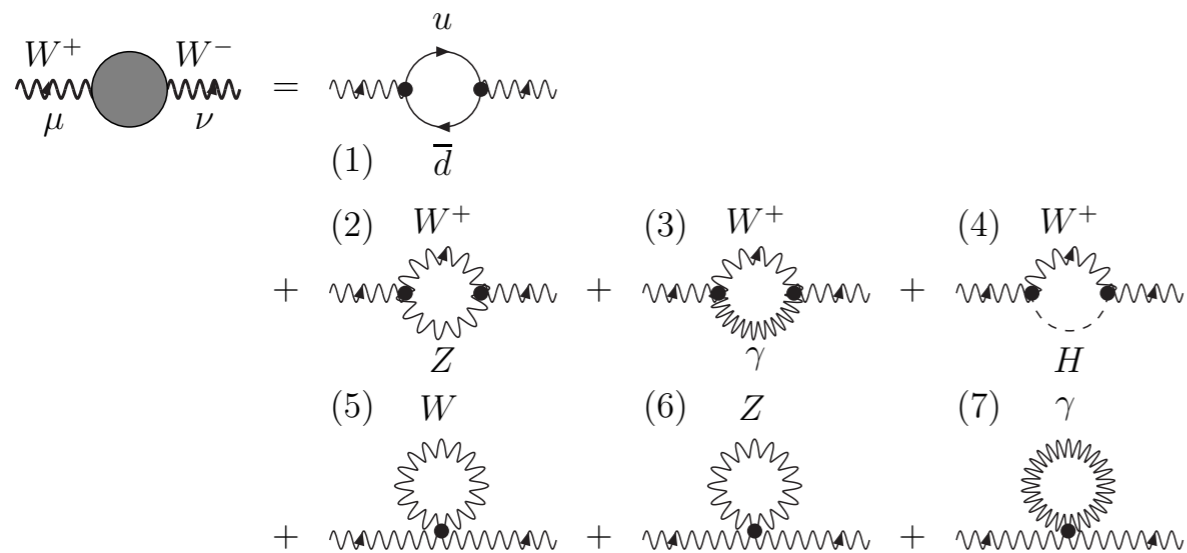
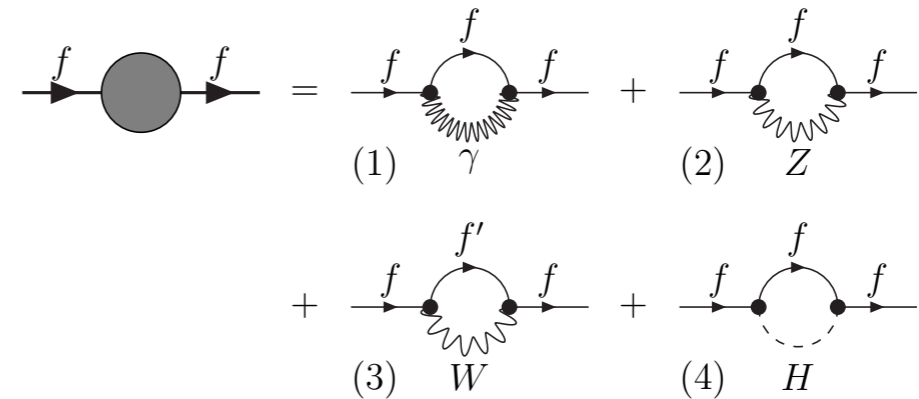
calculation performed using DIZET library 6.21

Parameter	Value	Description
Measured		
$m_Z$	91.1876 GeV	Mass of Z boson
$m_H$	125.0 GeV	Mass of Higgs boson
$m_t$	173.0 GeV	Mass of top quark
$m_b$	4.7 GeV	Mass of b quark
$1/\alpha(0)$	137.0359895(61)	QED coupling constant in Thomson limit
$G_\mu$	$1.166389(22) \cdot 10^{-5} \text{ GeV}^{-2}$	Fermi constant from muon lifetime
Calculated		
$m_W$	80.353 GeV	Mass of W boson
$\sin^2 \theta_W$	0.22351946	On mass-shell-value of weak mixing angle
$\alpha(m_Z^2)$	0.00775995	
$1/\alpha(m_Z^2)$	128.86674175	
$ZPAR(6) - ZPAR(8)$	0.23175990	$\sin^2 \theta_{eff}^\ell(m_Z^2)$ ( $e, \mu, \tau$ )
$ZPAR(9)$	0.23164930	$\sin^2 \theta_{eff}^u(m_Z^2)$ (up quark)
$ZPAR(10)$	0.23152214	$\sin^2 \theta_{eff}^d(m_Z^2)$ (down quark)

## Z → ff corrections



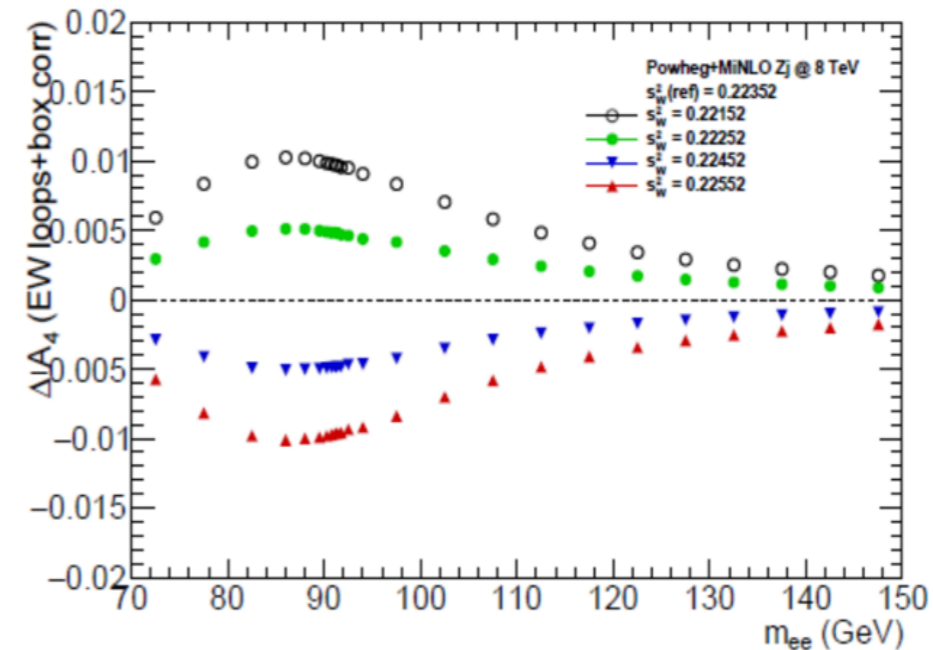
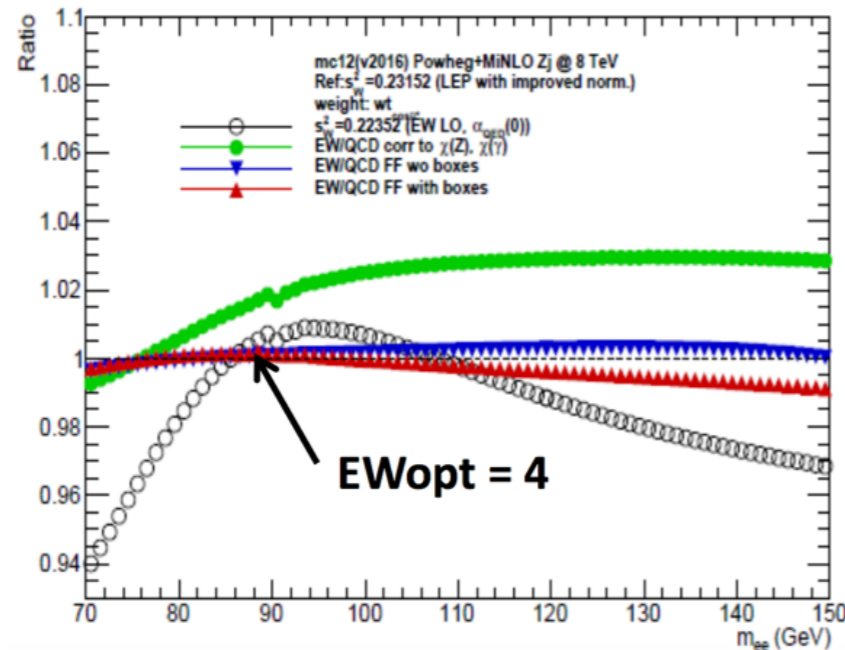
## fermionic self-energy corrections



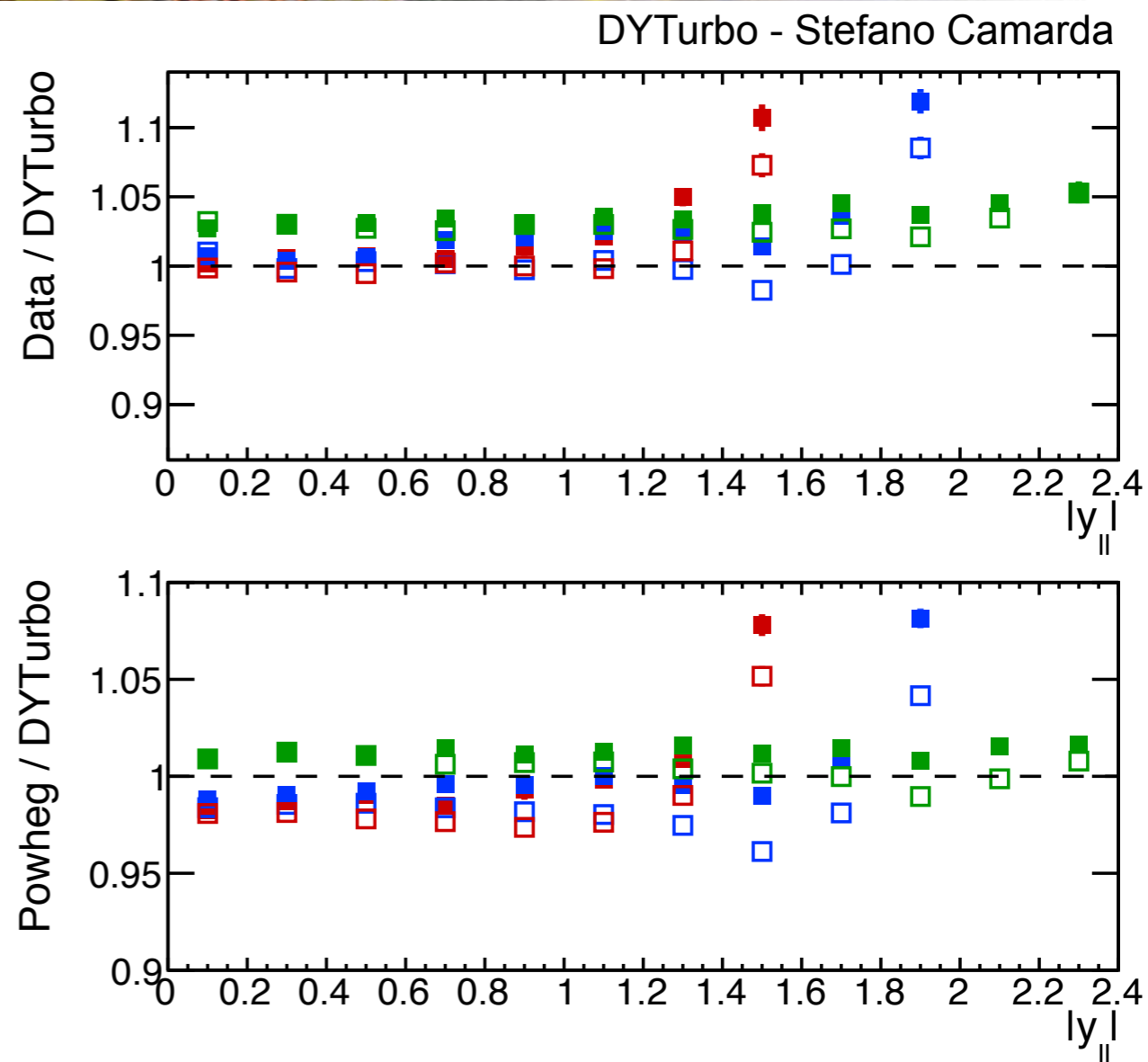
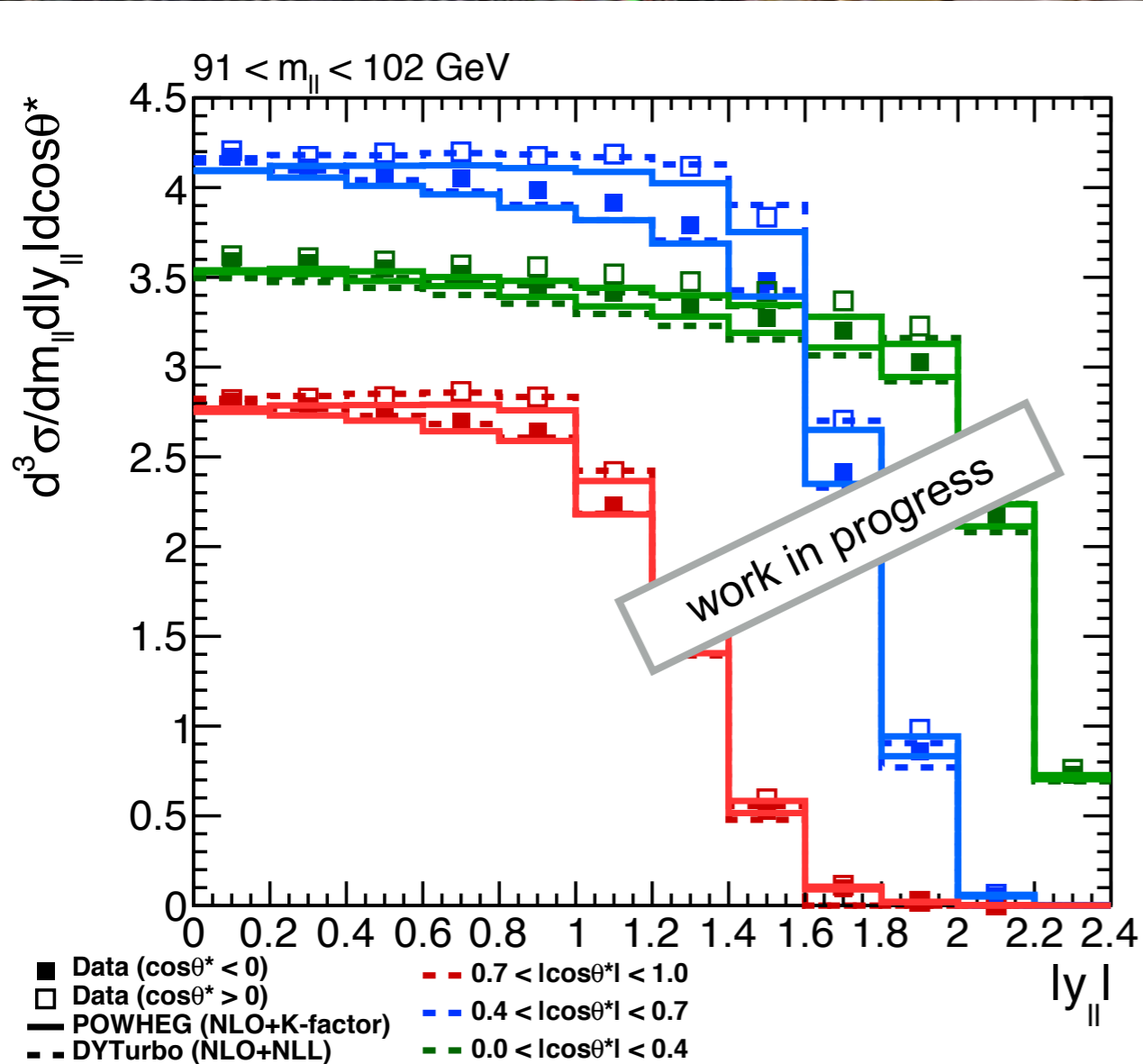
## boson self-energy corrections

## W and Z box diagrams

## EW Corrections



- Existing MC samples used for analysis are missing higher order EW corrections
- Factorize gauge invariant set of EW corrections from QCD, and interface with existing MC samples via “after-burn” approach
  - Interface to DIZET and KKMC libraries adapted to pp collisions, developed for LEP, to compute EW form factors
    - Exact  $O(\alpha)$  + higher order terms
    - Dependent on event kinematics  $s, t = s*(1-\cos\theta)/2$
  - Insert as event weights in MC sample
- Weights can also be embedded for effective (LO) EW scheme
  - Difference between this and EW FF is quoted to be  $\sim 22*10^{-5}$  for Tevatron and CMS (studies ongoing to confirm)
- Allows us to study EW effects at the per-mil level, and scan  $\sin^2\theta_w$  within a single MC sample
- Studies to be done cross-checking with PowhegEW generator
- More detailed info [here](#). **Will have dedicated talk in next meeting from Elzbieta!**



Initial comparisons of DYTURBO (NLO+NLL) with Powheg (NLO x [NNLO  $\otimes$  NLO EW] k-factor)

Prediction code needs tuning / optimisation for:

- integration time & precision for fiducial  $d^3\sigma$
- large QCD scale  $\mu_R$  &  $\mu_F$  dependence observed in some kinematic regions
- optimisation of resummation scale  $\mu_{\text{Resum}}$  in NLL

Could indicate improved resummation is needed (move to NNLL?)





New ATLAS method:

Perform QCD & EW fit to  $d^3\sigma$  cross sections differential in  $m, |y|, \cos\theta^*$

Ai - Angular coefficient analysis (methodology used here [arXiv:1606.00689](https://arxiv.org/abs/1606.00689))

Target precision on  $\sin^2\theta_{\text{eff}}$  about  $30 \times 10^{-5}$  (total uncertainty) :

- use large  $20 \text{ fb}^{-1}$  luminosity data sample at  $\sqrt{s}=8\text{TeV}$
- include FCAL forward electron kinematic region - better sensitivity
- use unfolded  $d^3\sigma$  to gain PDF sensitivity
- perform simultaneous fit to PDFs and  $\sin^2\theta_{\text{eff}}$  on same data

Method combine best NNLO QCD & NLO EW predictions

Use xFitter framework to perform  $\chi^2$  fits

Use method of PDF profiling to optimise PDF eigenvalues [arXiv:1402.6623](https://arxiv.org/abs/1402.6623)

Account for correlated experimental systematics

Scan for optimum value of  $\sin^2\theta_{\text{eff}}$

Ingredient list: State-of-the-art fiducial QCD predictions  $\otimes$  NLO EW corrections

NLO ( + NLL ?)

$d^3\sigma$  measurement is inclusive in  $p_{T,Z}$

NNLO ( + NNLL ?)

but resummations may be important in some kinematic regions

Toolkits:

DYTurbo : (NLO + NLL) or (NNLO + NNLL) resummation for small  $p_{T,Z}$  predictions

MCFM: NLO

DYres:

Powheg: NLO + PS

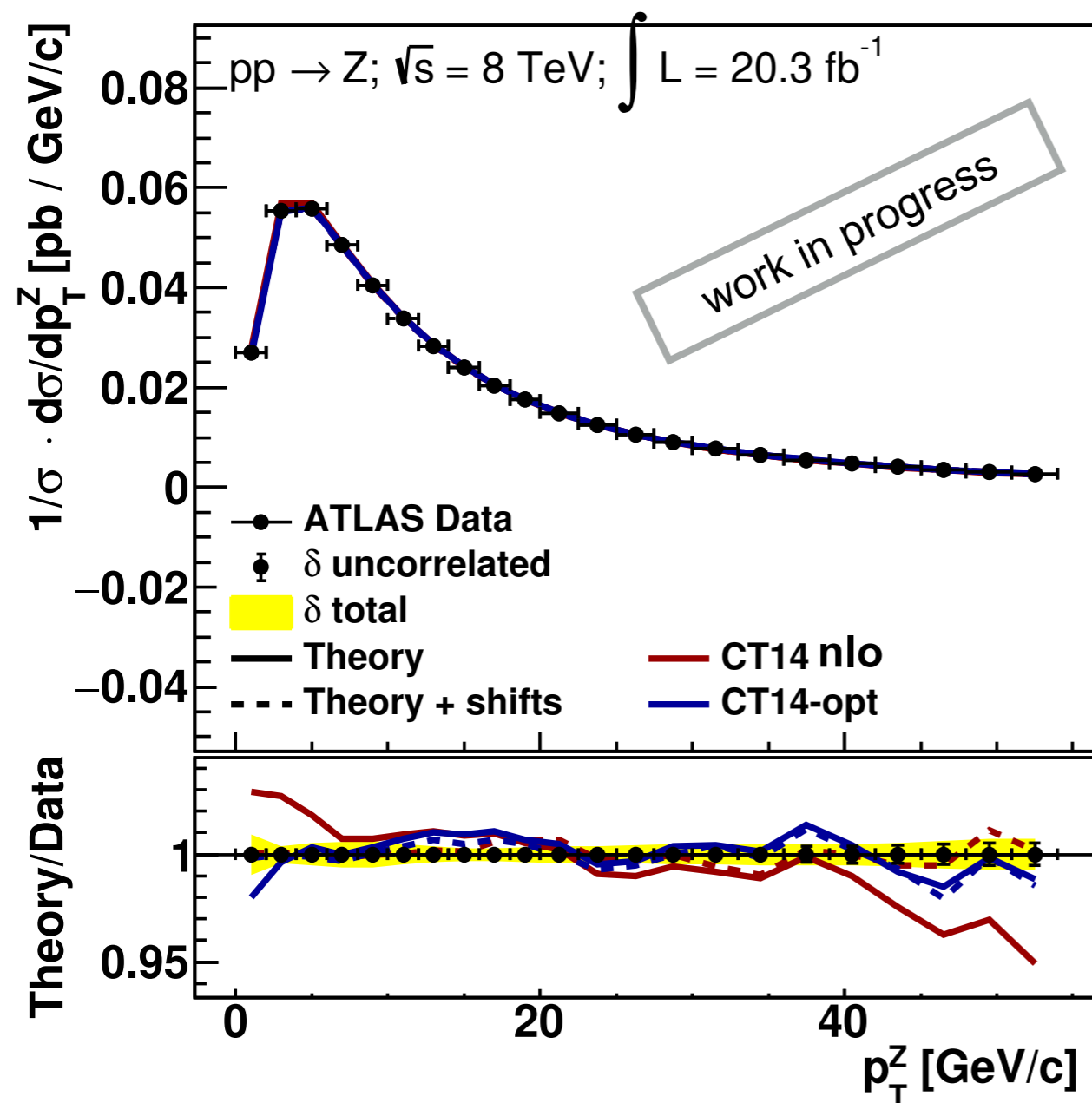
NNLOjet: ??

This is very much work-in-progress !

Can illustrate status since  $d^3\sigma$  cross sections are published



DYTurbo - Stefano Camarda  
 xFitter - Sasha Glazov & co



Tune resummation calculation on ATLAS 8 TeV Z  $p_T$  data

- non-perturbative parameter  $g$
- $\mu_R$  &  $\mu_F$  &  $\mu_{\text{Resum}}$

### Default settings

$$\mu_R = \mu_F = \mu_{\text{Resum}} = 0.5 m_{\parallel}$$

$$g = 1.0 \text{ GeV}^2$$

### Initial optimisation prediction

$$\mu_R = 0.34 \times m_{\parallel}$$

$$\mu_F = 0.49 \times m_{\parallel}$$

$$\mu_{\text{Resum}} = 0.41 \times m_{\parallel}$$

$$g = 1.04 \text{ GeV}^2$$

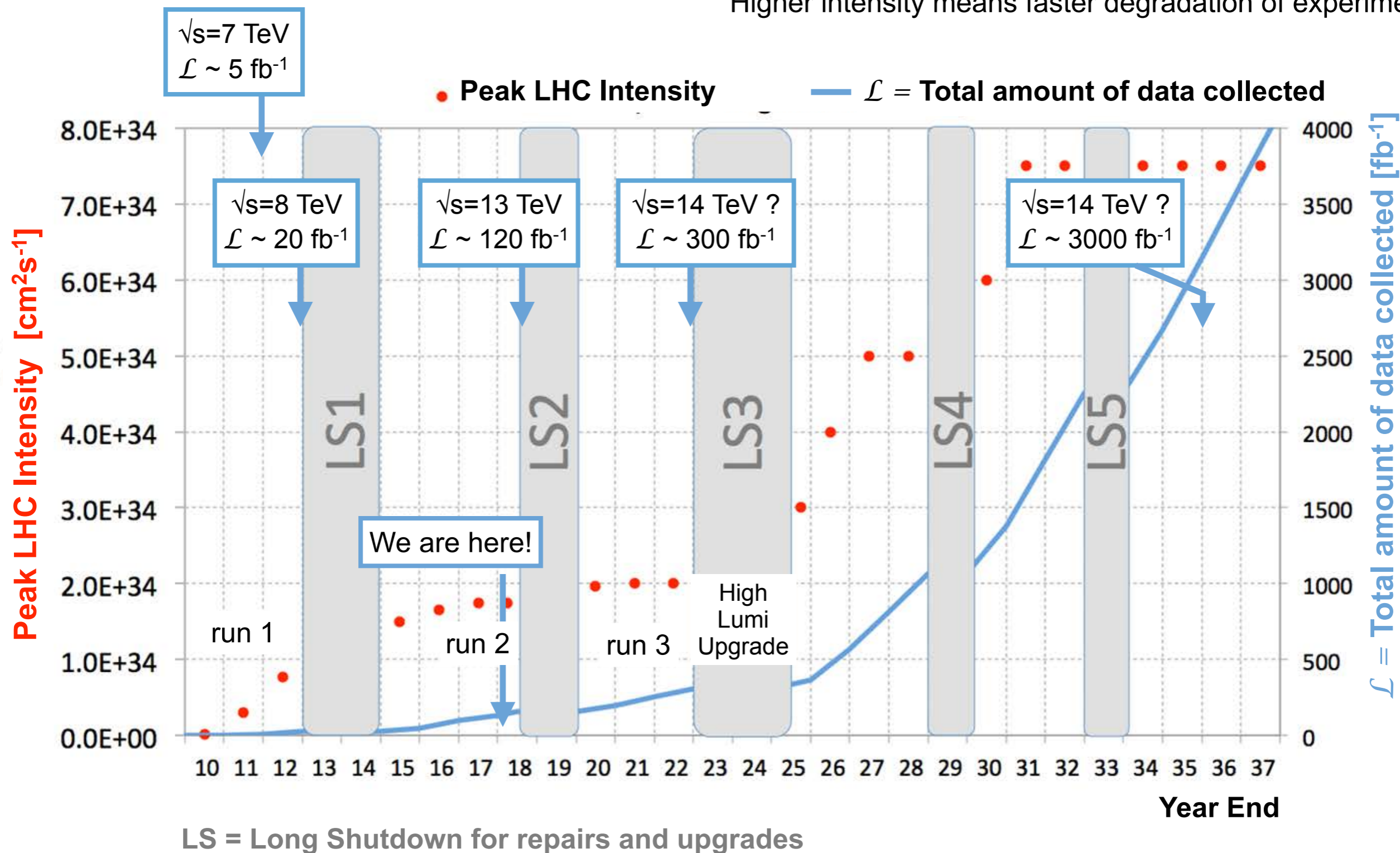
Alternatively - switch to using  $A_{\text{FB}}$  in some regions where scale errors are large?

# LHC Schedule to 2035



\* actual schedule slipped by 1 year  
e.g. LS3 starts 2023

Large increases in intensity  
Requires significant changes to LHC magnets  
Higher intensity means faster degradation of experiments

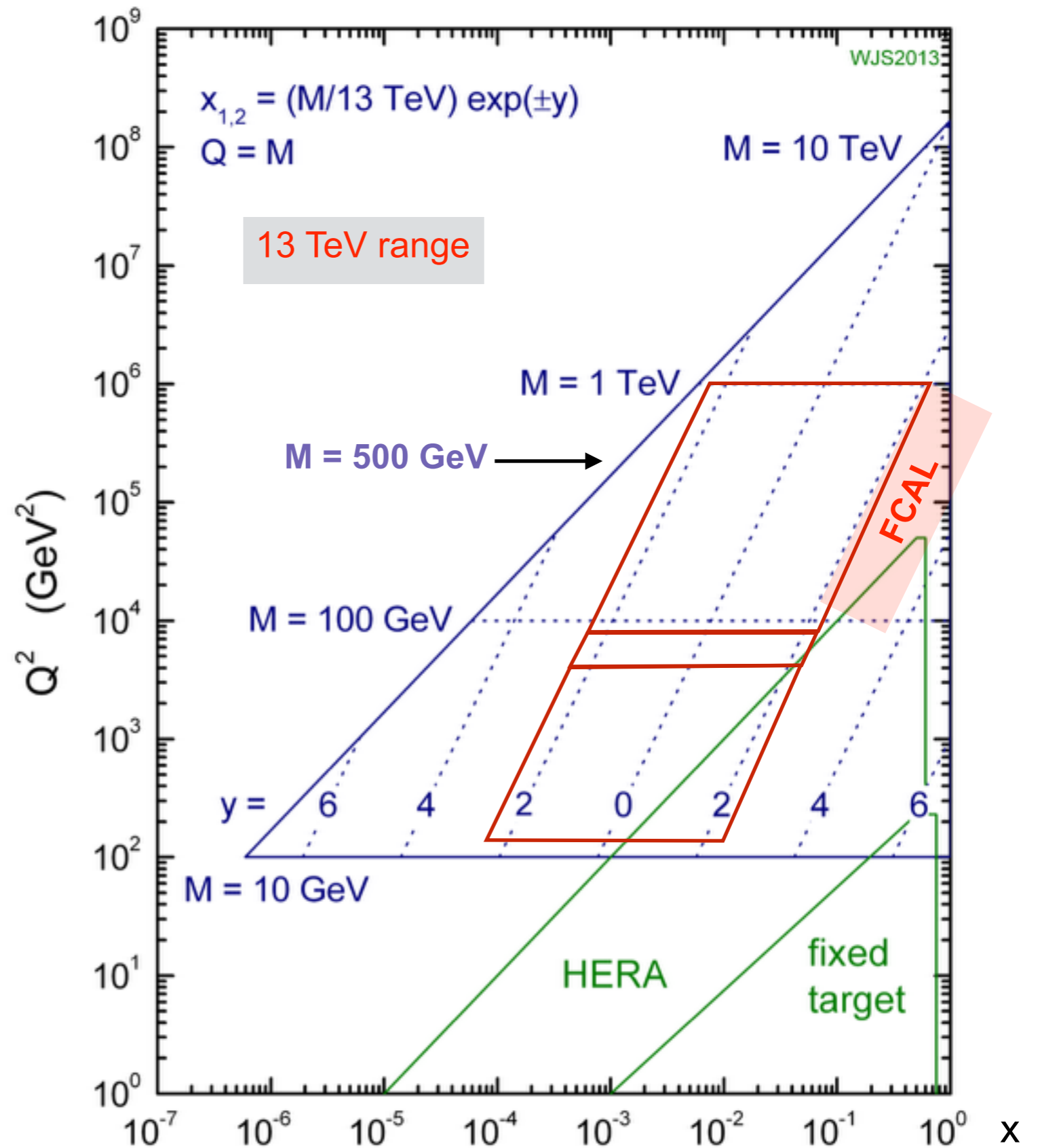




Classic problem: how to constrain PDFs at high x for BSM searches?

Measure cross sections at high rapidity

FCAL forward electrons  $\rightarrow$  PDF sensitivity up to  $x=1$  at  $m=500$  GeV





General models of new physics SM Lagrangian extended by dimension 6 operators

They describe new physics appearing at scale  $m > \sqrt{s}$

- ★ new EW vector bosons
- ★ new EW fermions
- ★ EW compositeness...

<https://arxiv.org/abs/1609.08157>

Effective field theory (EFT) attempts to encapsulate this  
For DY production 4 propagator form-factors introduced:

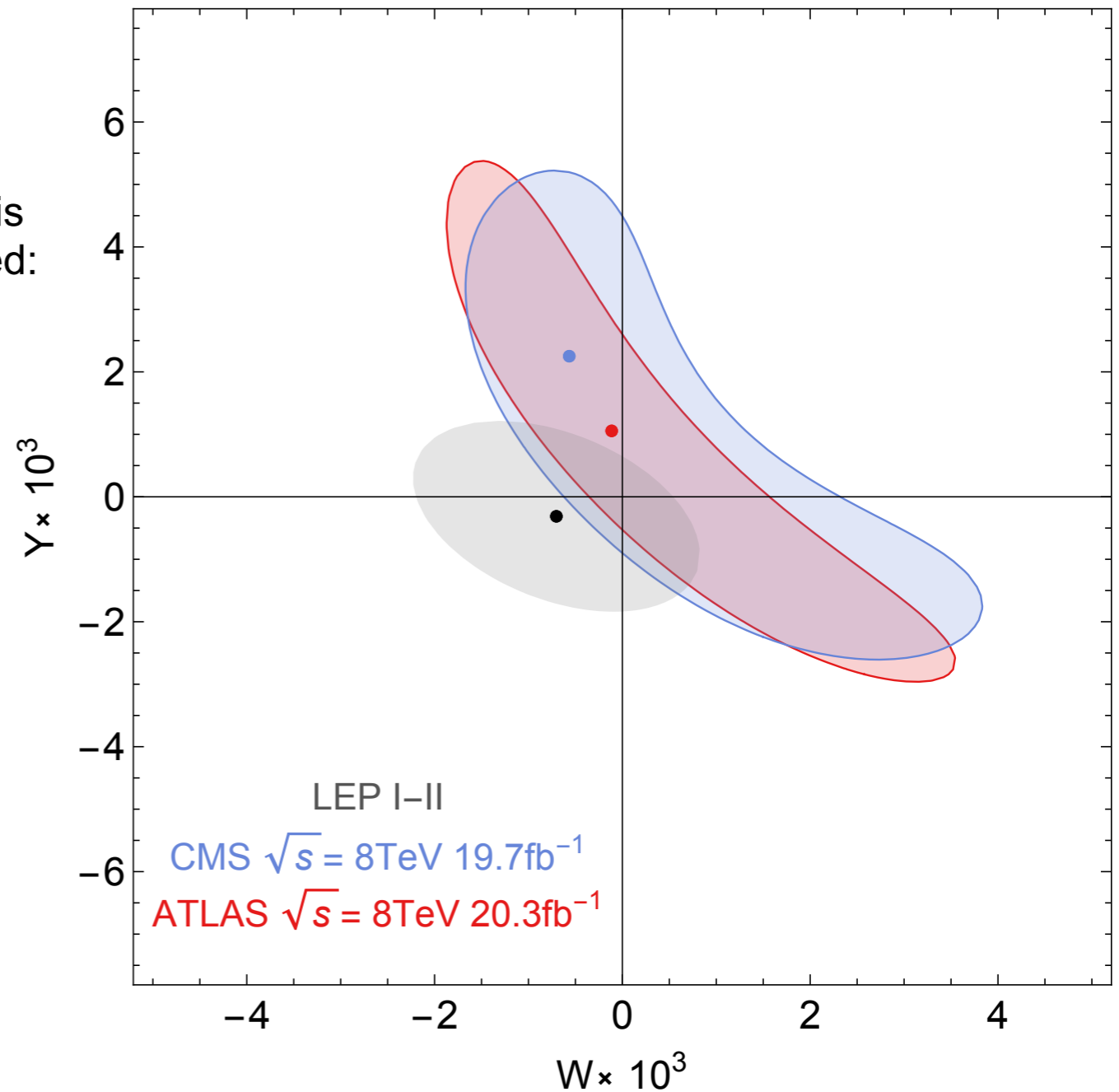
S, T, Y, W

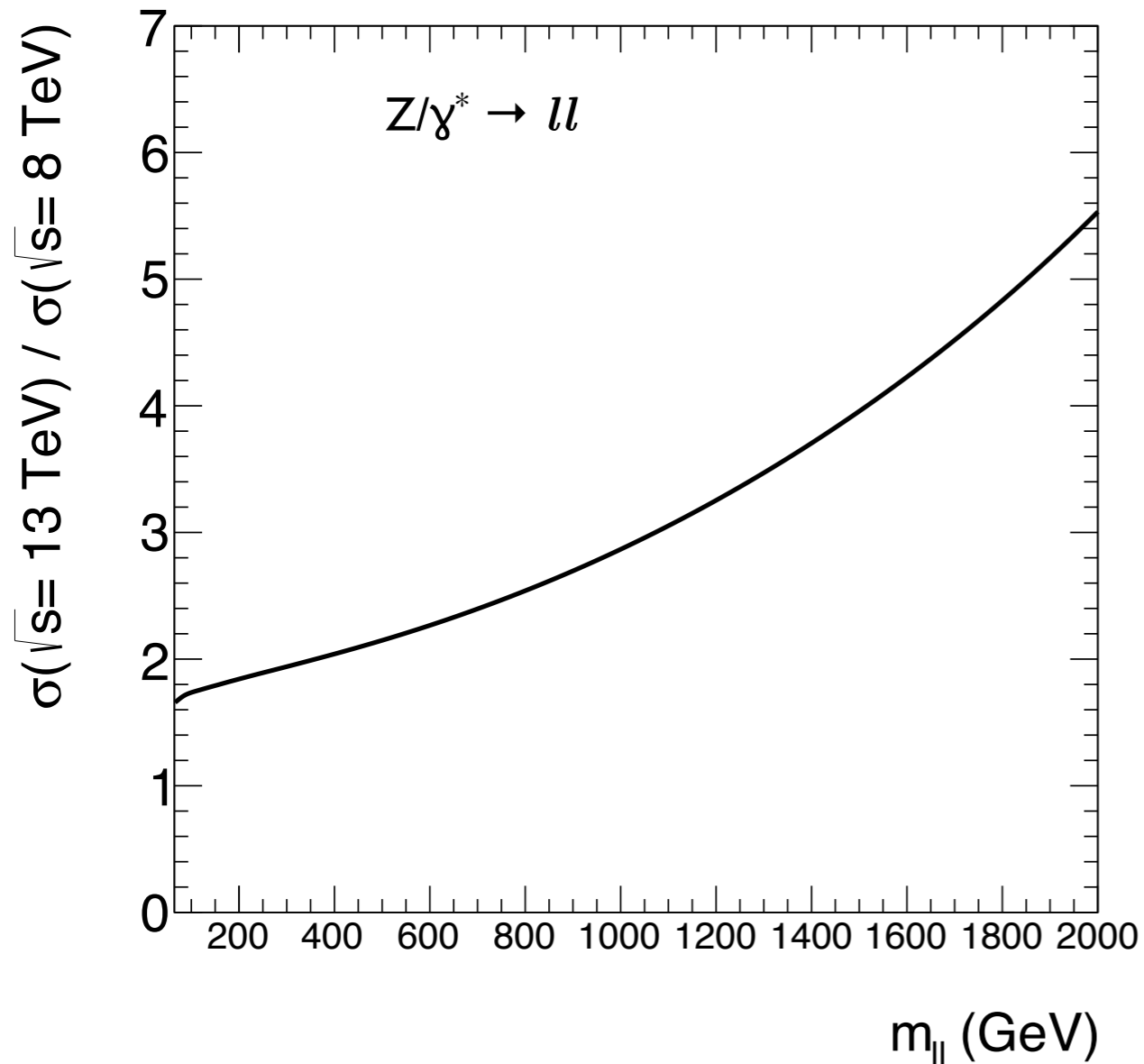
- Y and W increase with  $\sqrt{s}$
- S and T do not grow with  $\sqrt{s}$

LHC data can help constrain Y & W

Current constraints based on neutral current HMDY 8 TeV data

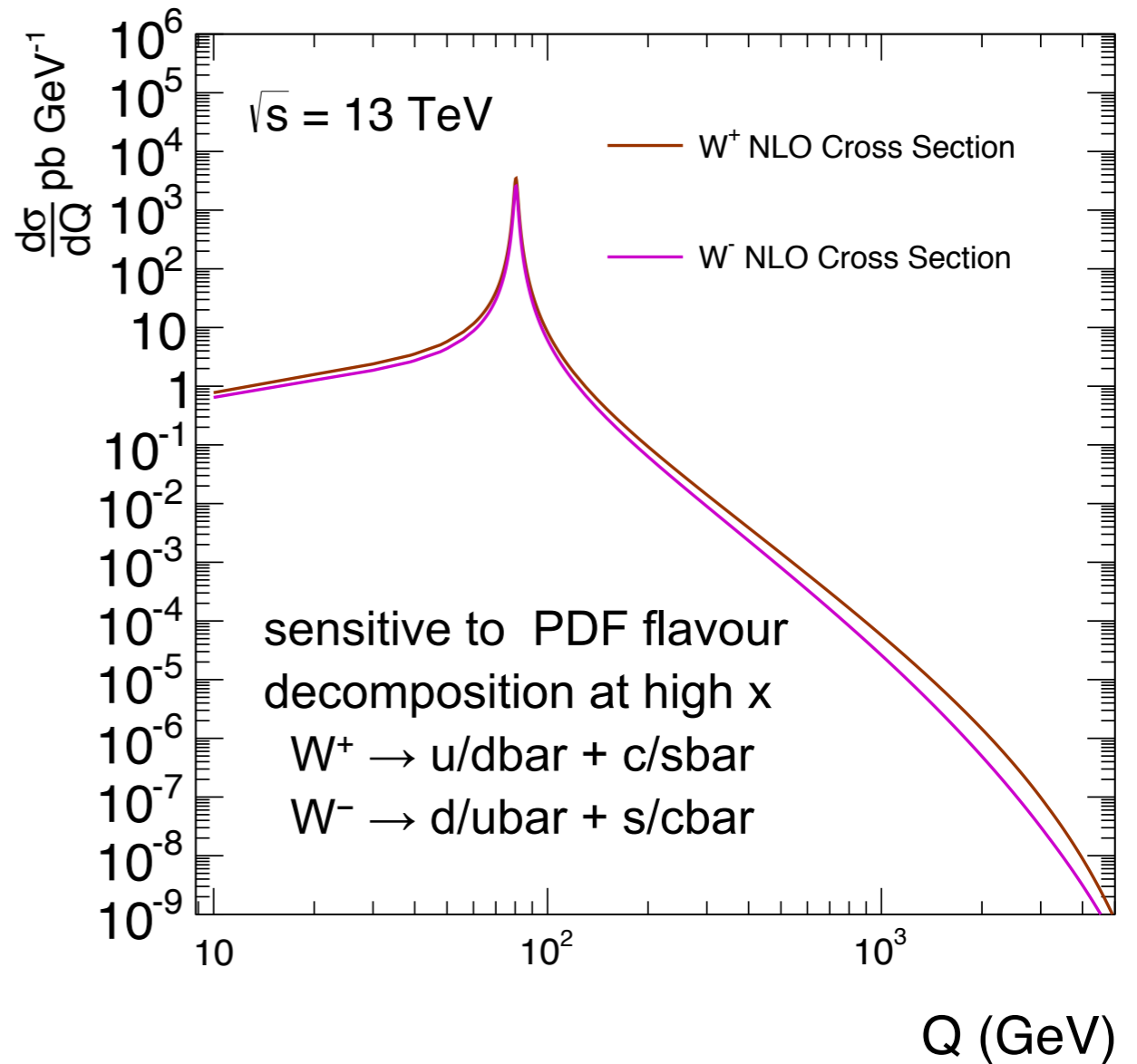
⇒ Cannot yet compete with LEP





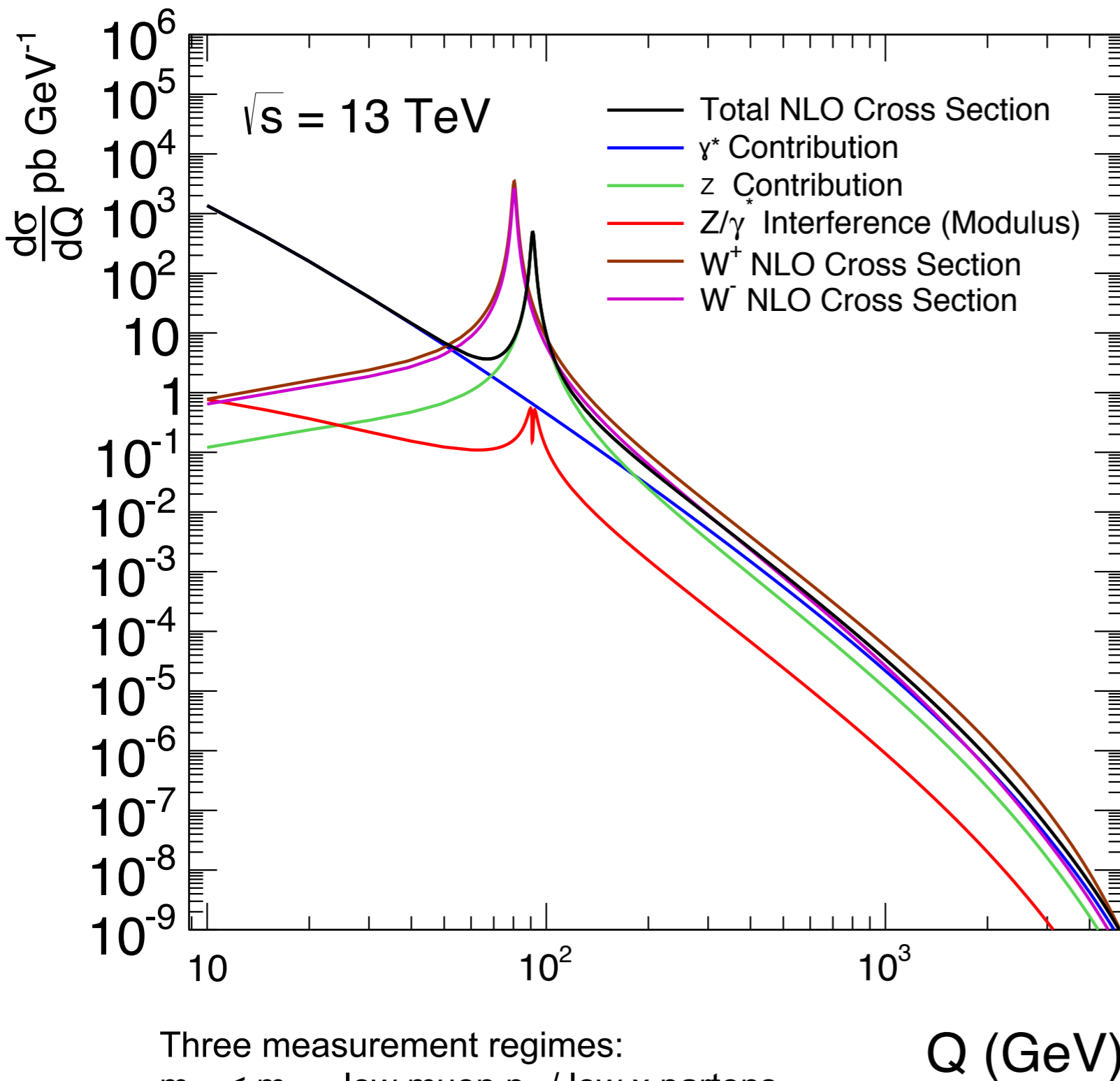
## Neutral current

Cross section enhancement > factor 5 at large  $m_{||}$   
 Similar for charged current



## Charged current

First measurement off-shell high  $m_T$   $W^\pm$  production  
 Analogous to neutral current  $Z/\gamma^*$  measurement



- At large Q  $\sigma(W^+) > \sigma(W^-) \geq \sigma(\gamma^*)$  by  $\sim$  factor 2
- Run-II total  $\int L \sim 120$  fb $^{-1}$
- Lumi  $\sim$  4-5 times larger than Run-I
- Factor  $>2$  larger cross section at 13 TeV  $\Rightarrow$  order of magnitude more data

High mass DY reaches high x region

Factor 5 higher x than on-shell Z at 8 TeV

At  $M=300-500$  can achieve  $\sim 2\%$  precision for  $|y| < 1$

Three measurement regimes:

- $m_{\mu\mu} < m_Z$  – low muon  $p_T$  / low x partons
- $m_{\mu\mu} = m_Z$  – ultra-high precision
- $m_{\mu\mu} > m_Z$  – high muon  $p_T$  / new physics / high x partons



Stringent constraints on  $Y$  &  $W$  from LEP  
 100 fb<sup>-1</sup> of NC data  $Z/\gamma^* \rightarrow l^+l^-$  reaches LEP precision  
 20 fb<sup>-1</sup> of CC data  $W \rightarrow lv$  surpasses LEP by factor 4!

<https://arxiv.org/abs/1609.08157>

Discussions with Andrea / Riccardo et al  
 Request for unfolded cross sections  
 Additional gains in NC channel measuring decay angles  
 $\cos \theta^*$   
 $y_{ll}$   
 $m_{ll}$   
 → triple differential cross sections

Started analysis of high mass DY cross sections  
 in run-II @  $\sqrt{s}=13$  TeV

Simultaneous measurement in NC & CC channels

